

Kaonic

nuclei

Kaonic nuclei

at J-PARC

Prototype system = $K^- pp$



大強度陽子加速器施設 (J-PARC)

ハドロン実験施設

Kaonic nuclei at J-PARC

Theoretical situation of K^-pp study

KEK Theory Center / IPNS/ J-PARC branch

Akinobu Doté

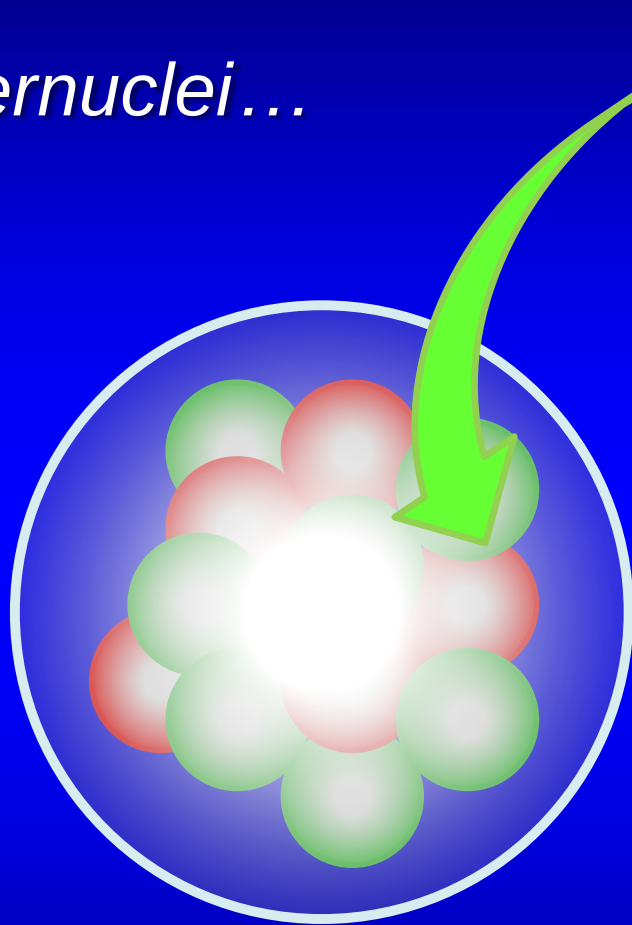
1. *Introduction*
2. *Recent theoretical studies of K^-pp*
3. *K^-pp studied with
“coupled-channel Complex Scaling Method + Feshbach method”*
4. *Issues on kaonic nuclei*
5. *Summary*

Workshop on J-PARC hadron physics in 2014,
10.Feb.'14 @ IQBRC, Tokai, Ibaraki, Japan

1. Introduction

Strange nuclear physics

Hypernuclei...



Nucleus



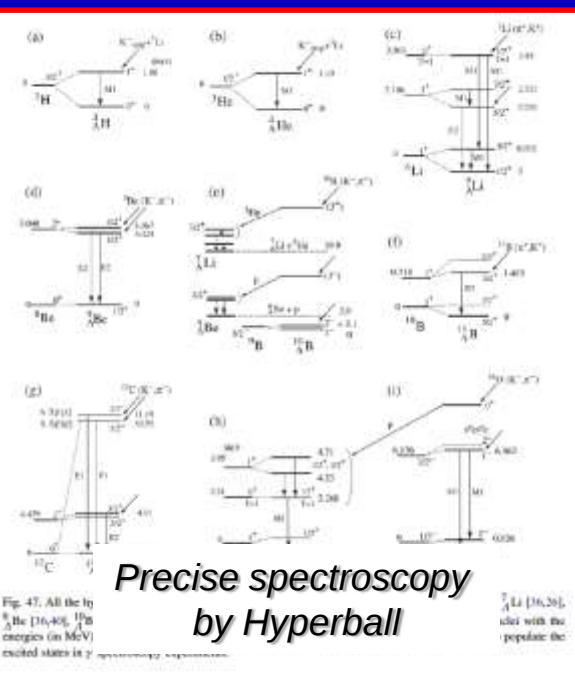
Hyperon

baryon = qqq

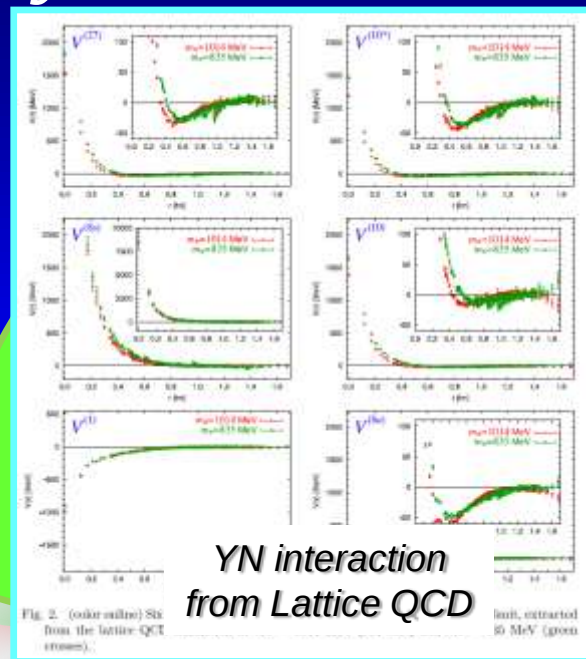
*Strangeness is introduced
through baryons.*

Strange nuclear physics

Hypernuclei...

