

# **Workshop on Future Prospects of Hadron Physics at J-PARC and Large Scale Computational Physics**

February 9 – 11, 2012

Ibaraki Quantum Beam Research Center

## **Production of medium and heavy hypernuclei**

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(Osaka E-C)

# CONTENTS

## 1. Introduction / Basic motivations

production and decays of hypernuclei is  
a window to disclose B-B interactions

## 2. Reaction spectroscopy

some theoretical predictions beyond p-shell  
(e,e'K+) at JLab, (K-, $\pi$ ) and ( $\pi$ ,K) at JPARC

## 3. Production of strangeness -2 hypernuclei

suggest (K-,K+) reaction on odd-Z sd-shell  
targets

## 4. Summary

$${}^6\text{Li} = 3p + 3n$$

$${}^7_{\Lambda}\text{Li} = 3p + 3n + \Lambda$$

$${}^{11}\text{C} = 6p + 5n$$

$${}^{12}_{\Lambda}\text{C} = 6p + 5n + \Lambda$$

$${}^{12}_{\Xi}\text{B} = 6p + 5n + \Xi^-$$

M-B-I, PTP Suppl. 81 (1986)  
B-M-Z, Int.J. Mod.Phys.A5 (1990)

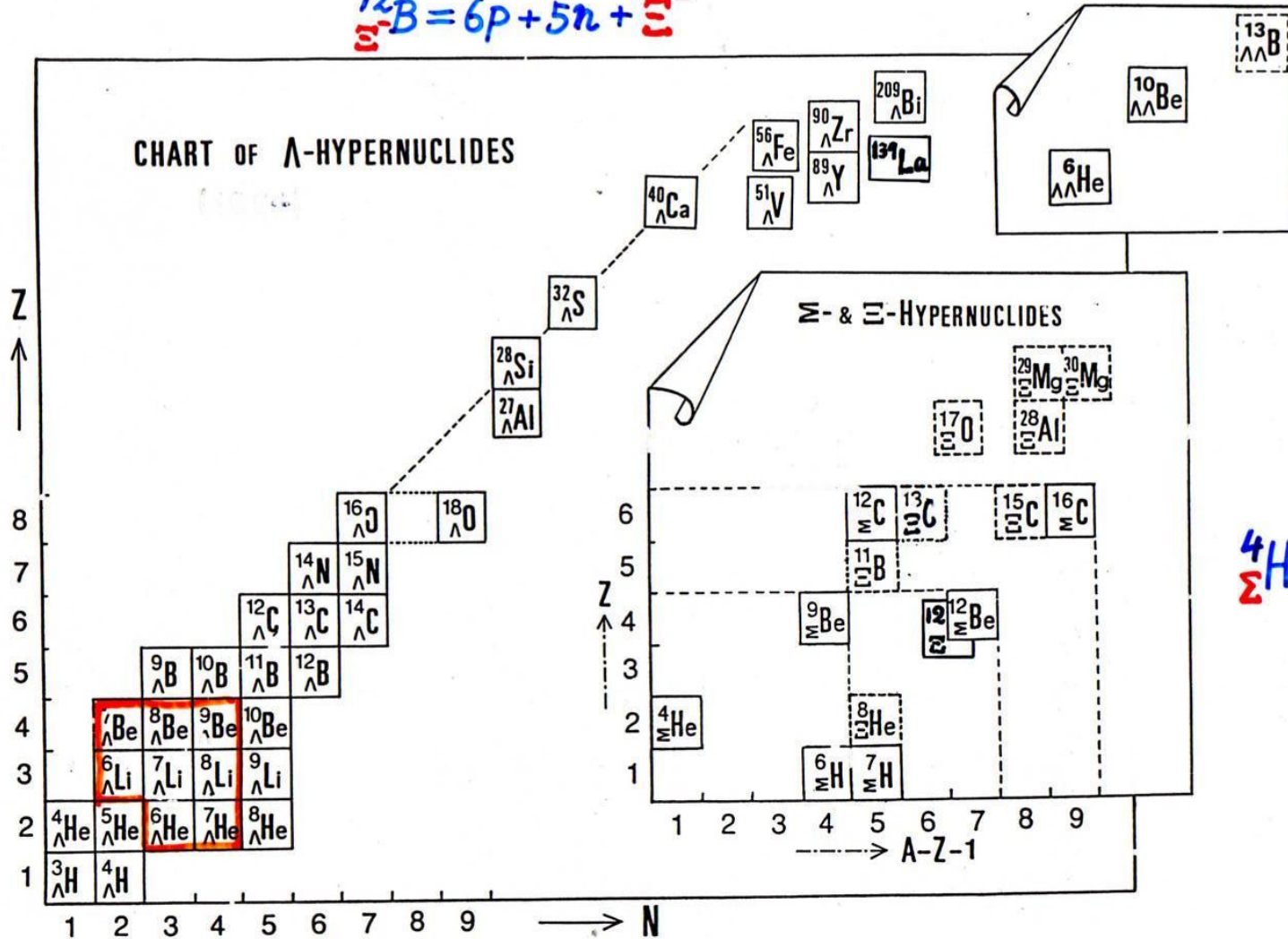


Fig. 1-1

$N_u \sim N_d \sim N_s$

(H. Tamura)

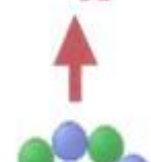
Strangeness in neutron stars ( $\rho > 3 - 4 \rho_0$ )

Strange hadronic matter ( $A \rightarrow \infty$ )

$S = -\infty$

$p, n, \Lambda, \Xi^0, \Xi^-$

↑ higher density



Strangeness

$S = -2$

$S = -1$

$\Lambda\Lambda, \Xi$  hypernuclei

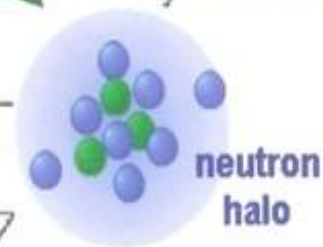
$\Lambda, \Sigma$  hypernuclei

$\Lambda N$  interaction

Proton-rich nuclei

Neutron-rich nuclei

non-strange nuclei



neutron number

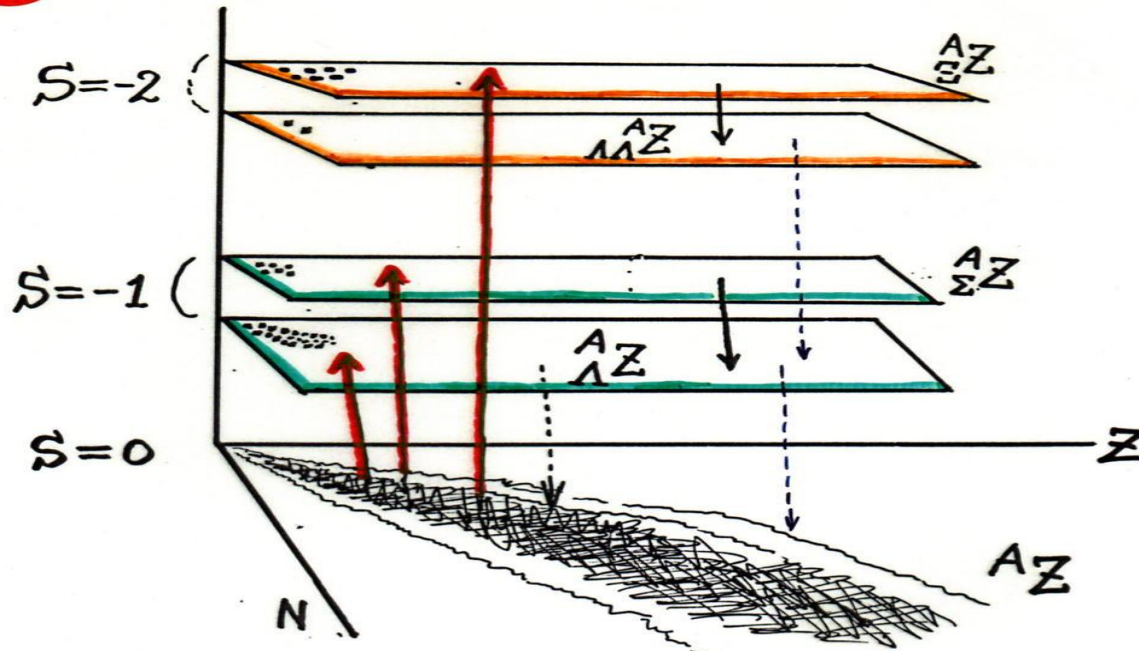
3-dimensional nuclear chart

# Old slide for many-body system with **S**

● “Strangeness を持った多体系の物理”

## First 3D-Chart

T.Motoba, Talk at  
**Program on Big  
Hadron Project  
Physics,**  
INS, U. Tokyo  
(1986),  
Genshikaku  
Kenkyu **32**, No.2.  
97 (1987)



**Intense  $K^-$  beam** + **Good  $K^+$  detector** + 人

1. Pauli-free な  $Y$  を probe して 核多体系の response
2. Hy 核 ----- new symmetry  
genuinely hypernuclear states
3. Interaction -----  $NN$ ,  $YN$ ,  $YY$
4. 多体系の論理  
{ elementary  $\leftrightarrow$  q-based physics

# 1. Introduction

## *Why medium-heavy hypernuclei ?*

- 1)  $\Lambda$  single-particle energies: not well known
- 2) More chances of high-spin selectivity to see new aspects such as hyperon coupled with rotational and/or vibrational states.
- 3) Recent (e,e'K+) experiments encourage to go to sd- and fp-shell regions.
- 4) Expect to find “stable”  $\Xi$ -hypernuclear states, which leads to hyperon-mixing phenomena

# Single-particle energies of $\Lambda$

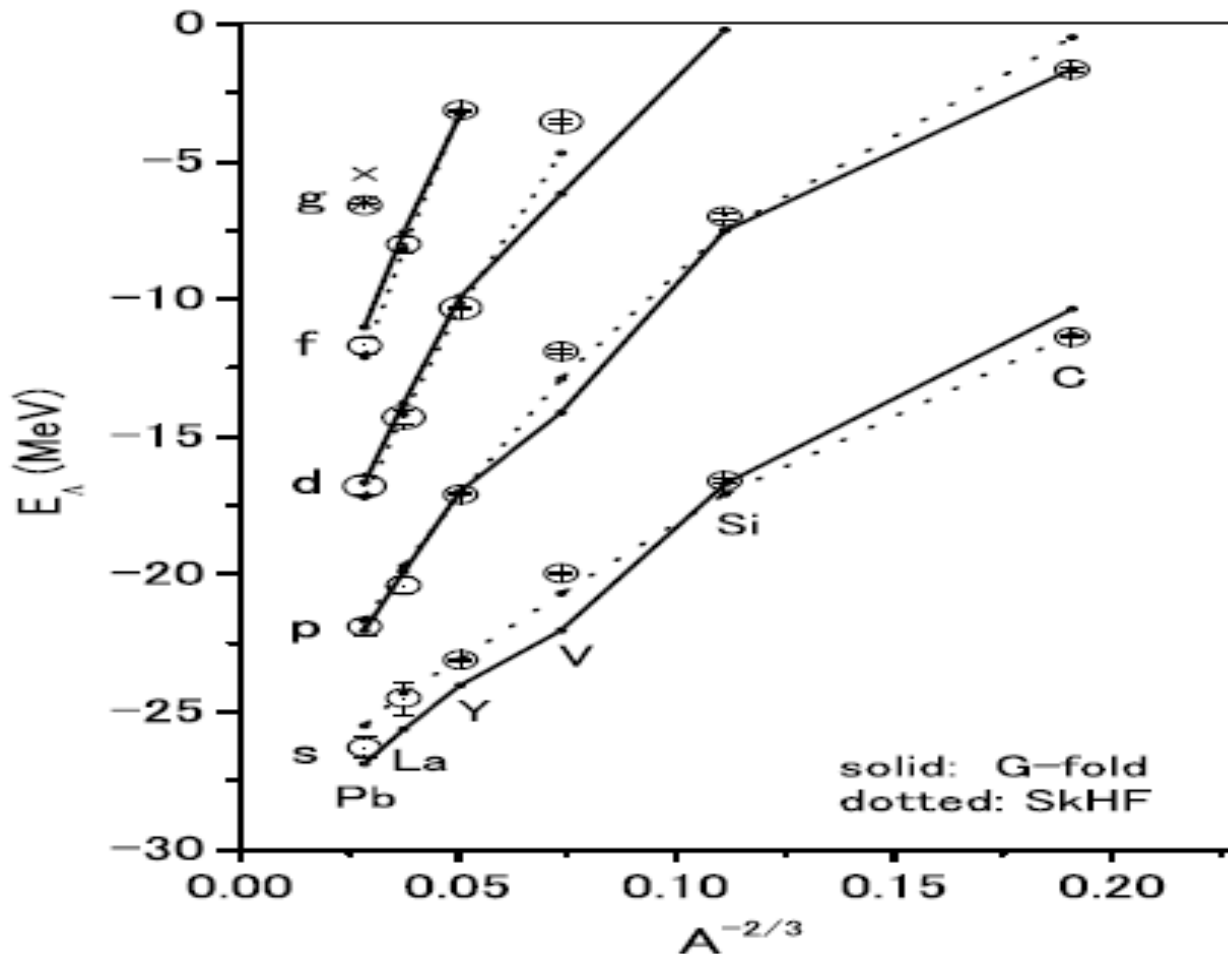


Fig. 1. Energy spectra of  $^{13}_{\Lambda}\text{C}$ ,  $^{28}_{\Lambda}\text{Si}$ ,  $^{51}_{\Lambda}\text{V}$ ,  $^{89}_{\Lambda}\text{Y}$ ,  $^{139}_{\Lambda}\text{La}$  and  $^{208}_{\Lambda}\text{Pb}$  are given as a function of  $A^{-2/3}$ ,  $A$  being mass numbers of core nuclei. Solid (dashed) lines show calculated values by the G-matrix folding model derived from ESC08a (the Skyrme-HF model). Open circles denote the experimental values taken from Ref. 17).

