Dynamical coupled-channels study of S = -1 hyperon resonances

Hiroyuki Kamano (RCNP, Osaka U.)

[HK, Nakamura, Lee, Sato, PRC90(2014)065204; arXiv:1506.01768 (to appear in PRC)]

KEK理論センター J-PARC分室、JAEA先端基礎研究センター共催研究会 「ストレンジネス核物理の発展方向」, KEK東海キャンパス, Aug. 3-5, 2015

Introduction: Baryon spectroscopy

Discovery of the Δ baryon (1952)

Total Cross Sections of Positive Pions in Hydrogen*

H. L. ANDERSON, E. FERMI, E. A. LONG,[†] AND D. E. NAGLE Institute for Nuclear Studies, University of Chicago, Chicago, Illinois (Received January 21, 1952)



FIG. 1. Total cross sections of negative pions in hydrogen (sides of the rectangle represent the error) and positive pions in hydrogen (arms of the cross represent the error). The cross-hatched rectangle is the Columbia result. The black square is the Brookhaven result and does not include the charge exchange contribution.



PDG2014 reports ~ 150 baryons. (NOTE: Existence is still uncertain for ~ 50% of baryons.)

Baryon Spectroscopy: Understanding nature of baryons and their excitations

- Mass, width, spin, parity ...?
- Internal structure (form factors)?
- How produced in reaction processes?
- How interact with other particles?



Approaches to light-quark baryon $(N^*, \Delta^*, \Lambda^*, \Sigma^*)$ spectroscopy



Current situation of Y*(= Λ^* , Σ^*) **spectroscopy**

PDG listing

	Λ*			Σ*		
\checkmark Y [*] (= Λ [*] , Σ [*]) resonances are much less understood than N* & Λ*	Particle	J^P	Overall status	Particle	J^P	Overall status
		$\frac{1/2+}{1/2-}$	**** ****	$\frac{\Sigma(1193)}{\Sigma(1385)}$	$\frac{1/2+}{3/2+}$	**** ****
For example:	$egin{array}{l} \Lambda(1520) \ \Lambda(1600) \ \Lambda(1670) \ \Lambda(1690) \end{array}$	3/2 - 1/2 + 1/2 - 3/2 -	**** *** **** ****	$\Sigma(1480) \\ \Sigma(1560) \\ \Sigma(1580) \\ \Sigma(1620)$	3/2 - 1/2	* ** * *
Even low-lying resonances are not well determined.	$\Lambda(1800)$ $\Lambda(1810)$	1/2 - 1/2 +	***	$\begin{array}{c} \Sigma(1660) \\ \Sigma(1670) \end{array}$	$\frac{1/2+}{3/2-}$	*** ****
N* & Δ* spectra are well established in the mass range 0.94 < M < ~ 1.8 GeV.	$ \begin{array}{c} \Lambda(1820) \\ \Lambda(1830) \\ \Lambda(1890) \\ \Lambda(2000) \\ \Lambda(2020) \end{array} $	5/2+ 5/2- 3/2+ 7/2+	**** **** * *		1/2- 1/2+ 5/2- 3/2+	** *** * *
Before 2012, PDG listed only Breit-Wigner (BW) mass and width. (+ "highly" model-dependent !!)	$\Lambda(2100) \ \Lambda(2110) \ \Lambda(2325) \ \Lambda(2350)$	7/2 - 5/2 + 3/2 -	**** *** *	$\Sigma(1880)$ $\Sigma(1915)$ $\Sigma(1940)$ $\Sigma(2000)$ $\Sigma(2020)$	1/2+ 5/2+ 3/2- 1/2- 7/2+	** **** ***
→ N* & Δ* case:	$\Lambda(2585)$		**	$\frac{\Sigma(2030)}{\Sigma(2070)}$	7/2+ 5/2+	**** *
Resonance parameters defined by poles of scattering amplitudes are extensively studied: PDG lists BOTH pole and BW parameter	ers.			$\Sigma(2080) \\ \Sigma(2100) \\ \Sigma(2250) \\ \Sigma(2455)$	3/2+7/2-	** * *** **
### NOTE: Pole parameters should be "model-independent" and meaningful to compare !!	-			$\Sigma(2620) \\ \Sigma(3000) \\ \Sigma(3170)$		** * *