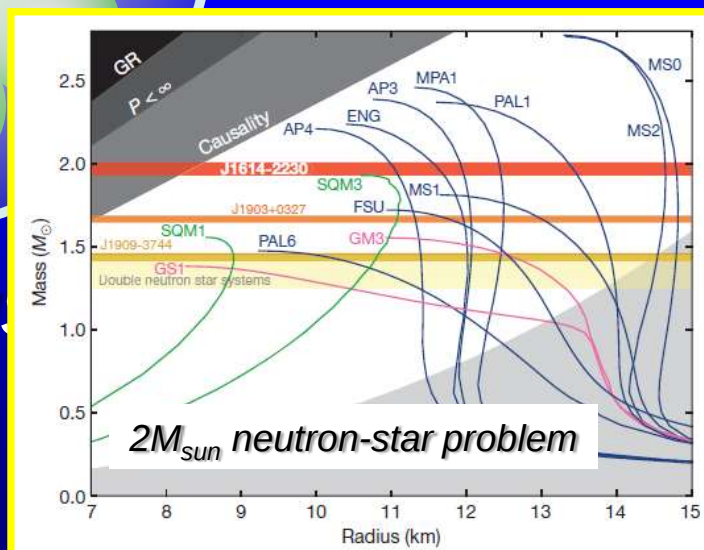


Nucleus



Hyperon

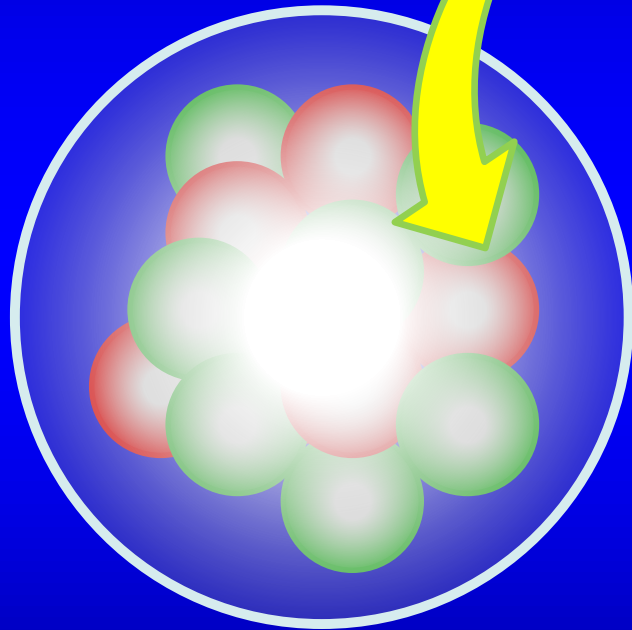
Hyperon = qqq



is introduced
ons.

Strange nuclear physics

Strangeness is introduced
through *mesons* ...



Nucleus



K^- meson
(anti-kaon, K^{bar})

$$\text{meson} = qq^{\text{bar}}$$

Kaonic nuclei !

Kaonic nuclei = Nuclear system with Anti-kaon “ K^{bar} ”

Attractive $K^{\text{bar}}N$ interaction!

Excited hyperon $\Lambda(1405)$ = a quasi-bound state of K^- and proton

- ✓ Difficult to explain by naive 3-quark model
- ✓ Consistent with “repulsive nature” indicated by $K^{\text{bar}}N$ scattering length and 1s level shift of kaonic hydrogen atom

$\Lambda(1405) \ 1/2^-$

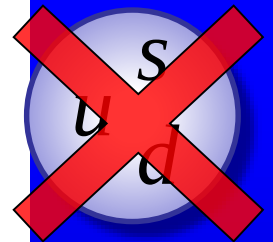
$$I(J^P) = 0(\frac{1}{2}^-)$$

Mass $m = 1405.1^{+1.3}_{-1.0}$ MeV

Full width $\Gamma = 50 \pm 2$ MeV

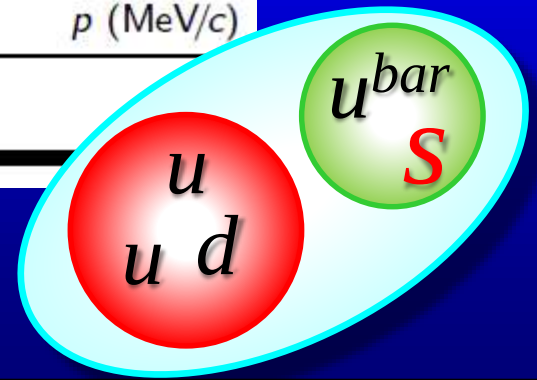
Below $\overline{K}N$ threshold

$K^{\text{bar}}N$ threshold = 1435 MeV



$\Lambda(1405)$ DECAY MODES

	Fraction (Γ_i/Γ)	p (MeV/c)
$\Sigma \pi$	100 %	



Kaonic nuclei = Exotic system !?