## 2.1 Developments in reaction spectroscopy of hypernuclear production

$${}^A_Z(J_i T_i \tau_i)(K^-, \pi^-) {}^A_Z(J_f T_f \tau_f) \quad \begin{array}{l} (\pi^+, K^+) \\ (\gamma, K^+) \\ (K^-, K^+) \end{array}$$

**Theoretically reliable analyses take account of:**

1. DW effects (DW vs. PW),
2. Microscopic treatment with elem. amplitudes,
3. Nuclear core excitation effects

$(K^-, \pi^-)$

$(\pi^+, K^+)$

played a great role of  
exciting high-spin series

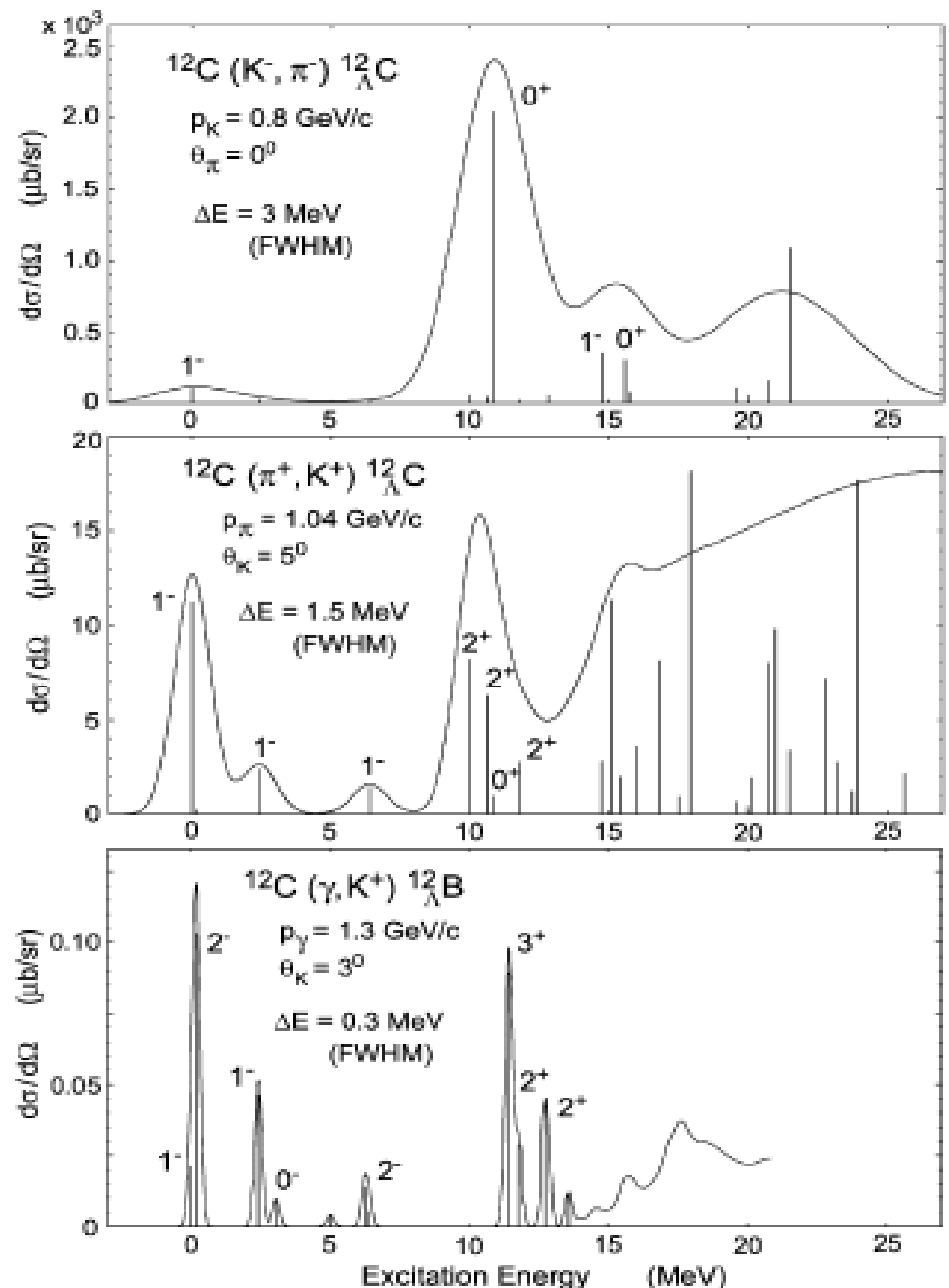
$\Gamma = 1.5 \text{ MeV (best)}$

$(e, e'K^+), (\gamma, K^+)$

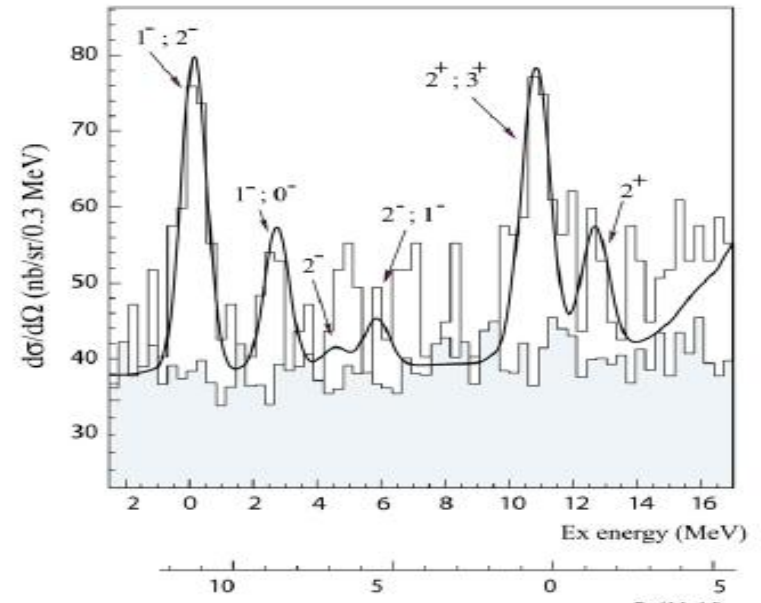
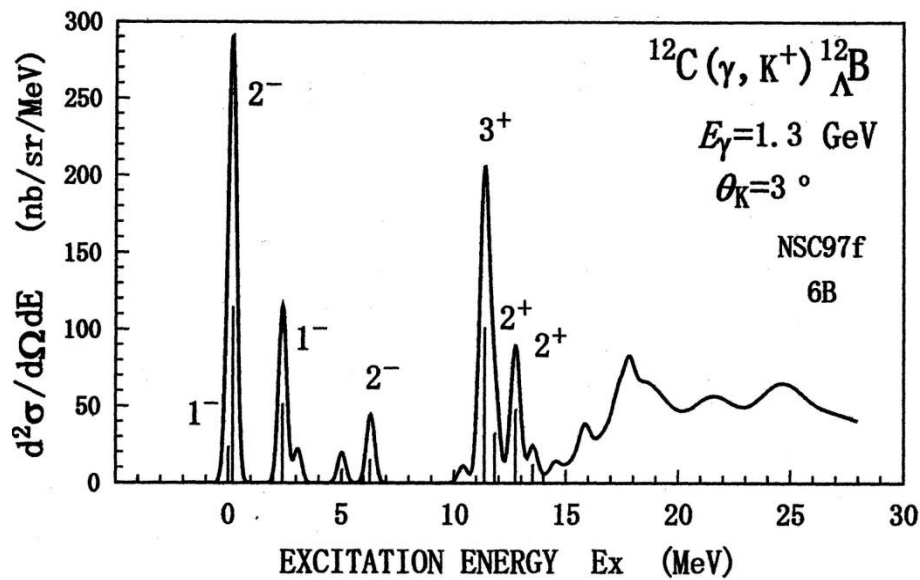
Motoba, Sotona, Itonaga,  
Prog.Theor.Phys.S.117(1994)

T.M. Mesons & Light Nuclei  
(2000) updated w/NSC97f.

-----  
JLab Exp't :  $\Gamma = 0.5 \text{ MeV}$



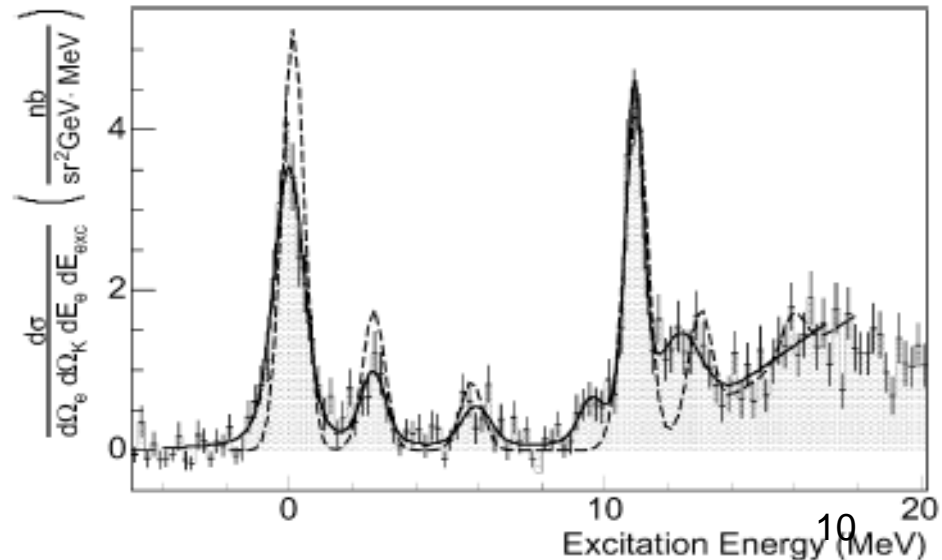
# Theor. prediction vs. (e,e'K<sup>+</sup>) experiments



Motoba, Sotona, Itonaga,  
*Prog.Theor.Phys.Sup.***117** (1994)  
 T.M. *Mesons & Light Nuclei* (2000)  
 updated w/NSC97f.

----- Sotona's Calc.----->

Hall C (up) T. Miyoshi et al.  
*P.R.L.***90** (2003) 232502.  $\Gamma=0.75 \text{ MeV}$   
 Hall A (bottom), J.J. LeRose et al.  
*N.P. A***804** (2008) 116.  $\Gamma=0.67 \text{ MeV}$



# (e,e'K+) cross sections confirmed.

Table II. Comparison of the experimental energy levels and cross sections for  $^{12}\text{C}(e, e'K^+)^{12}_A\text{B}$  with the theoretical estimates. The systematic experimental errors are not shown. Taken from Ref. 3).

$E_x^{\text{Exp}}$ (MeV)	Width (MeV)	Cross section (nb/sr <sup>2</sup> /GeV)	$E_x^{\text{Cal}}$ (MeV)	Main structure $^{11}\text{B}[J_c] \otimes j^A$	$J_f^\pi$	Cross section (nb/sr <sup>2</sup> /GeV)
0.0±0.03	1.15±0.18	4.48±0.29	0.0	$[3/2^-; \text{g.s.}] \otimes s_{1/2}^A$	1 <sup>-</sup>	1.02
			0.14	$[3/2^-; \text{g.s.}] \otimes s_{1/2}^A$	2 <sup>-</sup>	3.66
2.65±0.10	0.95±0.43	0.75±0.16	2.67	$[1/2^-; 2.12] \otimes s_{1/2}^A$	1 <sup>-</sup>	1.54
5.92±0.13	1.13±0.29	0.45±0.13	5.74	$[3/2^-; 5.02] \otimes s_{1/2}^A$	2 <sup>-</sup>	0.58
			5.85	$[3/2^-; 5.02] \otimes s_{1/2}^A$	1 <sup>-</sup>	0.18
9.54±0.16	0.93±0.46	0.63±0.20	—	—	—	—
10.93±0.03	0.67±0.15	3.42±0.50	10.48	$[3/2^-; \text{g.s.}] \otimes p_{3/2}^A$	2 <sup>+</sup>	0.24
			10.52	$[3/2^-; \text{g.s.}] \otimes p^A$	1 <sup>+</sup>	0.12
			10.98	$[3/2^-; \text{g.s.}] \otimes p_{1/2}^A$	2 <sup>+</sup>	1.43
			11.05	$[3/2^-; \text{g.s.}] \otimes p_{3/2}^A$	3 <sup>+</sup>	2.19
			12.95	$[1/2^-; 2.12] \otimes p_{3/2}^A$	2 <sup>+</sup>	0.91
12.36±0.13	1.58±0.29	1.19±0.36	13.05	$[1/2^-; 2.12] \otimes p^A$	1 <sup>+</sup>	0.27

$(\pi^+, K^+)$

PW vs. DW

(1)

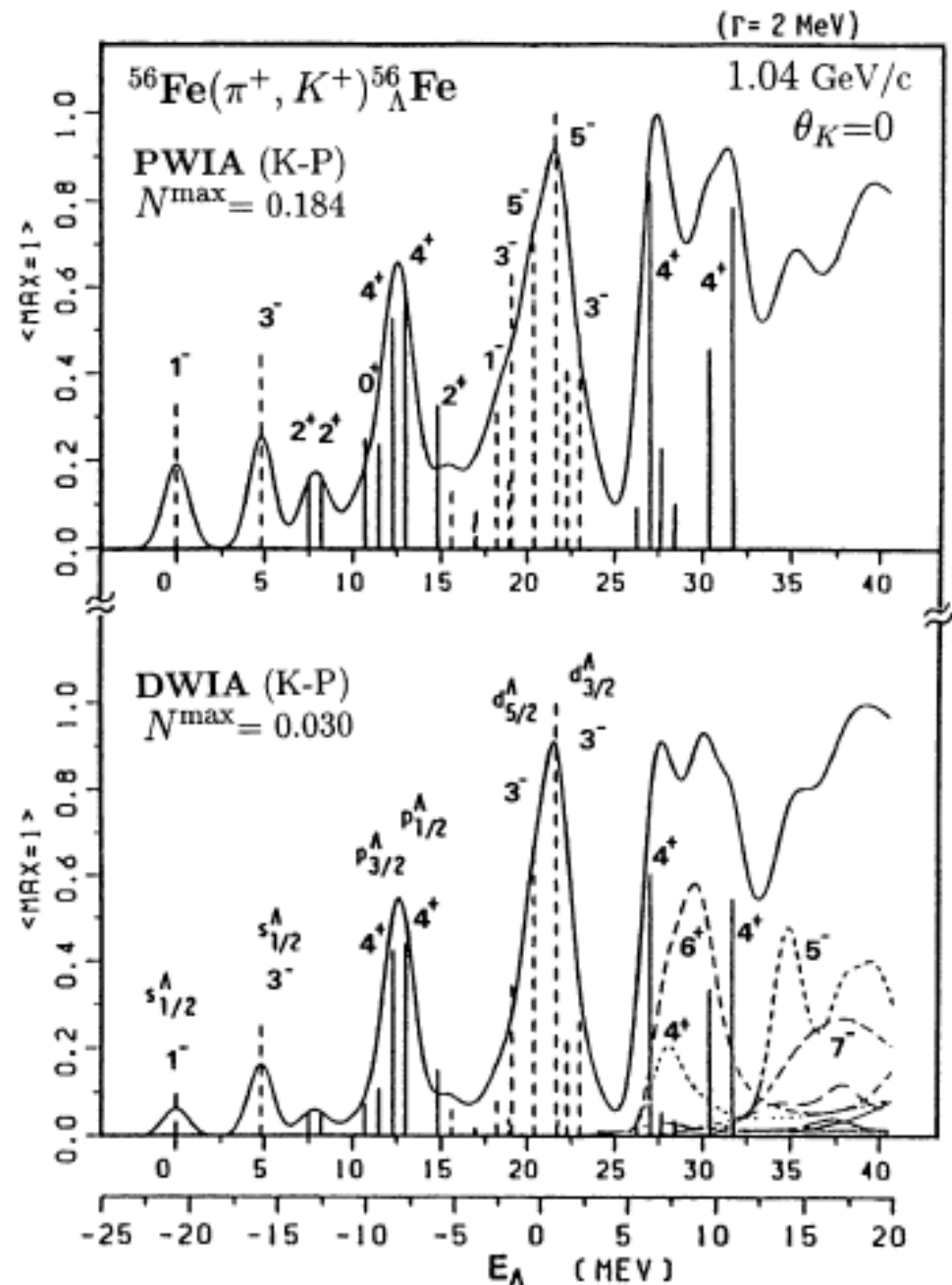
In a typical  $(\pi^+, K^+)$ :