Current situation of $Y^*(= \Lambda^*, \Sigma^*)$ spectroscopy

PDG listing

Comprehensive partial-wave analysis			۸*			Σ*	
of K ⁻ p reactions to extract Y* defined	Λ(13XX)1/2 ⁻ ??	Particle	J^P	Overall status	Particle	J^P	Overall status
by poles has been accomplished <i>just recently</i> :	above	$ \begin{array}{c} \Lambda(1116) \\ \Lambda(1405) \\ \Lambda(1520) \end{array} $	1/2+ 1/2- 3/2-	**** **** ****	$\Sigma(1193) \\ \Sigma(1385) \\ \Sigma(1480)$	$\frac{1/2+}{3/2+}$	**** ****
	KN threshold	$\begin{array}{c} \Lambda(1600) \\ \Lambda(1670) \\ \Lambda(1600) \end{array}$	1/2+ 1/2-	*** ****	$\frac{\Sigma(1560)}{\Sigma(1580)}$	3/2 -	** *
Kent State University (KSU) group	·	$\Lambda(1690)$	3/2-	****	$\Sigma(1620)$ $\Sigma(1660)$	$\frac{1}{2}$	**
(→ 2013, "KSU on-shell parametrization" o Zhang et al., PRC88(2013)035204, 035205.	f S-matrix)	$\Lambda(1800)$ $\Lambda(1810)$ $\Lambda(1820)$ $\Lambda(1830)$ $\Lambda(1800)$	1/2- 1/2+ 5/2+ 5/2- 3/2+	*** *** **** ****	$\Sigma(1660)$ $\Sigma(1670)$ $\Sigma(1690)$ $\Sigma(1750)$ $\Sigma(1770)$	1/2+ 3/2- 1/2- 1/2- 1/2+	*** *** ** *
Our group (→ 2014-2015, dynamical coupled-channels	s approach)	$\Lambda(1390)$ $\Lambda(2000)$ $\Lambda(2020)$ $\Lambda(2100)$	$\frac{3}{2+}$ $\frac{7}{2+}$	**** * *	$\Sigma(1775) \\ \Sigma(1840) \\ \Sigma(1880)$	5/2 - 3/2 + 1/2 +	**** * **
HK, Nakamura, Lee, Sato, PRC90(2014)065204; a	rXiv:1506.01768	$\begin{array}{c} \Lambda(2100) \\ \Lambda(2110) \\ \Lambda(2325) \\ \Lambda(2350) \\ \Lambda(2585) \end{array}$	5/2+ 3/2-	*** * *** **	$ \begin{array}{l} \Sigma(1915) \\ \Sigma(1940) \\ \Sigma(2000) \\ \Sigma(2030) \\ \Sigma(2070) \end{array} $	5/2+ 3/2- 1/2- 7/2+ 5/2+	**** ** * * *
					$ \begin{array}{l} \Sigma(2080) \\ \Sigma(2100) \\ \Sigma(2250) \\ \Sigma(2455) \\ \Sigma(2620) \\ \Sigma(3000) \\ \Sigma(3170) \end{array} $	3/2+7/2-	** * *** ** * *

HK, Nakamura, Lee, Sato, PRC90(2014)065204

channels

Coupled-channels integral equations for partial-wave amplitudes of a \rightarrow b reaction:

$$T_{b,a}^{(LSJ)}(p_b, p_a; E) = V_{b,a}^{(LSJ)}(p_b, p_a; E) + \sum_c \int_0^\infty q^2 dq V_{b,c}^{(LSJ)}(p_b, q; E) G_c(q; E) T_{c,a}^{(LSJ)}(q, p_a; E)$$

$$\frac{\mathsf{CC}}{\mathsf{effect}} \quad \mathsf{off-shell}_{\mathsf{effect}}$$

Reaction channels:

$$a, b, c = (\bar{K}N, \pi\Sigma, \pi\Lambda, \eta\Lambda, K\Xi, \pi\Sigma^*, \bar{K}^*N, \cdots)$$

Transition Potentials:

$$V_{a,b} = v_{a,b} + \sum_{\substack{Y^* \\ Y^* = 0}} \frac{\Gamma_{Y^*,a}^{\dagger} \Gamma_{Y^*,b}}{E - M_{Y^*}}$$

Exchange potentials Bare Y* states

HK, Nakamura, Lee, Sato, PRC90(2014)065204

Coupled-channels integral equations for partial-wave amplitudes of a \rightarrow b reaction:



HK, Nakamura, Lee, Sato, PRC90(2014)065204



HK, Nakamura, Lee, Sato, PRC90(2014)065204

✓ Coupled-channels integral equations for partial-wave amplitudes of $a \rightarrow b$ reaction:



 Momentum integral takes into account off-shell rescattering effects in the intermediate processes.