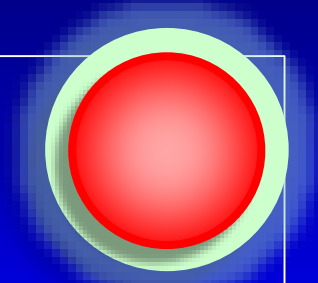
Phenomenologically derived $I=0$ $K^{\text{bar}}N$ potential ... Very attractive



Deeply bound (Total B.E. $\sim 100\text{MeV}$)

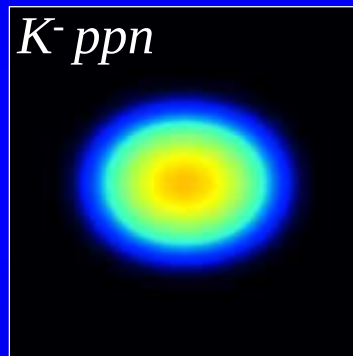
Highly dense state formed in a nucleus

Interesting structures that we have never seen in normal nuclei...



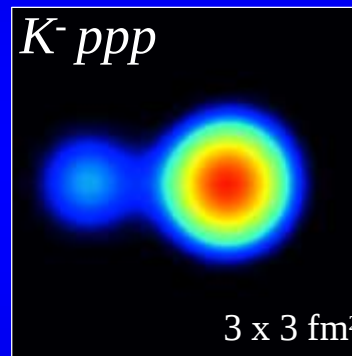
0.0

0.15



0.0

2.0



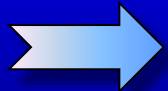
0.0

2.0

$[\text{fm}^{-3}]$

*Antisymmetrized Molecular Dynamics
method with
a phenomenological $K^{\text{bar}}N$ potential*

*A. Dote, H. Horiuchi, Y. Akaishi and
T. Yamazaki, PRC70, 044313 (2004)*



Relate to various interesting physics such as ...

- Restoration of chiral symmetry in dense matter
- Interesting structure
- Neutron star

K^-

$\Lambda(1405)$

“Building block of kaonic nuclei”

... used to determine $K^{\text{bar}}N$ interaction

Proton

K^-pp

P

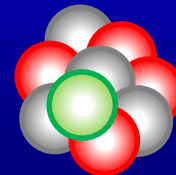
K^-

P

Most essential kaonic nucleus

*“Prototype of
kaonic nuclei”*

Nuclear many-body system with K^-



$^3\text{He}K^-$, $pppK^-$,
 $^4\text{He}K^-$, $pppnK^-$,
..., $^8\text{Be}K^-$, ...