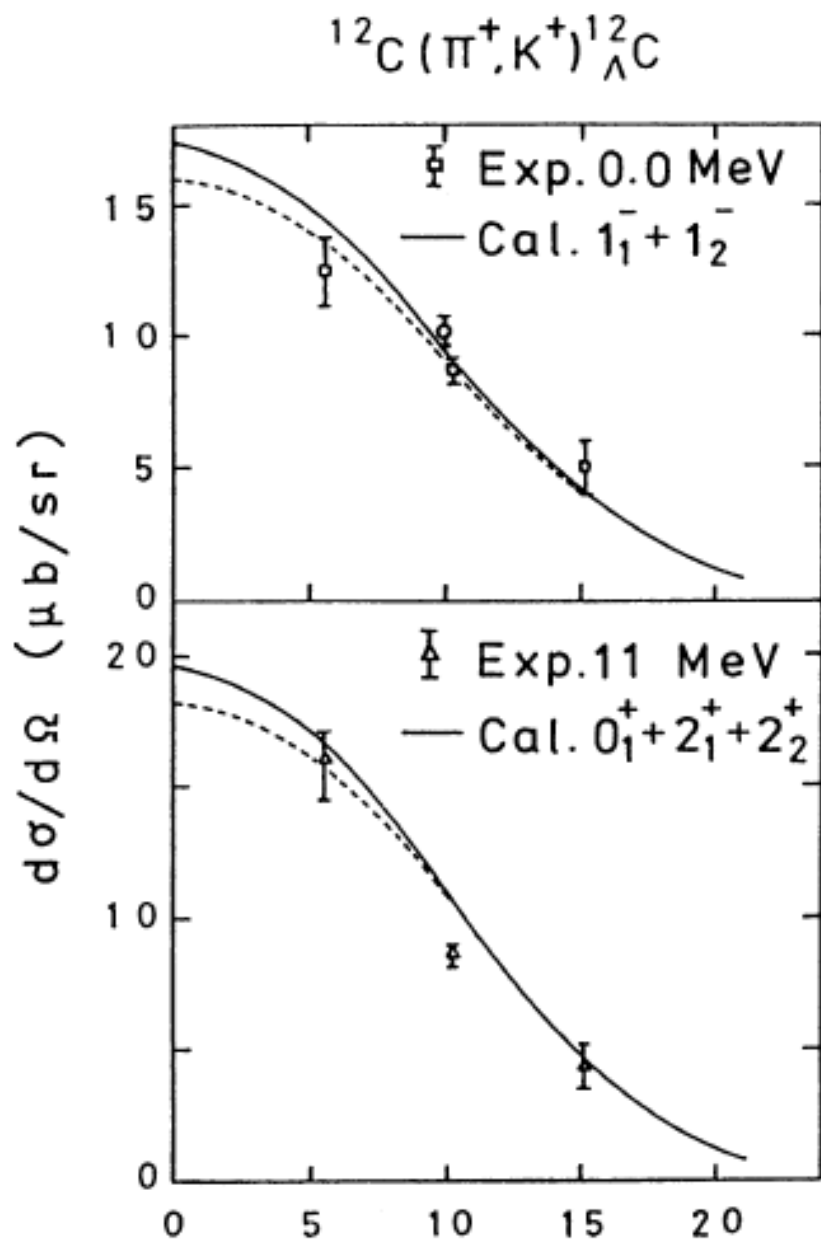
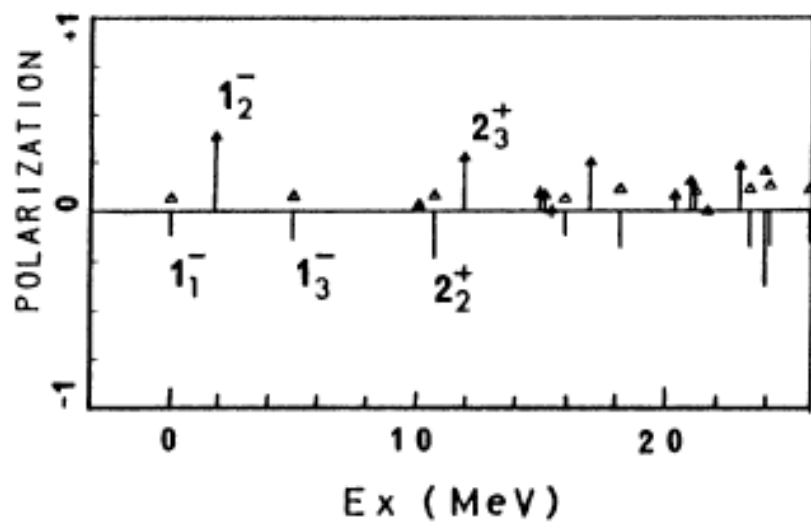
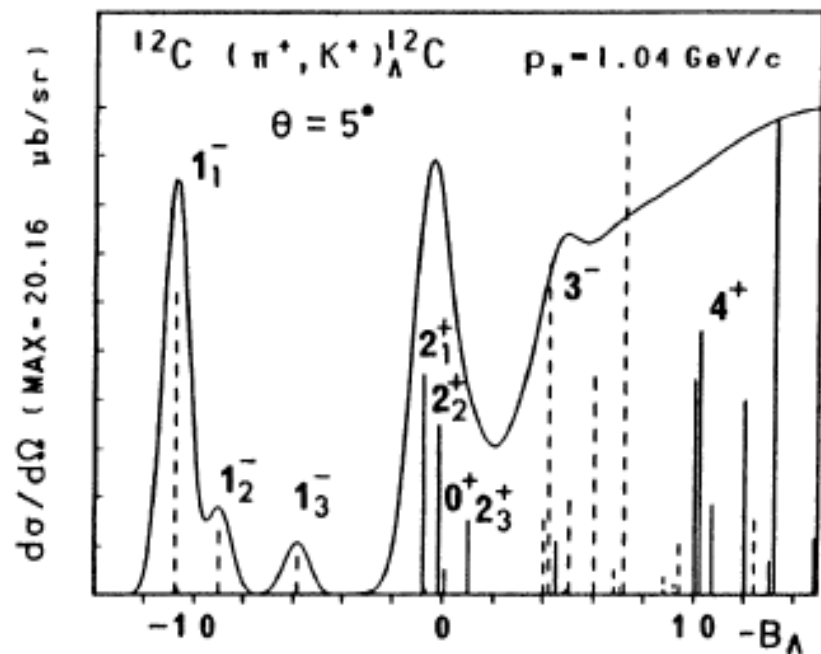
$N_{\text{eff}} = 0.184$  (PW)

$0.030$  (DW)  $1/6$

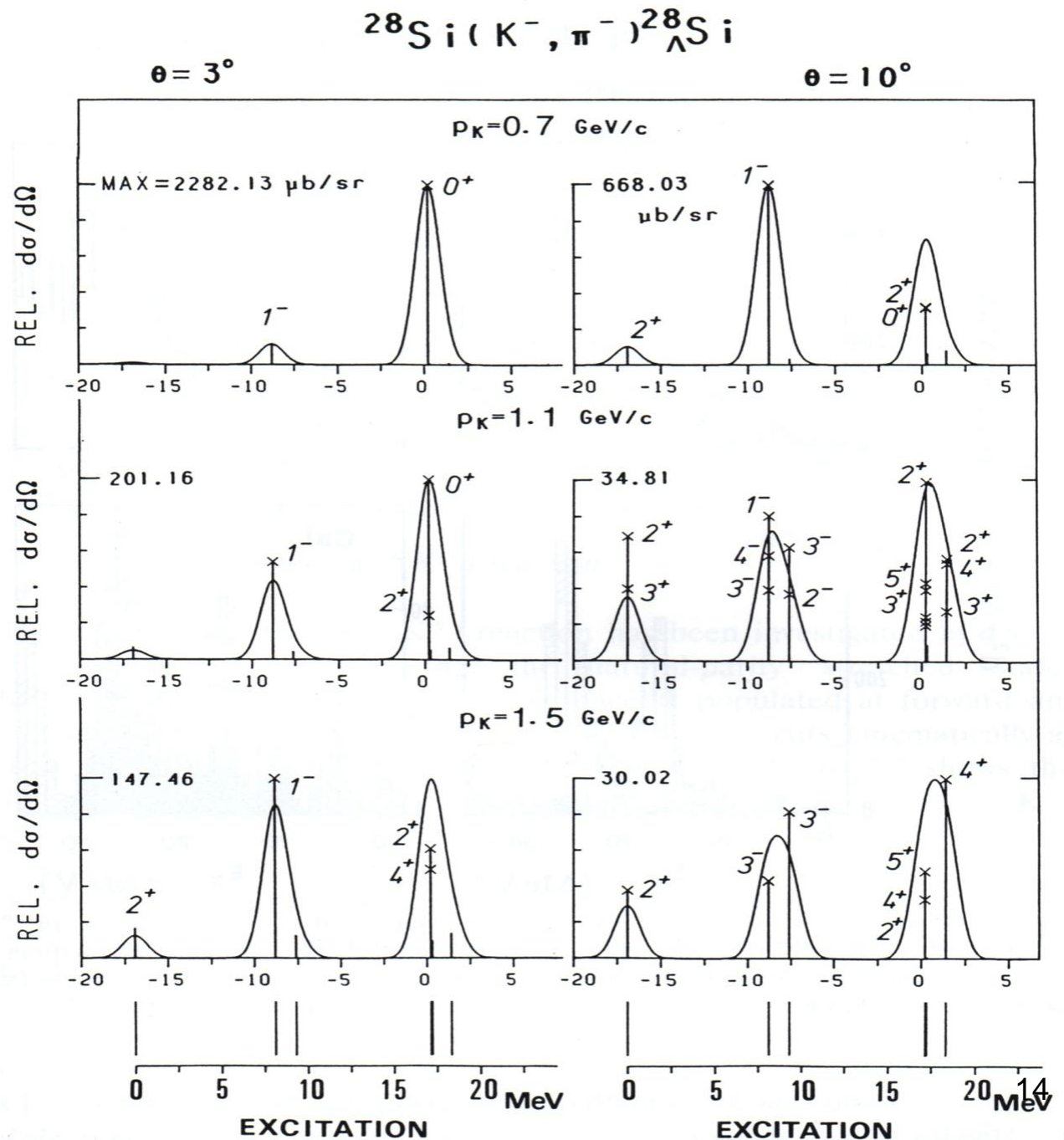
(2) XS to low-J states are much more reduced, resulting in the sharper peaks

(3) Interesting DW effect (mechanism)



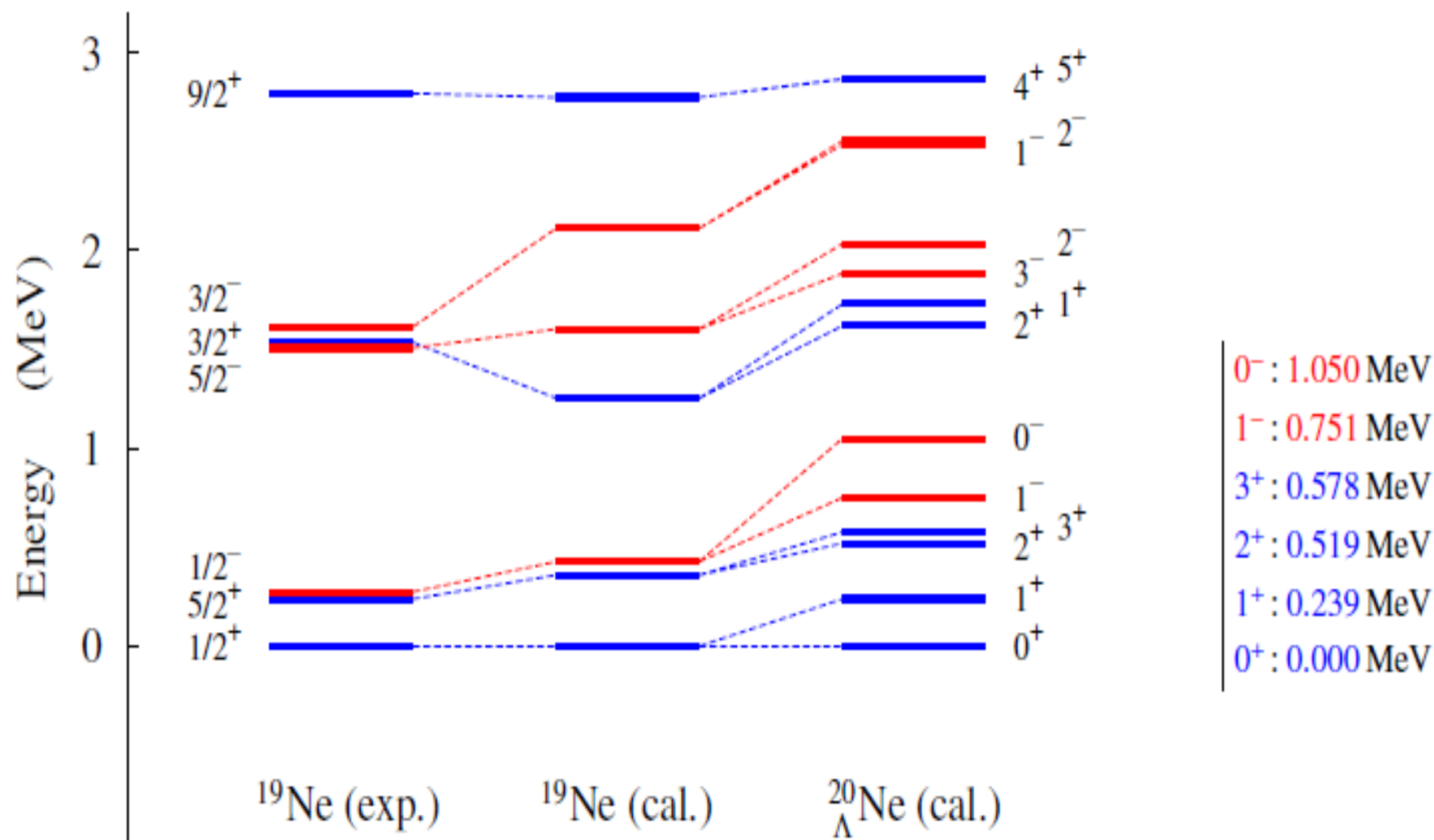


$(K^-, \pi^-)$   
 on  $^{28}\text{Si}$   
 at  
 $p=0.7$   
 $p=1.1$   
 $p=1.5$   
 $\text{GeV}/c$