HK, Nakamura, Lee, Sato, PRC90(2014)065204

✓ Coupled-channels integral equations for partial-wave amplitudes of $a \rightarrow b$ reaction:

Physical Y*s will be a "mixture" of the two pictures:



Why coupled-channels dynamics is important ??

Multichannel unitary condition for $T_{ab}(E) - T^{\dagger}_{ab}(E) = -2\pi i \sum_{c} T^{\dagger}_{ac} \delta(E - E_c) T_{cb}(E)$ T-matrix elements:

1) Ensures "conservation of probabilities" in multichannel reaction processes

- Key to comprehensive & simultaneous analysis of various reactions in a reliable way !!
- 2) Ensures "proper analytic structure" for scattering amplitudes in the complex energy plane (branch points and cuts, etc.)
 - > Without this, it is likely to get false resonance signals.



3) Brings nontrivial features allowing ones to have new physics interpretations for various resonance properties

Dynamical origin of baryon resonances

> Suzuki, Julia-Diaz, HK, Lee, Matsuyama, Sato PRL104 065203 (2010)





What we have done so far

With the dynamical coupled-channels approach developed for the S= -1 sector, we made:

✓ Comprehensive analysis of all available data of K⁻ p → KN, πΣ, πΛ, ηΛ, KΞ up to W = 2.1 GeV.

[HK, Nakamura, Lee, Sato, PRC90(2014)065204]

- Needs supercomputers to accomplish extensive coupled-channels analysis
- > Successfully determined the partial-wave amplitudes of $\overline{K}N \rightarrow \overline{K}N$, $\pi\Sigma$, $\pi\Lambda$, $\eta\Lambda$, $K\Xi$ for S, P, D, and F waves !!

Argonne

 Extraction of Λ* and Σ* mass spectrum defined by poles of scattering amplitudes.

[HK, Nakamura, Sato, arXiv:1506.01768 (to appear in PRC)]

Database of our analysis (W < 2.1GeV)

HK, Nakamura, Lee, Sato, PRC90(2014)065204



Results of the fits



Results of the fits

$K^- p \rightarrow K^- p$ scattering

HK, Nakamura, Lee, Sato, PRC90(2014)065204



dσ/dΩ (1832 < W < 2100 MeV)



Comparison of extracted partial-wave amplitudes



S-wave dominance ??

$K^- p \rightarrow MB$ total cross sections near threshold



Solid: Full Dashed: S wave only

For K- $p \rightarrow \pi\Lambda$, $\eta\Lambda$, K \equiv , higher partial waves visibly contribute to the cross sections even in the threshold region.

→ consistent with the observation in Jackson et al., PRC91(2015)065208

Naïve expectation for S-wave dominance near the threshold sometimes does not hold !!

Extracting Y* resonance parameters

Definitions of



Extracted Λ* and Σ* mass spectrum



Branching ratios

V High-mass Y* have large branching ratio to $\pi\Sigma^*$ ($\pi\pi\Lambda$) & \overline{K}^*N ($\pi\overline{K}N$)

- > K⁻ p $\rightarrow \pi\pi\Lambda$, $\pi\overline{K}N$ data would play a crucial role for establishing high-mass Y^{*}.
 - → Similar to high-mass N* and Δ* case, where $\pi\pi$ N channel plays a crucial role. (e.g., measurement of π N → $\pi\pi$ N reactions at J-PARC E45)



S-wave resonances below KN threshold from the current analysis

HK, Nakamura, Lee, Sato, arXiv:1506.01768

NOTE: Further extensive analysis including the data below KN threshold is necessary to have *conclusive* results for the KN subthreshold region.





Summary

- ✓ Comprehensive analysis of K⁻ p → KN, πΣ, πΛ, ηΛ, KΞ up to W = 2.1 GeV has been accomplished for the first time within a dynamical coupled-channels approach.
- Partial-wave (S, P, D, and F) amplitudes &
 Λ* and Σ* resonance parameters (pole mass, width,...)
 have been successfully extracted.

Visible analysis dependence exits in extracted resonance parameters



Lack of the K- p reaction data for

polarization observables (P and β , R, A) the $\overline{K}N$ threshold region 3-body ($\pi\pi\Lambda$, $\pi\overline{K}N$, ...) production reaction

J-PARC is a unique facility to overcome this unsatisfactory situation !!