2. Recent theoretical studies of K^-pp

Theoretical studies of K^-pp

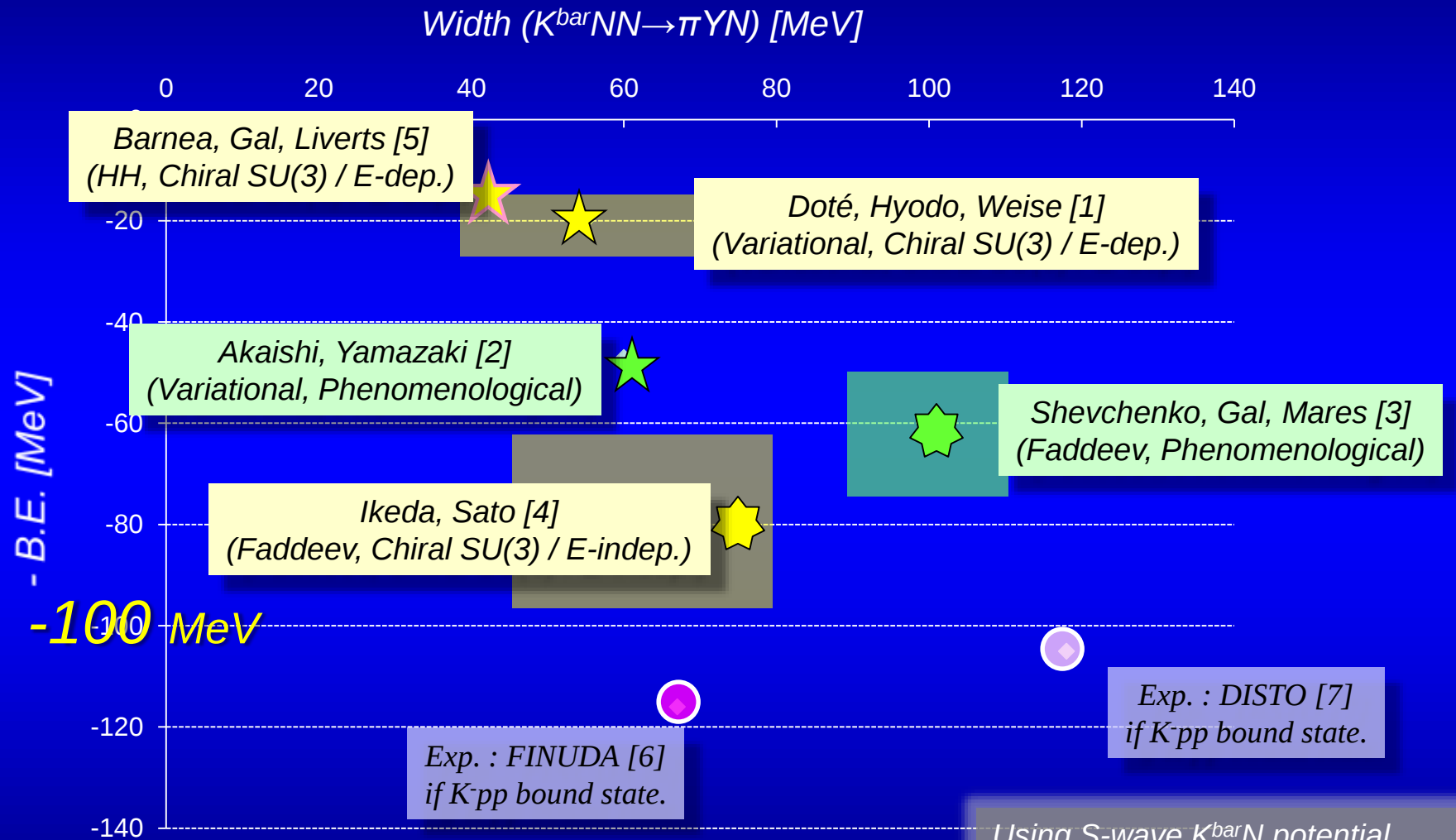
“ K^-pp ” = Prototype of K^{bar} nuclei

$K^{\text{bar}}NN - \pi\Sigma N - \pi\Lambda N$, $J^P=1/2^-$, $T=1/2$

- Doté, Hyodo, Weise PRC79, 014003(2009)
Variational with a chiral SU(3)-based $K^{\text{bar}}N$ potential
- Akaishi, Yamazaki PRC76, 045201(2007)
ATMS with a phenomenological $K^{\text{bar}}N$ potential
- Ikeda, Sato PRC76, 035203(2007)
Faddeev with a chiral SU(3)-derived $K^{\text{bar}}N$ potential
- Shevchenko, Gal, Mares PRC76, 044004(2007)
Faddeev with a phenomenological $K^{\text{bar}}N$ potential
- Barnea, Gal, Liverts PLB712, 132(2012)
Hyperspherical harmonics with a chiral SU(3)-based $K^{\text{bar}}N$ potential
- Wycech, Green PRC79, 014001(2009)
Variational with a phenomenological $K^{\text{bar}}N$ potential (with p-wave)
- Arai, Yasui, Oka/ Uchino, Hyodo, Oka PTP119, 103(2008)
 Λ^* nuclei model /PTPS 186, 240(2010)
- Nishikawa, Kondo PRC77, 055202(2008)
Skyrme model

All studies predict that K^-pp can be bound!

Typical results of theoretical studies of K^-pp



[1] PRC79, 014003 (2009)

[5] PLB94, 712 (2012)

[6] PRL94, 212303 (2005)

[2] PRC76, 045201 (2007)