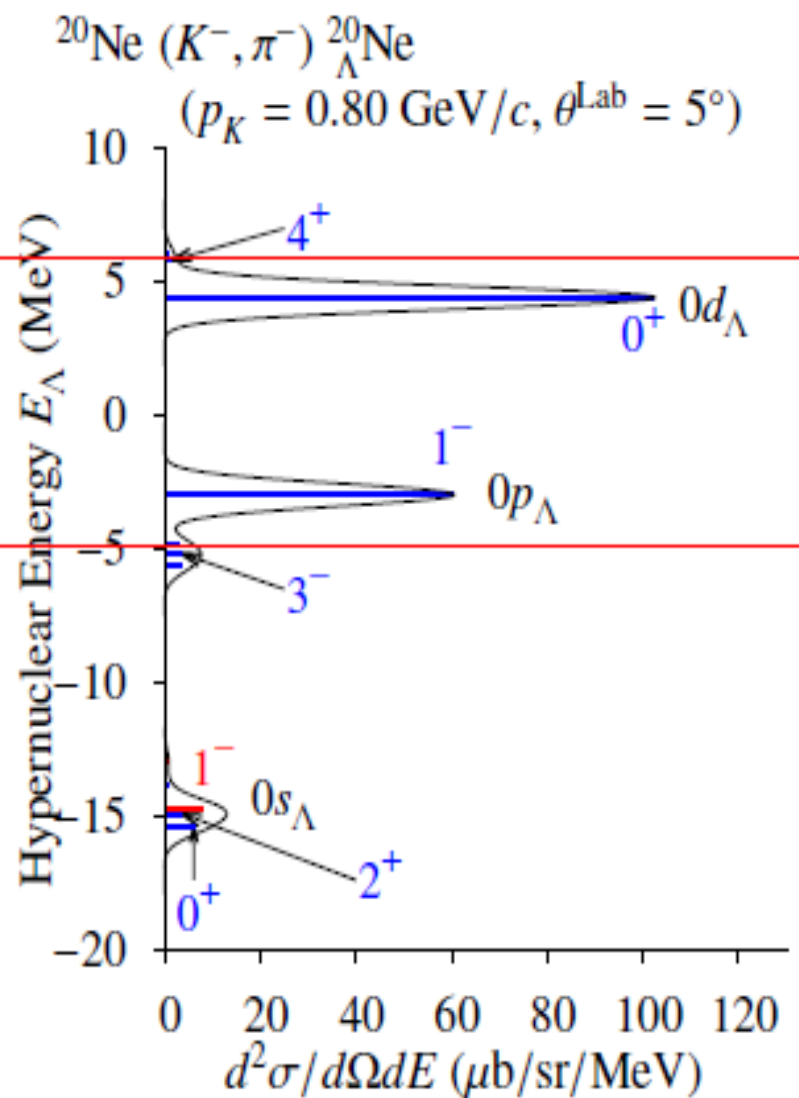
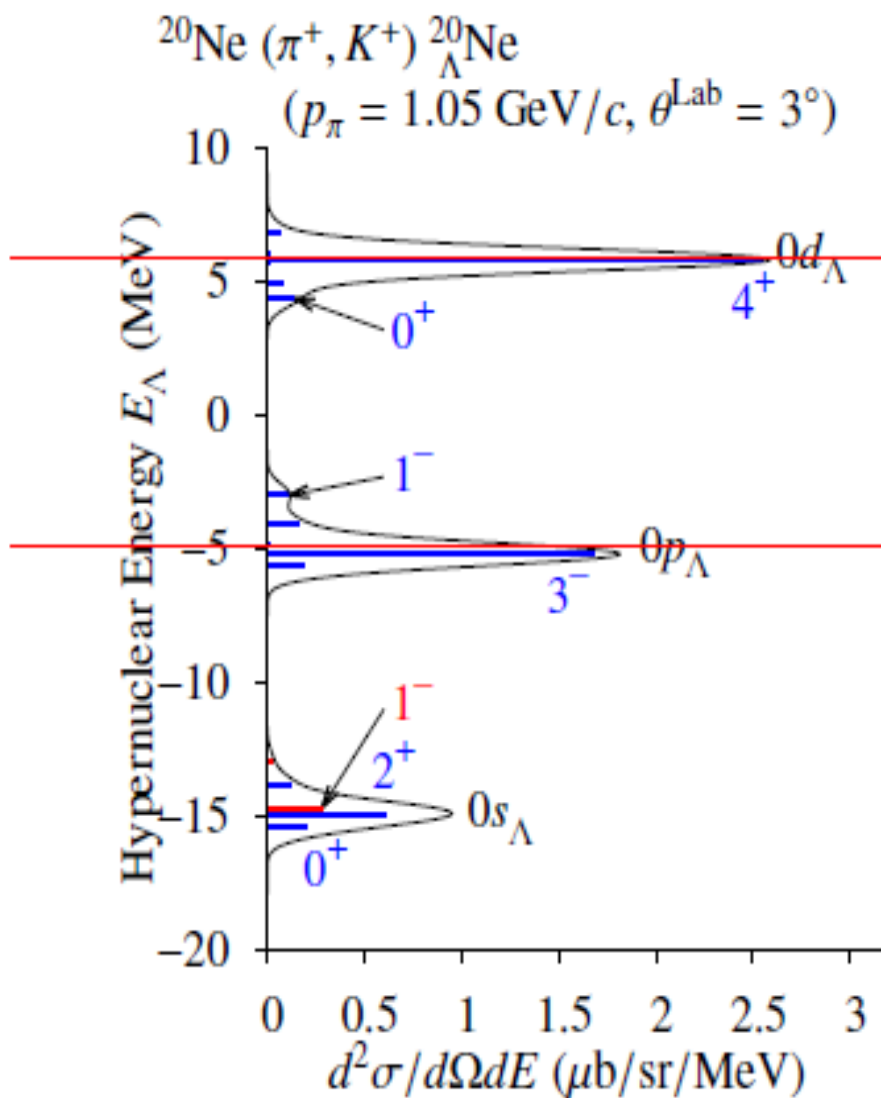
# Numerical Results : Low-lying energy levels of $^{20}_{\Lambda}\text{Ne}$ and $^{19}\text{Ne}$

(by Umeya)



# Numerical Results : Production cross sections

(by Umeya)



# **(K-,p) at $p \sim 1.1$ and $1.5$ GeV/c expected to be done at J-PARC**

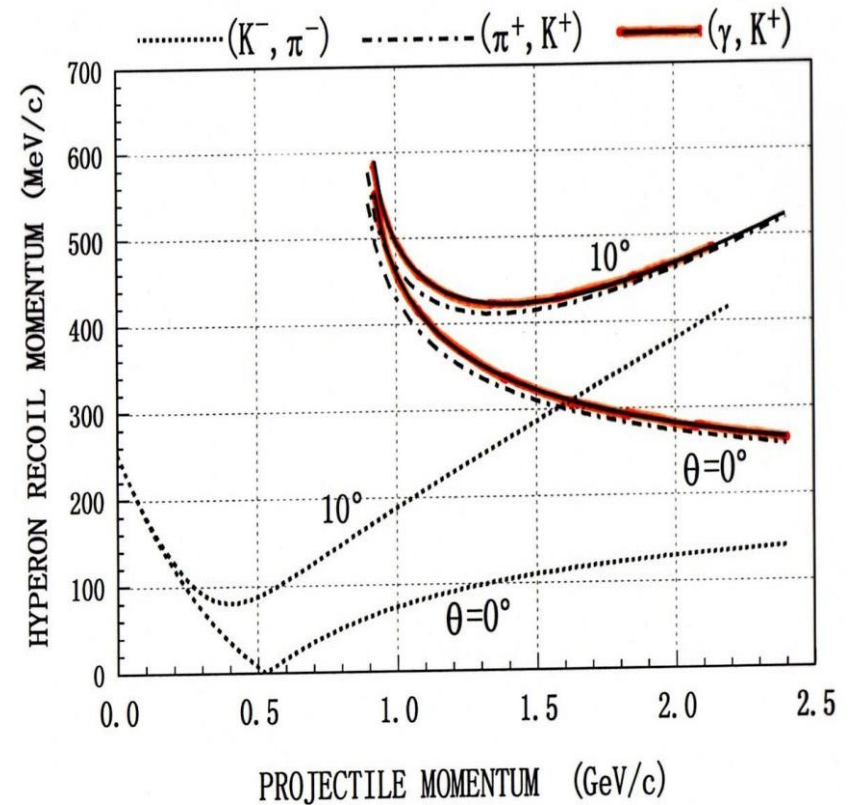
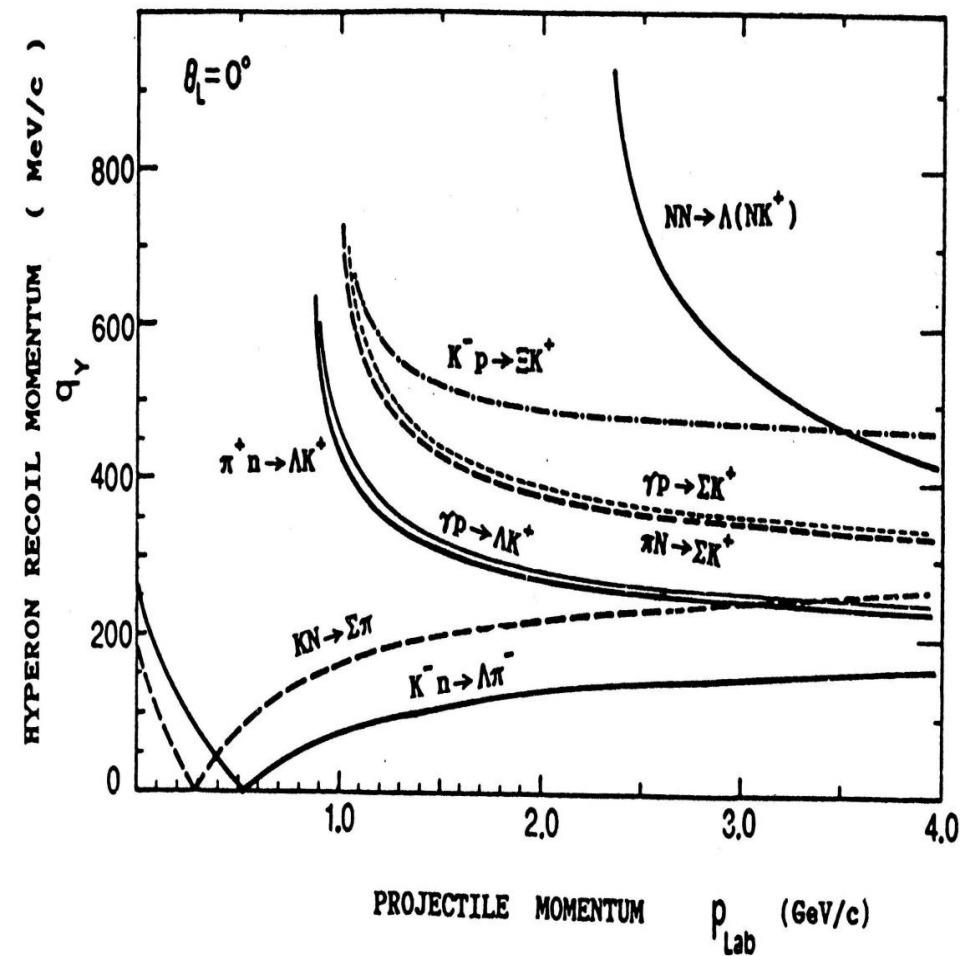
- can produce polarized hypernuclei
- Not only “substitutional ( $\Delta J=0$ )” states but also other states are excited at the same time.
- provide chances to see interplays between  $\Lambda$  and collective motions of nuclear core

## 2.2 Electro/photo-production of sd-shell ~ hypernuclei

(some cases for sd-shell: already done at Jlab, and expected to be done at MAMI)

- Microscopic based on elem. ampl.
- DW: solution of the Klein-Gordon eq.
- Emphasize the importance of taking account of nuclear core excitation effects

# Hyperon recoil momentum and the transition operator determine the reaction characteristics



**$q_\Lambda = 350-420$  MeV/c at  $E_\gamma = 1.3$  GeV**

# Lab $d\sigma/d\Omega$ for photoproduction (2Lab)

$$\left. \frac{d\sigma}{d\Omega} \right|_{2\text{Lab}} = \frac{(2\pi)^4 p^2 E_K E_\gamma E_A}{k \{ p(E_A + E_K) - k E_K \cos \theta_L \}} \left| \langle \mathbf{k} - \mathbf{p}, \mathbf{p} | t | \mathbf{k}, 0 \rangle_L \right|^2, \quad (2.4)$$

$$\langle \mathbf{k} - \mathbf{p}, \mathbf{p} | t | \mathbf{k}, 0 \rangle_L = a_1(\boldsymbol{\sigma} \cdot \boldsymbol{\epsilon}) + a_2(\boldsymbol{\sigma} \cdot \hat{\mathbf{k}})(\hat{\mathbf{p}} \cdot \boldsymbol{\epsilon}) + a_3(\boldsymbol{\sigma} \cdot \hat{\mathbf{p}})(\hat{\mathbf{p}} \cdot \boldsymbol{\epsilon}) + a_4\{(\hat{\mathbf{k}} \times \hat{\mathbf{p}}) \cdot \boldsymbol{\epsilon}\}. \quad (2.5)$$

$$\langle \mathbf{k} - \mathbf{p}, \mathbf{p} | t | \mathbf{k}, 0 \rangle_L = \epsilon_0 \underbrace{(f_0 + g_0 \sigma_0)}_{\text{Spin non-flip term}} + \epsilon_x \underbrace{(g_1 \sigma_1 + g_{-1} \sigma_{-1})}_{\text{Spin-flip terms}} \quad (2.11)$$

with definitions of the coefficients:

$$f_0 = a_4 \sin \theta_L,$$

$$g_0 = a_1,$$

$$g_{\pm 1} = \frac{1}{\sqrt{2}} \{ \mp (a_1 + a_3 \sin^2 \theta_L) - i \sin \theta_L (a_2 + a_3 \cos \theta_L) \}. \quad (2.12)$$

Spin-flip interaction are dominant

# Theor. x-section for $(d_{5/2})^6 (\gamma, K^+) [j_h - j_\Lambda] J$

DWIA		[ nb/sr ]											
Lambda=		s1/2L		p3/2L		p1/2L		1s1/2L		d5/2L		d3/2L	
		(-16.92)		(-8.40)		(-8.40)		(0.32)		(0.69)		(0.69)	
Proton hole d5/2 (-16.17)				1-	5.4					0+	0.0		
				2-	7.1					1+	26.0	1+	8.9
	2+	29.2		3-	4.2	2-	19.4	2+	2.2	2+	0.3	2+	34.9
	3+	63.8		4-	141.8	3-	76.2	3+	4.6	3+	26.7	3+	30.4
	(g.s.)									4+	0.5	4+	112.0
										5+	164.1		
p1/2 (-25.49)	0-	9.4				0+	0.0	0-	3.7				
	1-	30.5		1+	2.0	1+	28.3	1-	12.2			1-	1.4
				2+	66.9					2-	10.7	2-	43.5
										3-	76.9		
p3/2 (-29.84)				0+	0.0							0-	2.0
	1-	14.3		1+	8.9	1+	1.8	1-	5.9	1-	3.2	1+	5.7
	2-	59.1		2+	0.4	2+	62.5	2-	24.8	2-	4.5	2+	17.5
				3+	109.1					3-	4.5	3+	96.3
										4-	148.6		
s1/2 (-44.55)	0+	0.1				0-	7.3	0+	0.3				
	1+	19.2		1-	12.1	1-	23.7	1+	51.4			1+	16.5
				2-	50.0					2+	27.0	2+	40.1
										3+	58.1		

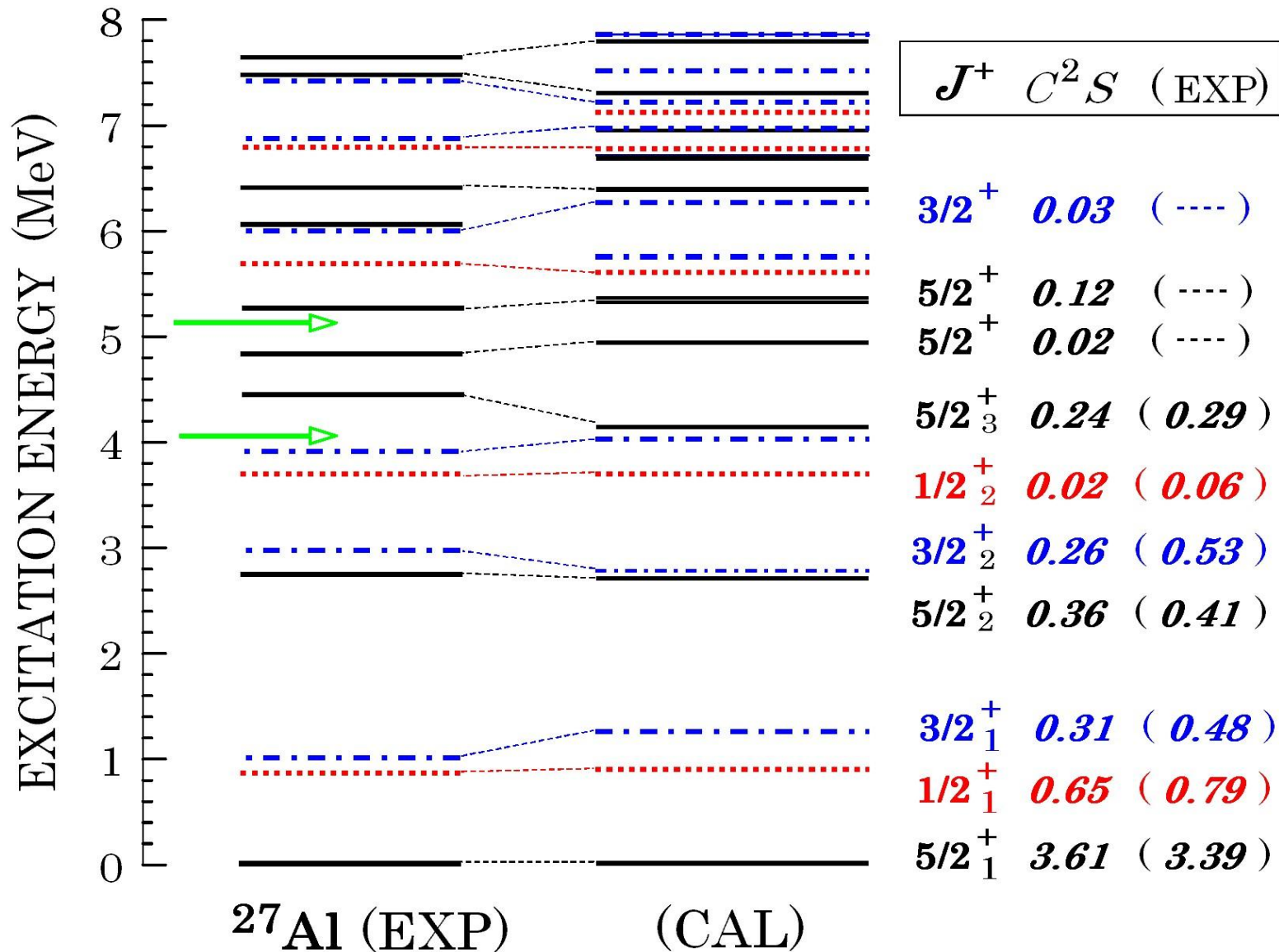
# Realistic prediction for $^{28}\text{Si} (\gamma, K^+)_{\Lambda} ^{28}\text{Al}$ with full sd-wf

By fully taking account of

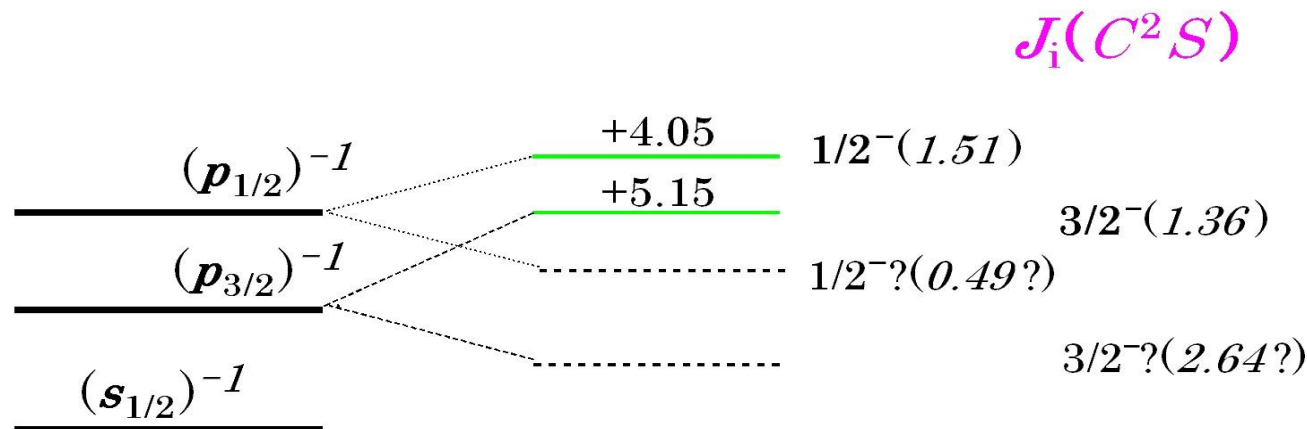
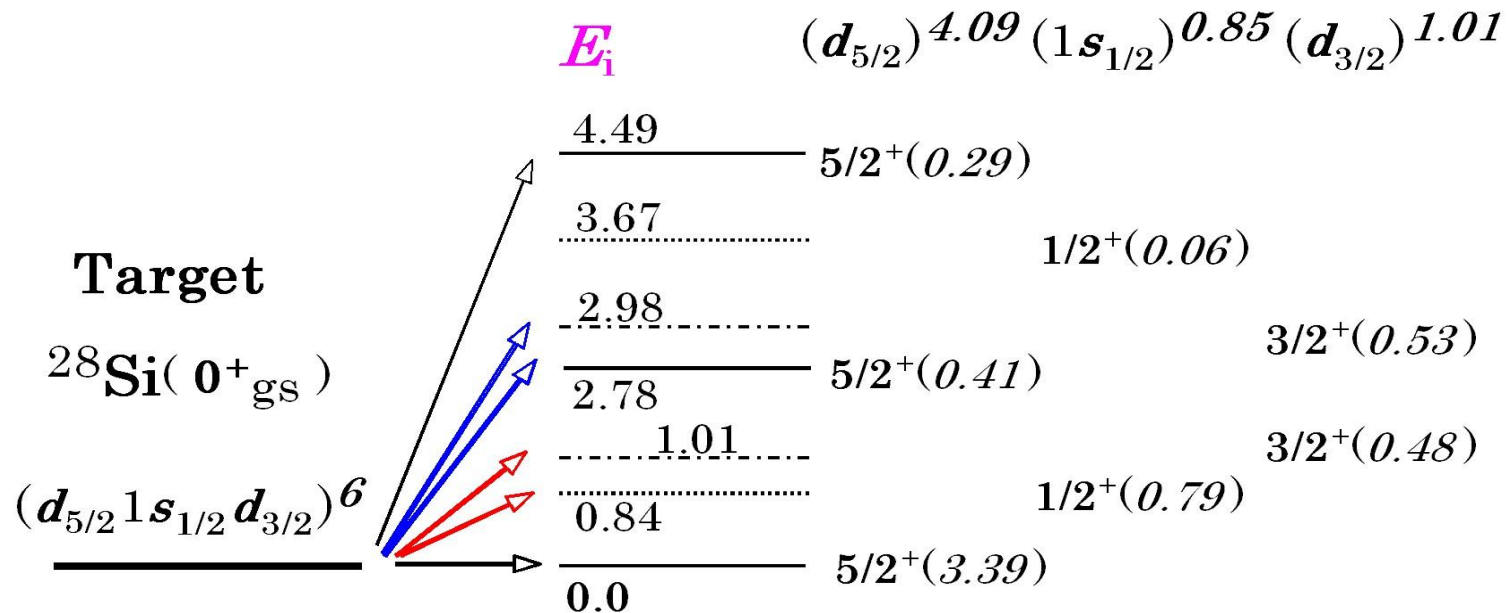
- full  $p(sd)^6.n(sd)^6$  configurations,
- **fragmentations** when a proton is converted,
- $^{27}\text{Al}$  core **nuclear excitation**
- $K^+$  wave distortion effects

→ **Comparison with the  $^{28}\text{Si} (e,e'K^+)$  exp.**

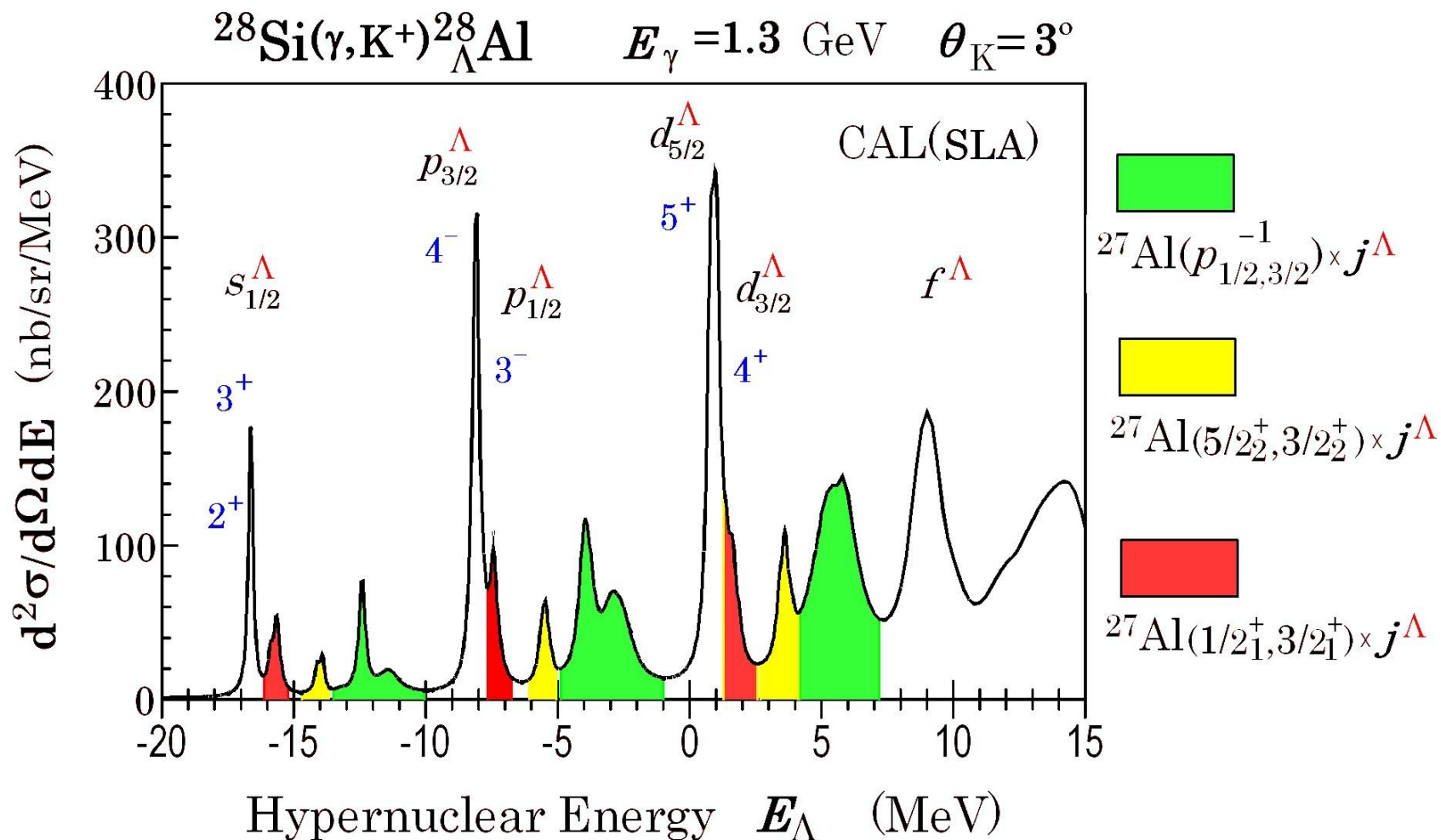
Proton pickup from  $^{28}\text{Si}(0^+): (sd)^6 = (d_{5/2})^{4.1} (1s_{1/2})^{0.9} (d_{3/2})^{1.0}$