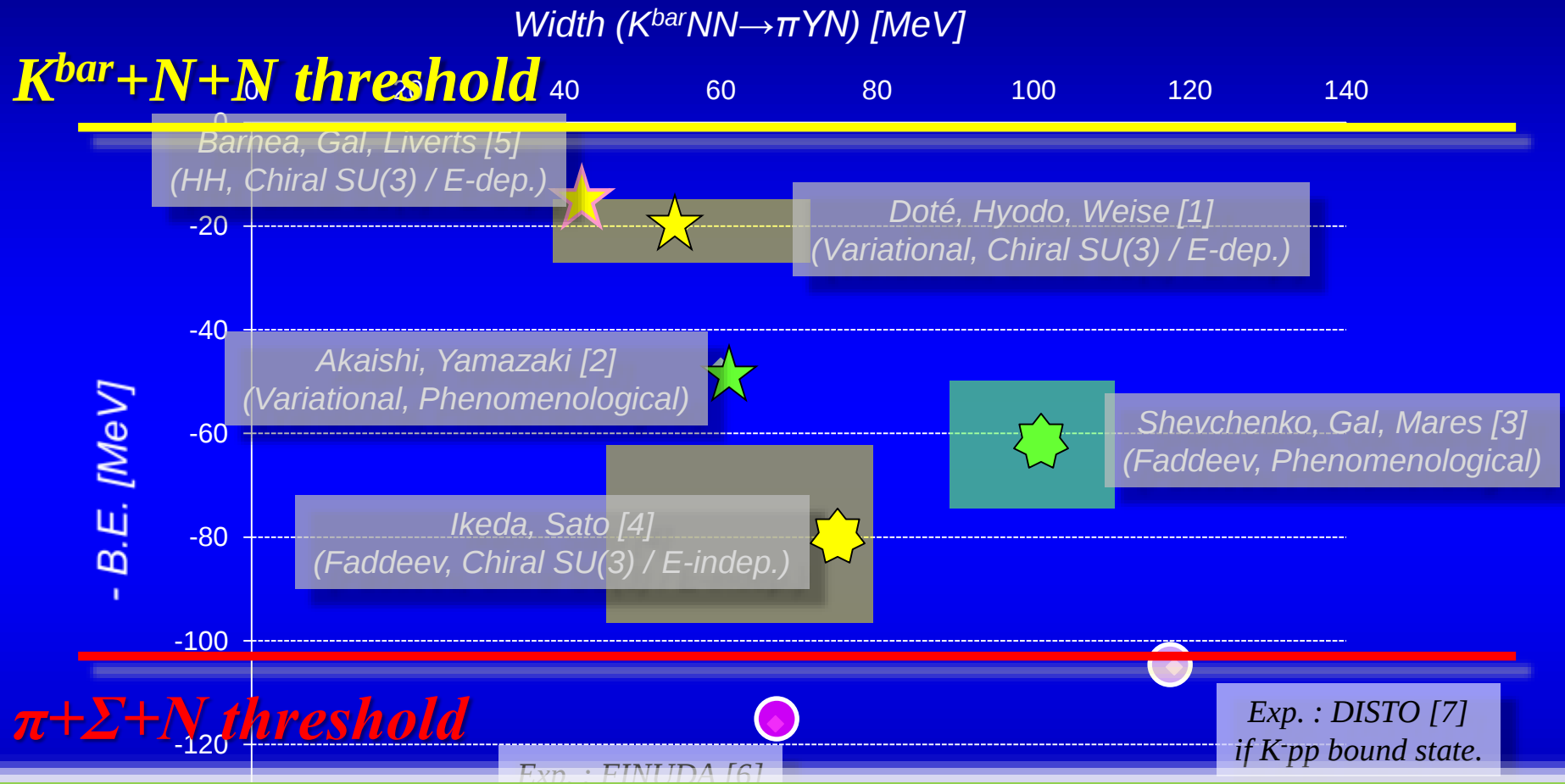
[7] PRL104, 132502 (2010)

[3] PRC76, 044004 (2007)

[4] PRC76, 035203 (2007)

Using S-wave $K^{\text{bar}}N$ potential
constrained by experimental data.
... $K^{\text{bar}}N$ scattering data,
Kaonic hydrogen atom data,
“ $\Lambda(1405)$ ” etc.

Typical results of theoretical studies of K^-pp



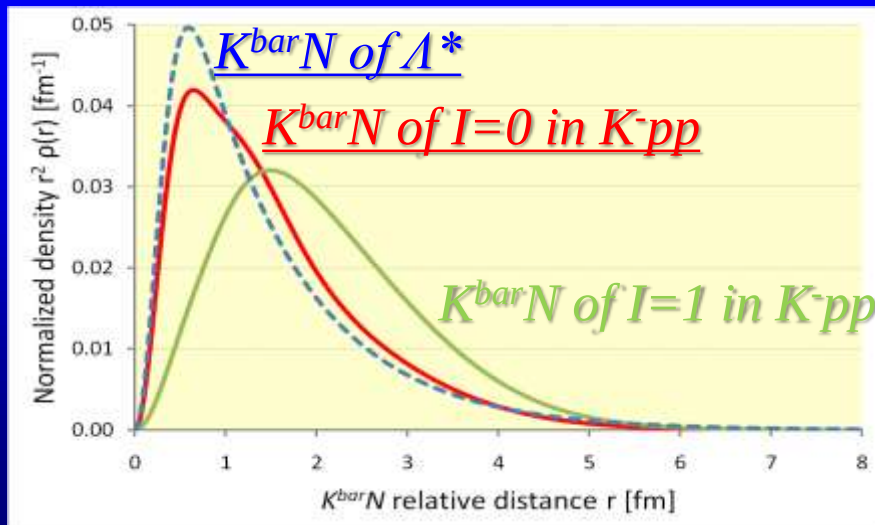
From theoretical viewpoint,
 K^-pp exists between $K^{\text{bar}}-N-N$ and $\pi-\Sigma-N$ thresholds!

Typical results of theoretical studies of K^-pp

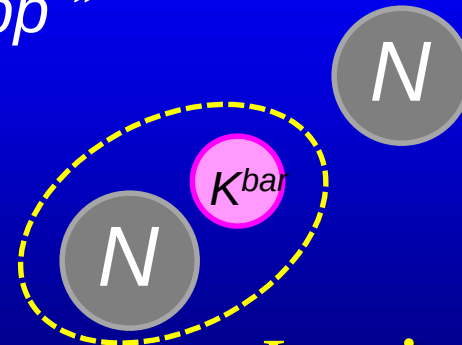
	DHW	AY	BGL	IS	SGM
B(K^-pp)	20 ± 3	47	16	60 ~ 95	50 ~ 70
Width Γ	40 ~ 70	61	41	45 ~ 80	90 ~ 110
Method	Variational (Gauss)	Variational (Gauss)	Variational (HH)	Faddeev-AGS	Faddeev-AGS
Potential	Chiral (E-dep.)	Pheno.	Chiral (E-dep.)	Chiral (E-indep.)	Pheno.
Kinematics	Non-rel.	Non-rel.	Non-rel.	Rel.	Non-rel.

DHW vs AY

- Difference of the potential: chiral (E-dep.) = shallow, phenomenological = deep
- “Survival of Λ^* in K^-pp ”



“ K^-pp ”



Isospin-0 $\sim \Lambda^*$!

Typical results of theoretical studies of K^-pp

	DHW	AY	BGL	IS	SGM
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DHW vs BGL

DHW is checked by BGL with a different base (Hyperspherical Harmonics)

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- Difference of the potential: chiral (E-dep.) = shallow, phenomenological = deep
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DHW vs BGL

DHW is checked by BGL with a different base (Hyperspherical Harmonics)

Variational (DHW, AY, BGL) vs Faddeev-AGS (IS, SGM)

- $\pi\Sigma N$ three-body dynamics included in Faddeev-AGS
- Decay width: perturbative in variational cal., non-perturbative in F-A

Typical results of theoretical studies of K^-pp

	DHW	AY	BGL	IS	SGM
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Pot.	Chiral (E-dep.)	Pheno.	Chiral (E-dep.)	Chiral (E-indep.)	Pheno.
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Further study of IS

Y. Ikeda, H. Kamano, T. Sato, PTP124, 533 (2010)

Typical results of theoretical studies of K^-pp

	DHW	AY	BGL	IS	SGM
B(K^-pp)	20 ± 3	47	16	44 ~ 58	50 ~ 70
Width Γ	40 ~ 70	61	41	34 ~ 40	90 ~ 110
Method	Variational (Gauss)	Variational (Gauss)	Variational (HH)	Faddeev-AGS	Faddeev-AGS
Pot.	Chiral (E-dep.)	Pheno.	Chiral (E-dep.)	Chiral (E-indep.)	Pheno.
Kinematics	Non-rel.	Non-rel.	Non-rel.	Non-rel.	Non-rel.

DHW vs AY

- Difference of the potential: chiral (E-dep.) = shallow, phenomenological = deep
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Further study of IS

Y. Ikeda, H. Kamano, T. Sato, PTP124, 533 (2010)

- Kinematics: Rel. $\rightarrow$ Non-rel. $\Rightarrow (B, \Gamma) = (44 \sim 58, 34 \sim 40)$

Typical results of theoretical studies of K^-pp

	DHW	AY	BGL	IS	SGM
B(K^-pp)	20 ± 3	47	16	9 ~ 16	50 ~ 70
Width Γ	40 ~ 70	61	41	34 ~ 46	90 ~ 110
Method	Variational (Gauss)	Variational (Gauss)	Variational (HH)	Faddeev-AGS	Faddeev-AGS
Potential	Chiral (E-dep.)	Pheno.	Chiral (E-dep.)	Chiral (E-dep.)	Pheno.
Kinematics	Non-rel.	Non-rel.	Non-rel.	Non-rel.	Non-rel.

DHW vs AY

- Difference of the potential: chiral (E-dep.) = shallow, phenomenological = deep
- “Survival of Λ^* in K^-pp ”

DHW vs BGL

DHW is checked by BGL with a different base (Hyperspherical Harmonics)

Variational (DHW, AY, BGL) vs Faddeev-AGS (IS, SGM)

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- Decay width: perturbative in variational cal., non-perturbative in F-A

Further study of IS

Y. Ikeda, H. Kamano, T. Sato, PTP124, 533 (2010)

- Kinematics: Rel. $\rightarrow$ Non-rel. $\Rightarrow (B, \Gamma) = (44 \sim 58, 34 \sim 40)$
- Energy-dependence of chiral potential: E-indep. $\rightarrow$ E-dep. $\Rightarrow (B, \Gamma) = (9 \sim 16, 34 \sim 46)$
 $(B, \Gamma) = (67 \sim 89, 244 \sim 320) \dots$ another pole

*3. K - pp studied with
“coupled-channel
Complex Scaling Method
+ Feshbach method”*

Takashi Inoue (Nihon univ.)