proton-state fragmentations should  
be taken into account *to be realistic*

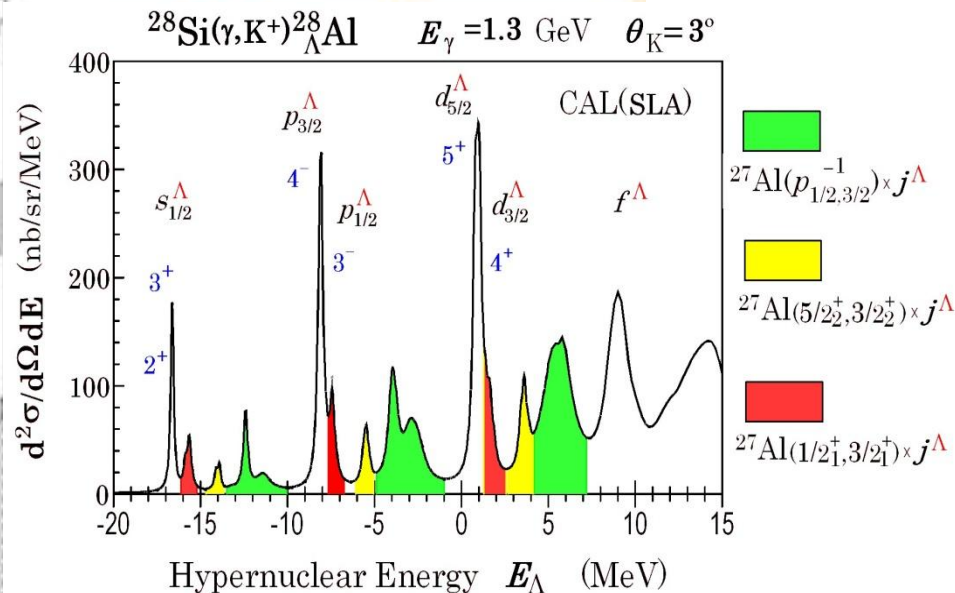
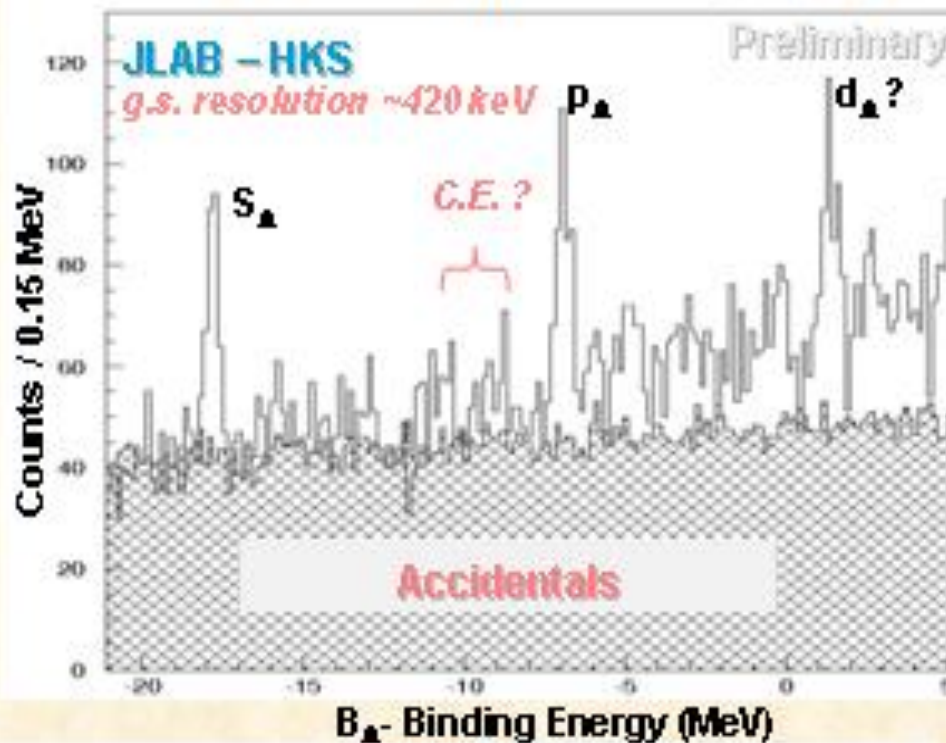


# Peaks can be classified by the characters

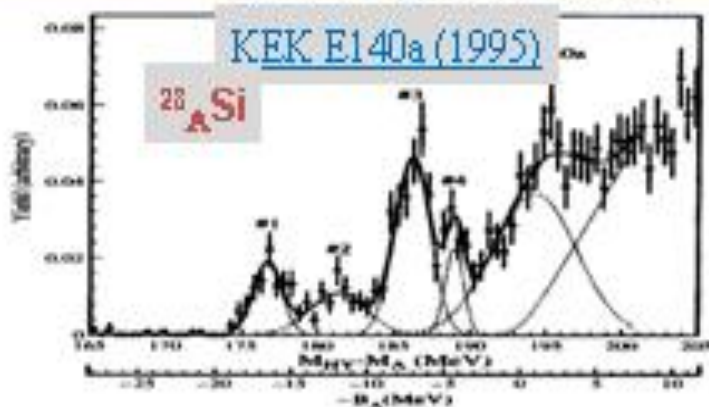


*Major peak series* :  $[^{27}\text{Al}(5/2_1^+) \times j^{\Lambda}]_J$  with  $j^{\Lambda} = s, p, d, \dots$

# $^{28}\text{Si}(e,e'K^+)^{28}_{\Lambda}\text{Al}$ – First Spectroscopy of $^{28}_{\Lambda}\text{Al}$



Major peak series :  $[^{27}\text{Al}(5/2_1^+) \times j^{\Lambda}]_J$  with  $j^{\Lambda} = s, p, d, \dots$

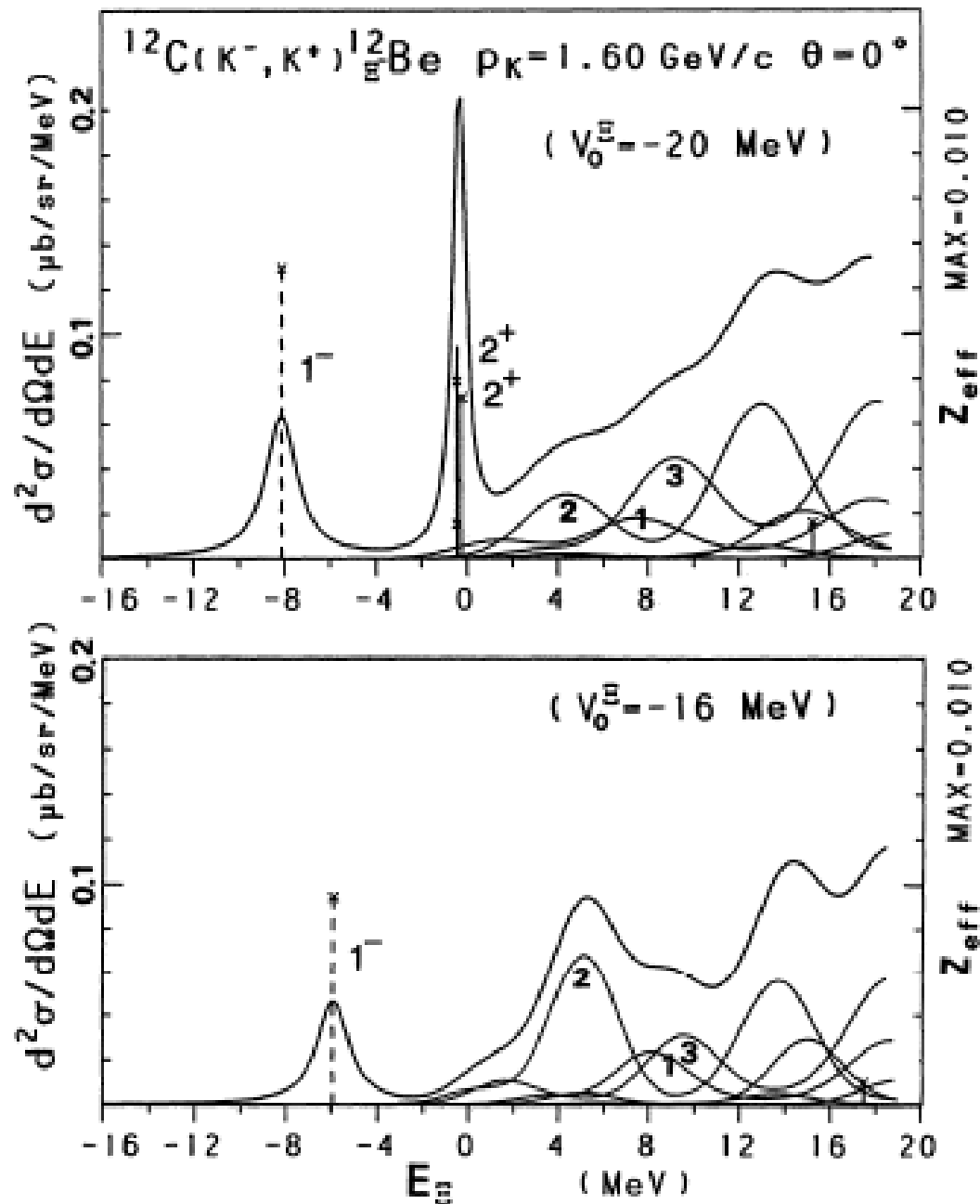


Far better energy resolution!

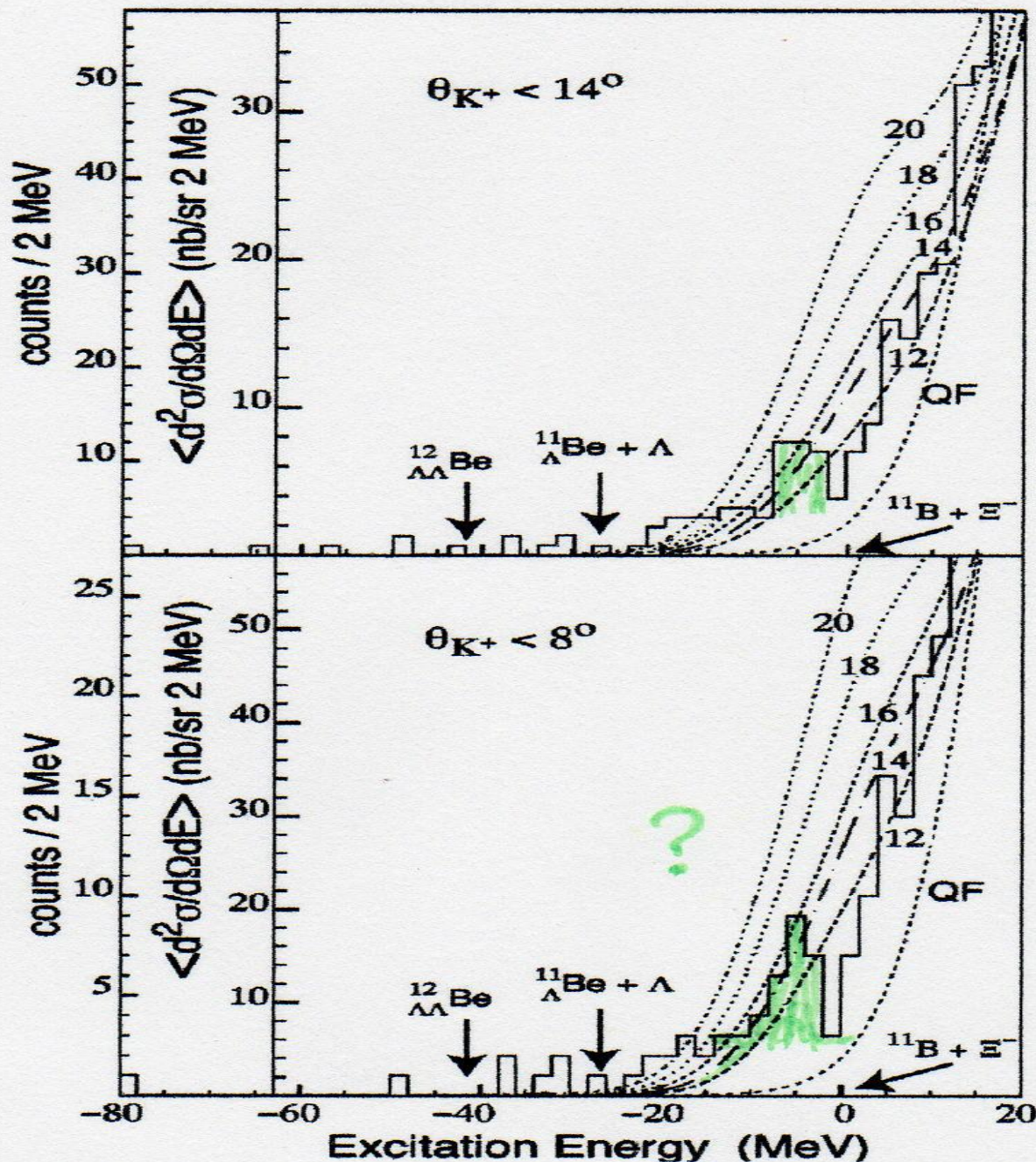
In comparison with  $(\pi^+, K^+)$  experiment from KEK

### 3. Production of $\Xi$ -hypernuclei (and $\Lambda\Lambda$ -hypernuclei )

- The similar theoretical framework has been applied to  $(K^-, K^+)$  reaction on  $^{12}\text{C}$ .
- DW: solution of the Klein-Gordon eq.
- taking account of nuclear core excitation effects



Arbitrary  
smearing  
widths are  
used in these  
figures



BNL-E885

P. Khaustov  
et al, PRC 61  
(2000)

BS strengths  
observed, but  
**peaks not  
confirmed.**

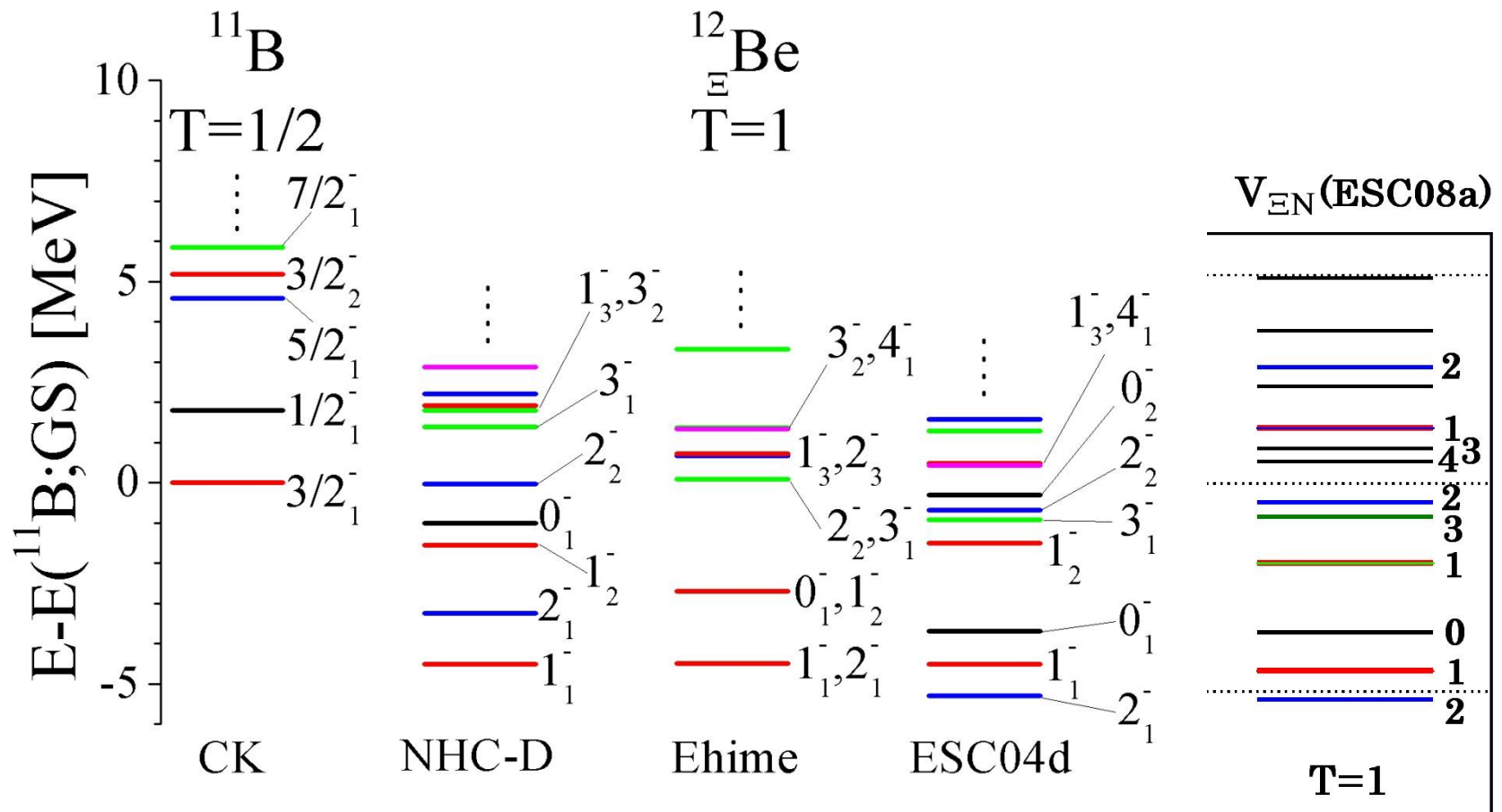
*Suggesting:*  
only WS-pot.  
depth:

$U=12-14$  MeV  
or less.

“shallow” 29

Reference position of  $J=1-(T=1)$  states.

The relative positions of  $J_n$  show the different  $\sigma.\sigma$  interaction nature of  $V_{\text{EN}}$ .



# $\Xi$ -N interactions used in Shell Model

Comparison in the form of  $V = -V_0 + \Delta(\sigma \cdot \sigma)$

		$V_0$	$\Delta$	$\Delta/V_0$
NE(ESC04d)	T=0	4.98	-15.81	-3.18
	T=1	0.30	-2.96	-9.88
NE(NHC-D)	T=0	2.14	4.75	2.23
	T=1	1.55	0.79	0.51
N $\Lambda$ (NSC97f)	T=1/2	1.05	0.04	0.04

$\sigma \cdot \sigma$  strengths are quite different for ESC and ND, so further trials and improvements are required.

# Sensitive interaction dependence

## Nijmegen NHC-D vs. ESC04d

- **Different partial-wave contributions**  
NHC-D (large p-state attraction) vs. ESC04d(s-state)
- **ESC04d (quite large spin- & isospin-dependence)**

Table 1:  $\Xi$  single particle energies  $U_{\Xi}$  and conversion widths  $\Gamma_{\Xi}$  at normal density calculated with ESC04d and NHC-D.  $S$ -state contributions in  $(TSLJ)$  states and total  $P$ -state contributions are also given. All entries are in MeV.

	$^{11}S_0$	$^{13}S_1$	$^{31}S_0$	$^{33}S_1$	$P$	$U_{\Xi}$	$\Gamma_{\Xi}$
ESC04d( $\alpha = 0$ )	6.4	-19.6	6.4	-5.0	-6.9	-18.7	11.4
ESC04d( $\alpha = .18$ )	6.3	-18.4	7.2	-1.7	-5.6	-12.1	12.7
NHC-D	-2.6	0.7	-2.3	-0.4	-16.8	-21.4	1.1

# Trying to use the most recent $\Xi$ -N interaction from Nijmegen (ESC08)

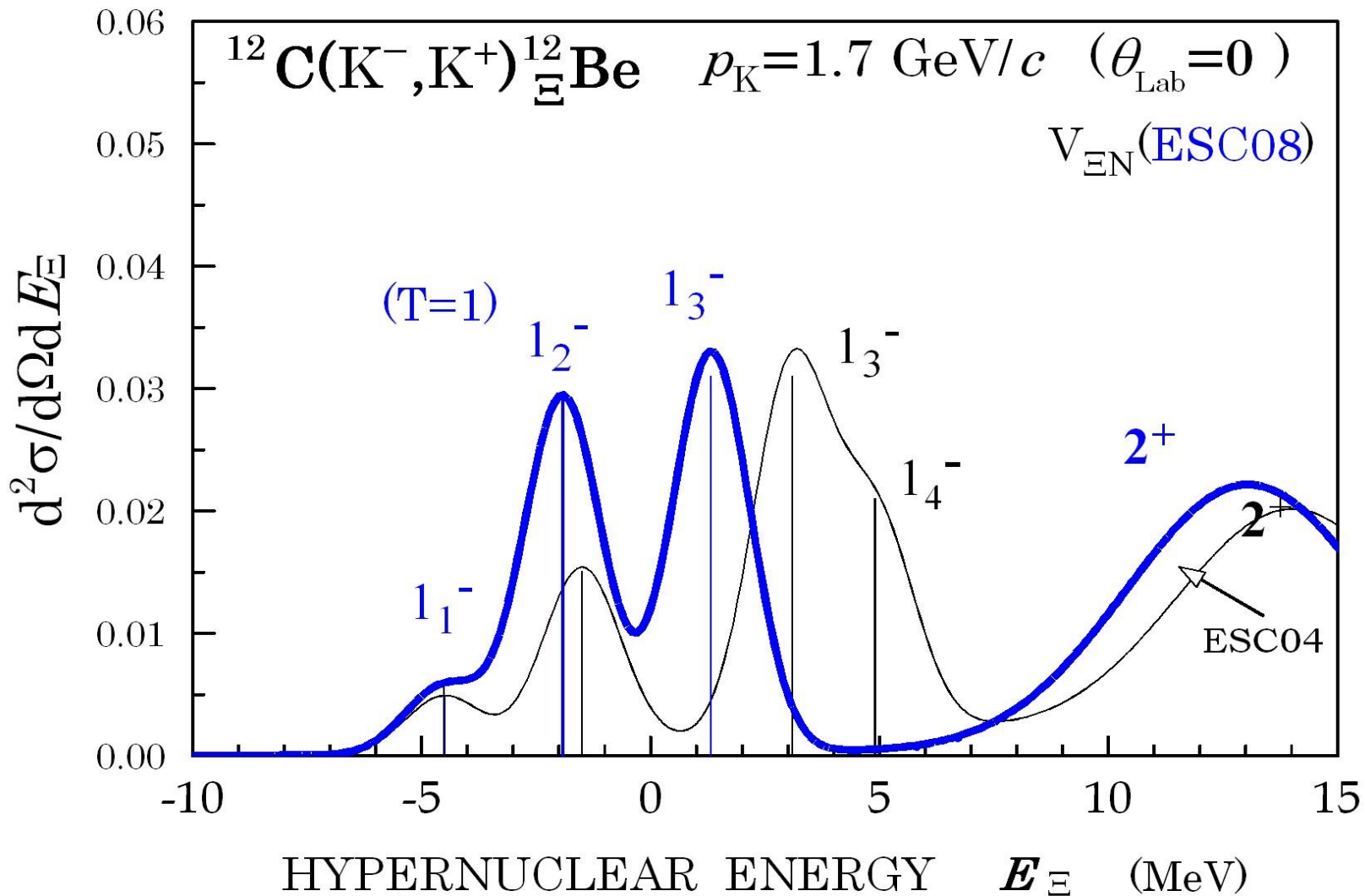
Table 1: Partial wave contributions to  $U_{\Xi}(\rho_0)$

model	$T$	$^1S_0$	$^3S_1$	$^1P_1$	$^3P_0$	$^3P_1$	$^3P_2$	$U_{\Xi}$	$\Gamma_{\Xi}$
ESC08	0	4.0	-2.7	0.2	-2.2	0.6	-1.1	-18.0	6.0
	1	7.0	-19.5	-0.3	0.1	-3.6	-0.7		
ESC04d	0	6.4	-19.6	1.1	1.2	-1.3	-2.0	-18.7	11.4
	1	6.4	-5.0	-1.0	-0.6	-1.4	-2.8		

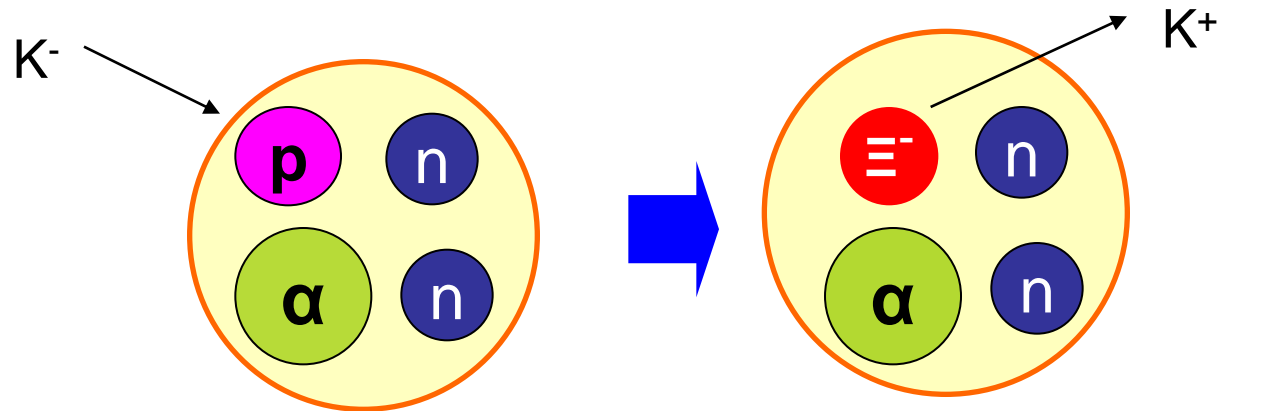
- Let us wait for the J-PARC Day-1 Exp. by Nagae et al, which will provide us a very important restriction on  $V(\Xi$ -N).

# ESC04 modified to be ESC08 !

(continuum bump should be larger)



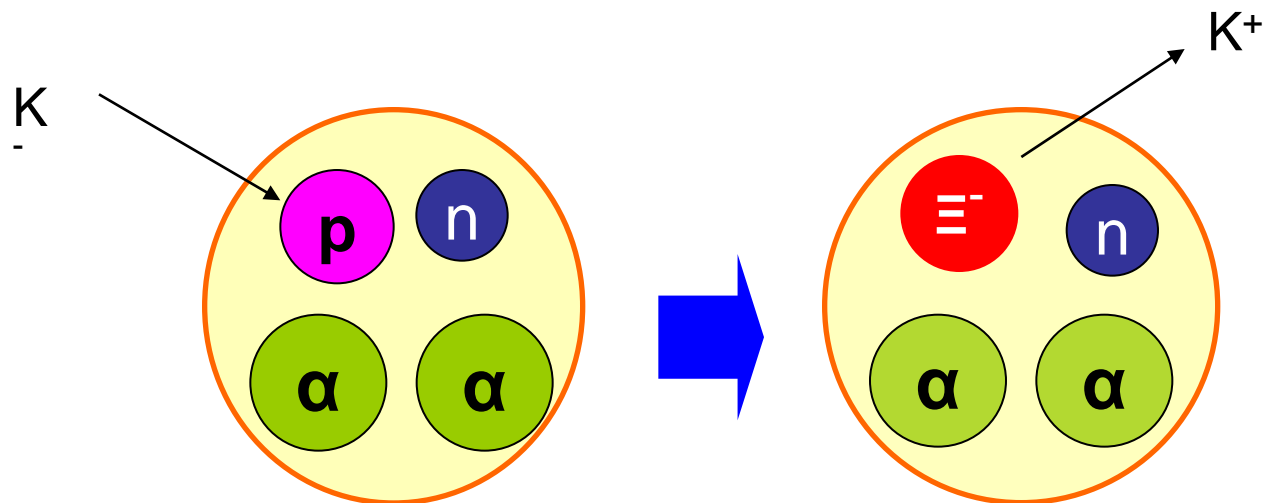
As the second best candidates to extract information about the spin-, isospin-independent term  $V_0$ , we propose to perform...



${}^7\text{Li} \ (T=1/2)$

${}^7_{\Xi^-}\text{H} \ (T=3/2)$

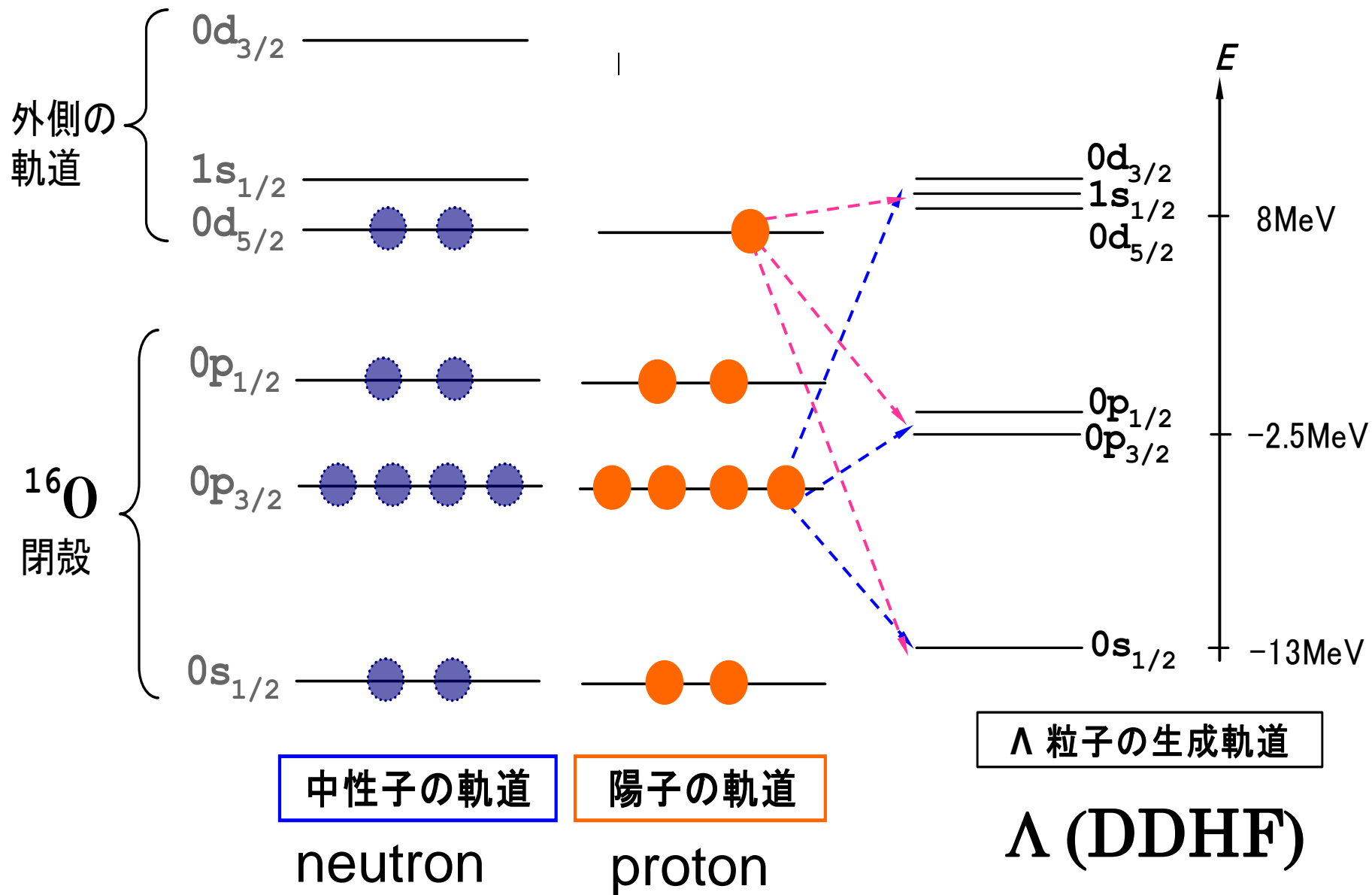
Why they are suited  
for investigating  $V_0$ ?