Takayuki Myo (Osaka Inst. Tech. univ.)

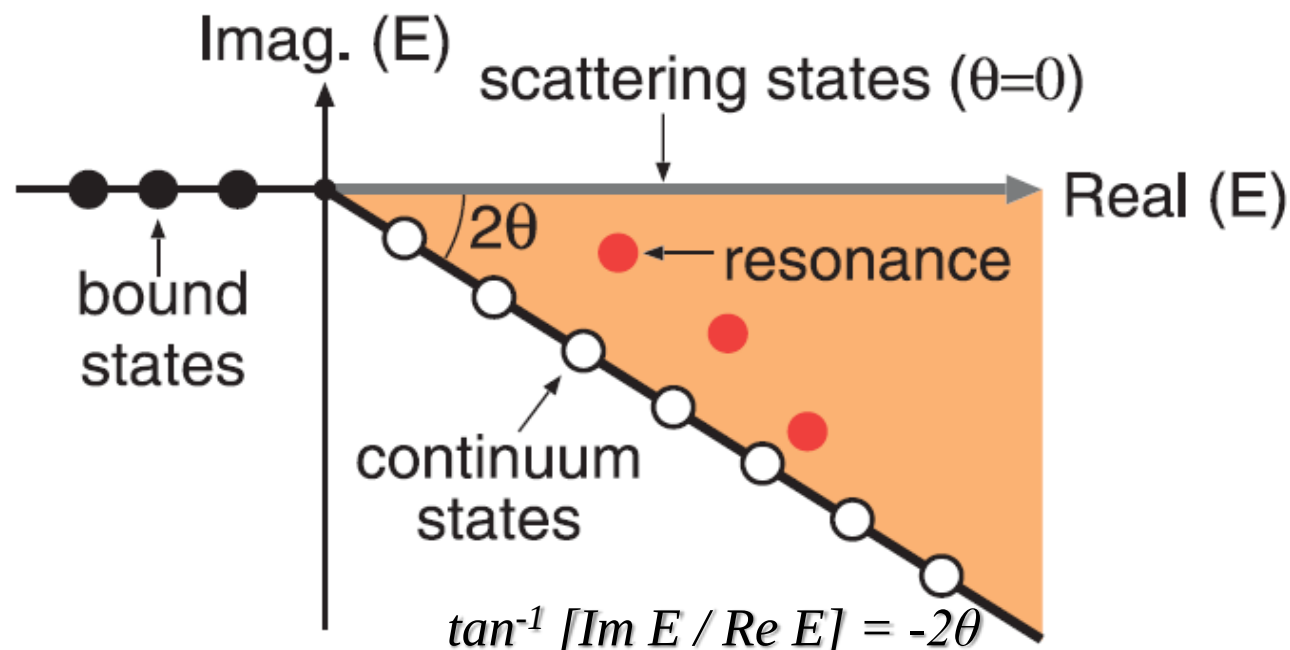
Complex Scaling Method

= Powerful tool for resonance study

Complex rotation of coordinate
(Complex scaling)

$$U(\theta): \mathbf{r} \rightarrow \mathbf{r} e^{i\theta}, \quad \mathbf{k} \rightarrow \mathbf{k} e^{-i\theta}$$

Diagonalize $H_\theta = U(\theta) H U^{-1}(\theta)$ with Gaussian base,



- Continuum state appears on 2ϑ line.
- Resonance pole is off from 2ϑ line, and independent of ϑ .

S. Aoyama, T. Myo, K. Kato
and K. Ikeda,
PTP116, 1 (2006)

ccCSM + Feshbach method

Elimination of channels by Feshbach method

Schrödinger eq.
in model space “P” and out of model space “Q”

$$\begin{pmatrix} T_P + v_P & V_{PQ} \\ V_{QP} & T_Q + v_Q \end{pmatrix} \begin{pmatrix} \Phi_P \\ \Phi_Q \end{pmatrix} = E \begin{pmatrix} \Phi_P \\ \Phi_Q \end{pmatrix}$$