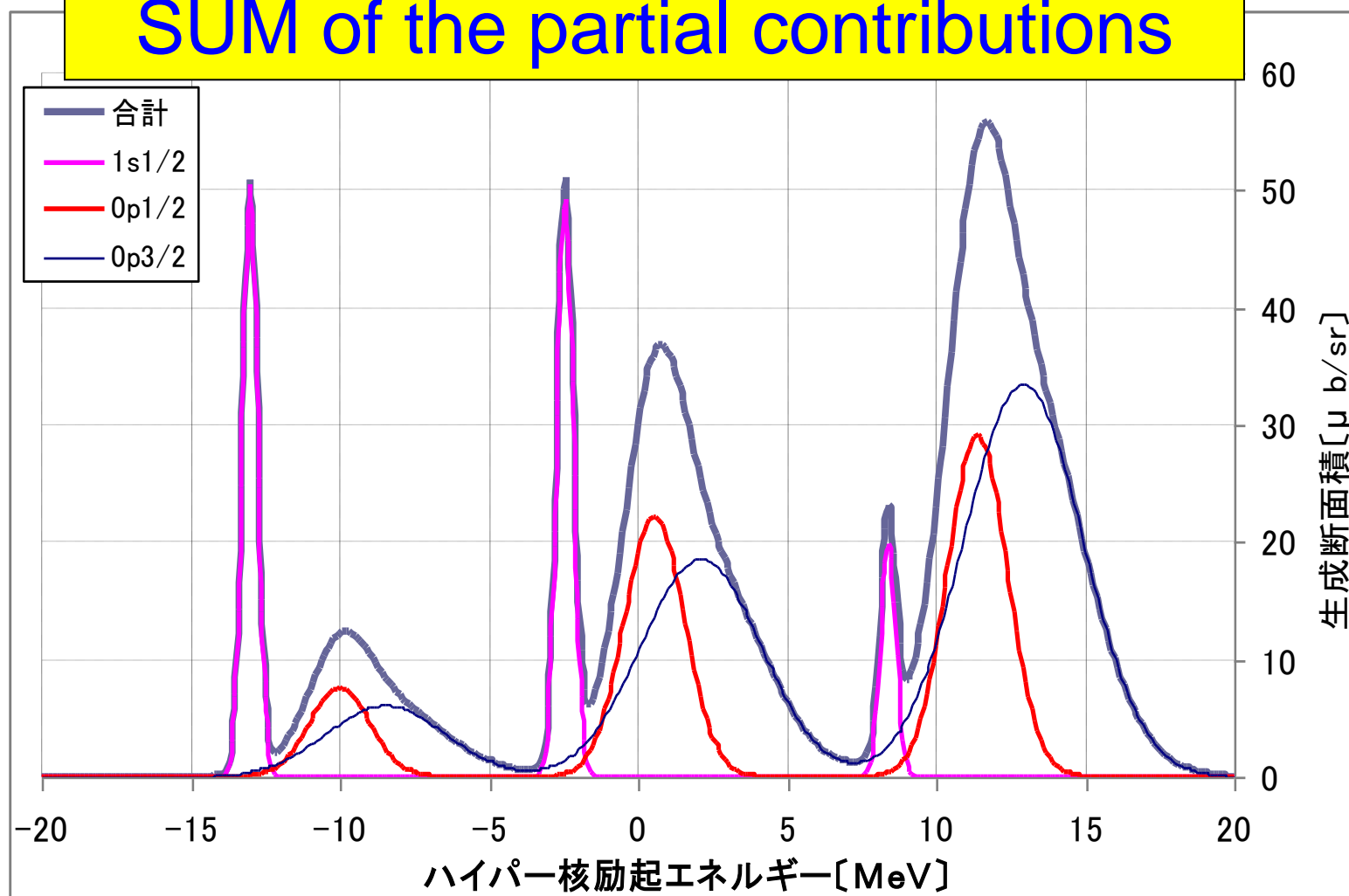
${}^{10}\text{B} \ (T=0)$

${}^{10}_{\Xi^-}\text{Li} \ (T=1)$

# Choose $^{19}\text{F}(1/2^+)$ target for demonstration

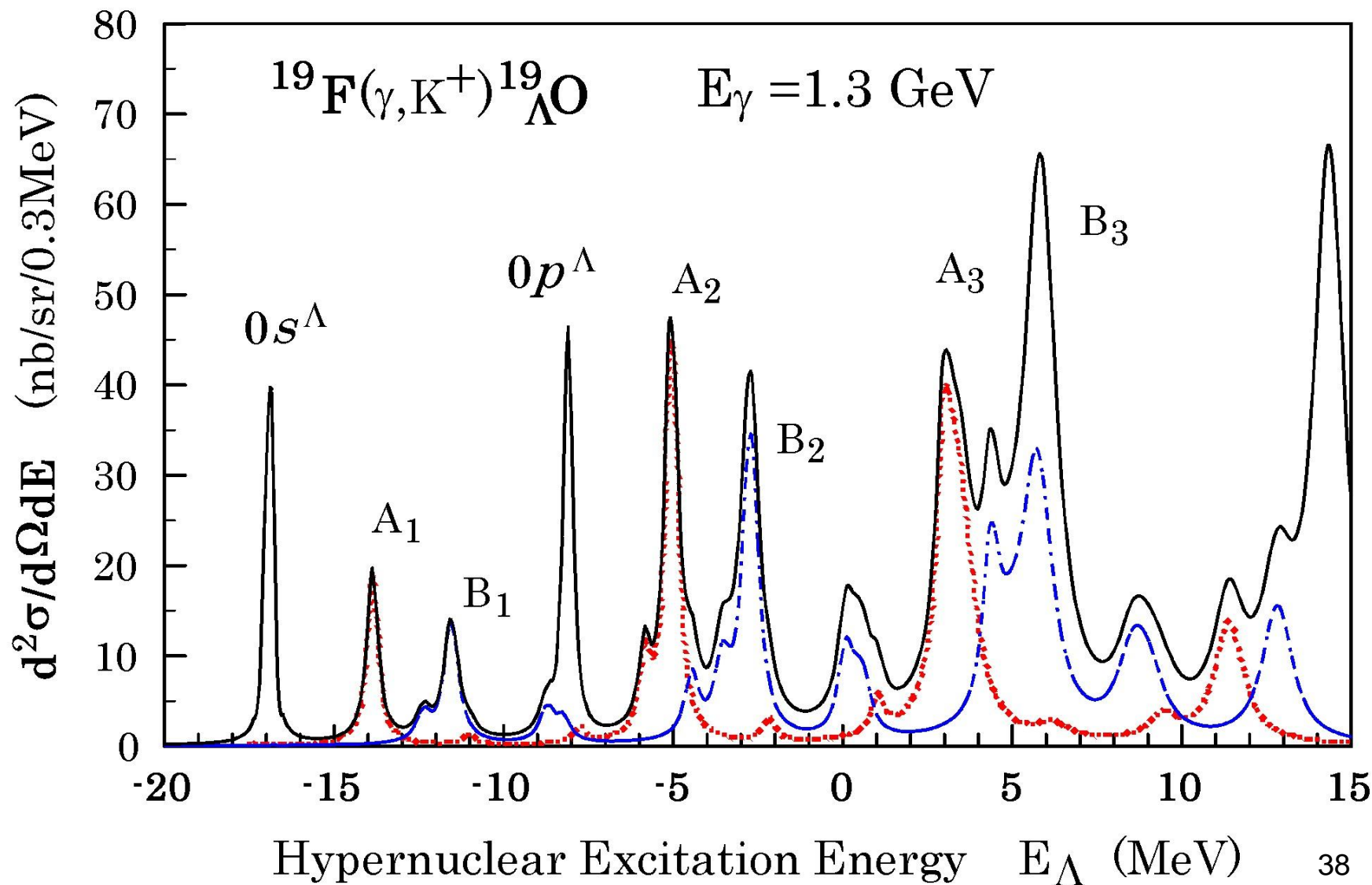


# $^{19}\text{F}(\gamma, K^+)_{\Lambda}^{19}\text{O}$ SUM of the partial contributions



As a “closed core ( $^{18}\text{O}$ )” +  $\Lambda$ , cf. SO-splitting(0p)=152+54 keV(C13)

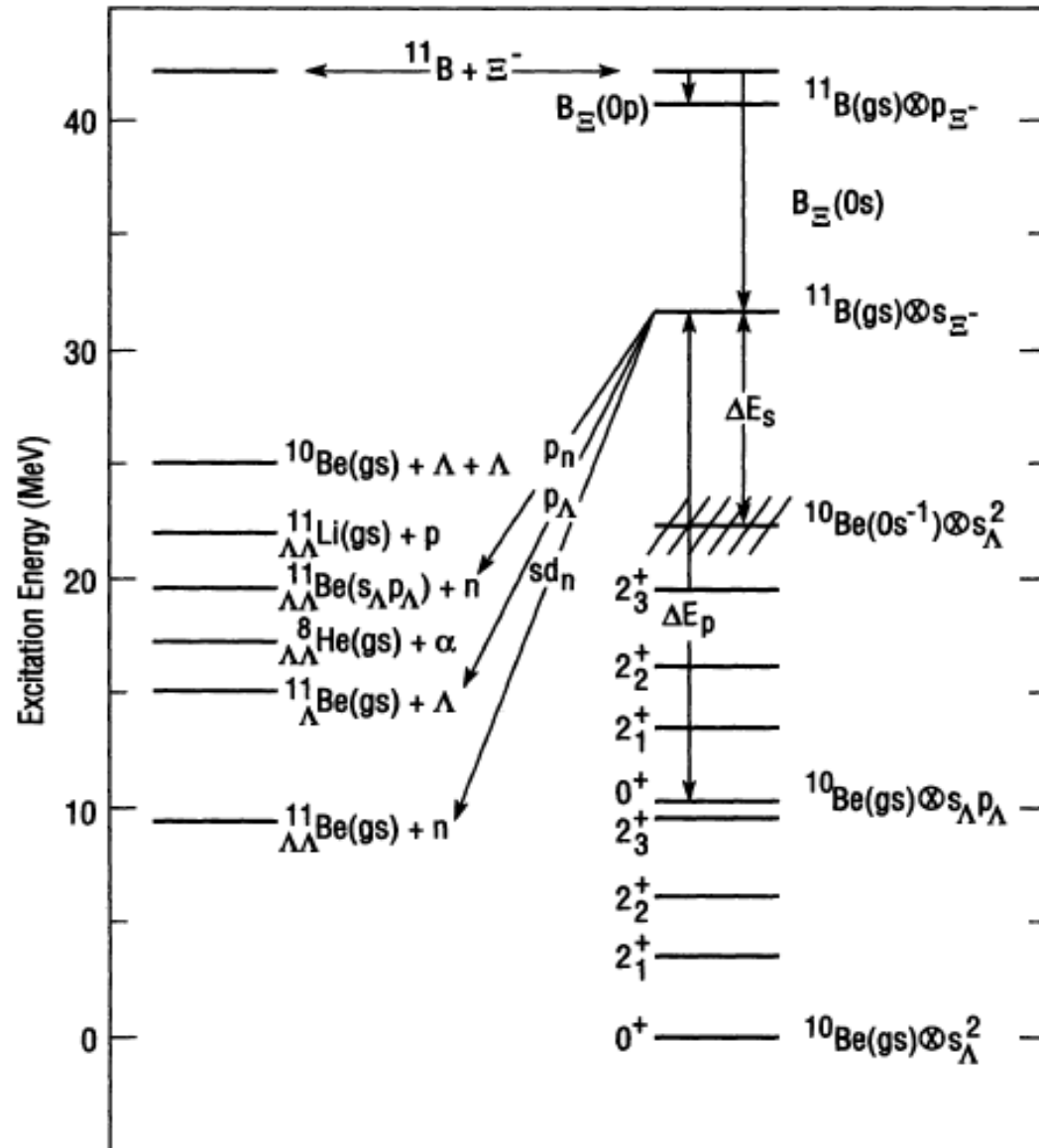
# Lightest sd-shell target: $^{19}\text{F}$



1-Skipped)  
discussion

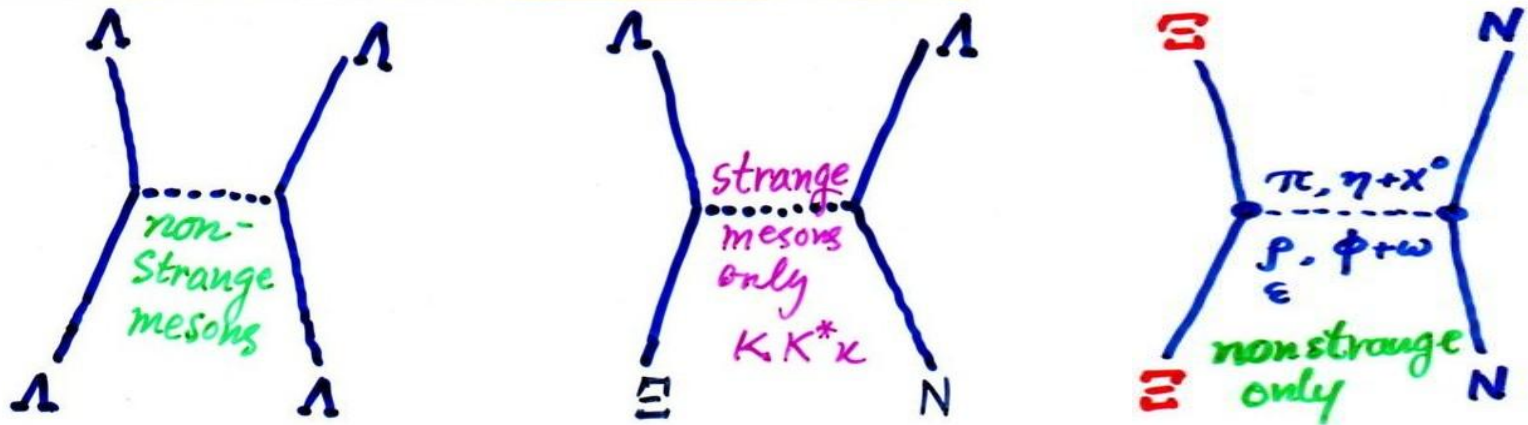
on

Widths of  
Xi states:  
Xi-LLmixing



# Why $\Xi$ -hypernuclei ?

1) They provide unique information on the  $S=-2$  B-B interactions inaccessible otherwise.



2) High-priority experiment at J-PARC (2009–  
E-05: “Spectroscopic study of X-hypernucleus via  
the  $^{12}\text{C}(K^-, K^+) \Xi^{12}\text{Be}$  reaction” by T. Nagae et al.

→ Realistic Calculations are required.

# Summary (proposal) for $\Xi$ -hypernuclear exp. programs to be done at J-PARC

important point is to extract a firm  
constraint on the  $\Xi$ -N **central** force  
(Meson theory shows strong spin-isospin  
dependence)

- To realize this purpose, I propose to  
choose odd-Z natural targets such as  
19F (18O+p), 27Al (26Mg+p),  
31P (30Si+p),..... (Cal. In progress