Schrödinger eq. in P space :

$$(T_P + U_P^{\text{Eff}}(E)) \Phi_P = E \Phi_P$$

Effective potential for P-space

$$U_P^{\text{Eff}}(E) = v_P + V_{PQ} G_Q(E) V_{QP}$$

Q-space Green function:

$$G_Q(E) = \frac{1}{E - H_{QQ}}$$

Extended Closure Relation in Complex Scaling

$$H_{QQ}^\theta |\chi_n^\theta\rangle = \varepsilon_n^\theta |\chi_n^\theta\rangle$$

$$H_{QQ}^\theta = U(\theta) H_{QQ} U^{-1}(\theta)$$

$$\int_C \sum_{R+B} |\chi_n^\theta\rangle \langle \chi_n^\theta| = 1$$

Diagonalize H_{QQ}^θ with Gaussian base,

$$\sum_n |\chi_n^\theta\rangle \langle \chi_n^\theta| \approx 1$$

Well approximated

T. Myo, A. Ohnishi and K. Kato, PTP99, 801 (1998)

Express the $G_Q(E)$ with Gaussian base using ECR

$$U_P^{\text{Eff}}(E) = v_P + V_{PQ} \underbrace{U^{-1}(\theta) G_Q^\theta(E) U(\theta)}_{G_Q(E)} V_{QP}$$

$$G_Q^\theta(E) \approx \sum_n |\chi_n^\theta\rangle \frac{1}{E - \varepsilon_n^\theta} \langle \chi_n^\theta|$$

$\{|\chi_n^\theta\rangle\}$: expanded with Gaussian base.

Apply ccCSM + Feshbach method to K^-pp

“ K^-pp ” ... $K^{bar}NN - \pi\Sigma N - \pi\Lambda N$ ($J^\pi=0^-, T=1/2$)

In the two-body system, $P = K^{bar}N$, $Q = \pi Y$

$$\begin{matrix} V(K^{bar}N - \pi Y; I=0,1) \\ V(\pi Y - \pi Y'; I=0,1) \end{matrix} \xrightarrow{\text{Feshbach + ccCSM}} U_{K^{bar}N(I=0,1)}^{Eff}(E)$$

- Schrödinger eq. for $K^{bar}NN$ channel :

$$\left(T_{K^{bar}NN} + V_{NN} + \sum_{i=1,2} U_{K^{bar}N_i(I)}^{Eff}(E_{K^{bar}N}) \right) \Phi_{K^{bar}NN} = E \Phi_{K^{bar}NN}$$

- Trial wave function

$$\begin{aligned} |"K^-pp"> &= \sum_a C_a^{(KNN,1)} \left\{ G_a^{(KNN,1)}(\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) + G_a^{(KNN,1)}(-\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) \right\} |S_{NN}=0> \left| [K[NN]_1]_{T=1/2} \right> && \text{Ch. 1: } K^{bar}NN, \quad NN: ^1E \\ &+ \sum_a C_a^{(KNN,2)} \left\{ G_a^{(KNN,2)}(\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) - G_a^{(KNN,2)}(-\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) \right\} |S_{NN}=0> \left| [K[NN]_0]_{T=1/2} \right> && \text{Ch. 2: } K^{bar}NN, \quad NN: ^1O \end{aligned}$$

- Basis function = Correlated Gaussian
...including 3-types Jacobi-coordinates

$$G_a^{(KNN,i)}(\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) = N_a^{(KNN,i)} \exp \left[-(\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) A_a^{(KNN,i)} \begin{pmatrix} \mathbf{x}_1^{(3)} \\ \mathbf{x}_2^{(3)} \end{pmatrix} \right]$$

Apply ccCSM + Feshbach method to K^-pp

“ K^-pp ” ... $K^{\text{bar}}NN - \pi\Sigma N - \pi\Lambda N$ ($J^\pi=0^-, T=1/2$)

Chiral SU(3)-based potential

A. D., T. Inoue, T. Myo, Nucl. Phys. A 912, 66 (2013)

$$V_{ij}^{(I=0,1)}(r) = -\frac{C_{ij}^{(I=0,1)}}{8f_\pi^2} (\omega_i + \omega_j) \sqrt{\frac{M_i M_j}{s \omega_i \omega_j}} g_{ij}(r) \quad g_{ij}(r) = \frac{1}{\pi^{3/2} d_{ij}^3} \exp\left[-(r/d_{ij})^2\right]$$

- Flux factor ansatz to make a potential from an interaction kernel

N. Kaiser, P. B. Siegel and W. Weise, NPA 594, 325 (1995)

- Based on Chiral SU(3) theory → **Energy dependence**
- WT term, **Gaussian form in r-space**
- Constrained by $K^{\text{bar}}N$ scattering length

$$a_{KN(I=0)} = -1.70 + i0.67 \text{ fm}, \quad a_{KN(I=1)} = 0.37 + i0.60 \text{ fm} \quad (\text{A.D.Martin, NPB179, 33(1979)})$$

Self-consistency for **complex** $K^{\text{bar}}N$ energy

A. D., T. Hyodo, W. Weise, PRC79, 014003(2009)

... Similarly to the DHW study

$$1. \text{ Kaon's binding energy: } B(K) \equiv -\left\{ \langle H \rangle - \langle H_{\text{Nucl}} \rangle \right\} \quad H_{\text{Nucl}} : \text{Hamiltonian of nuclear part}$$

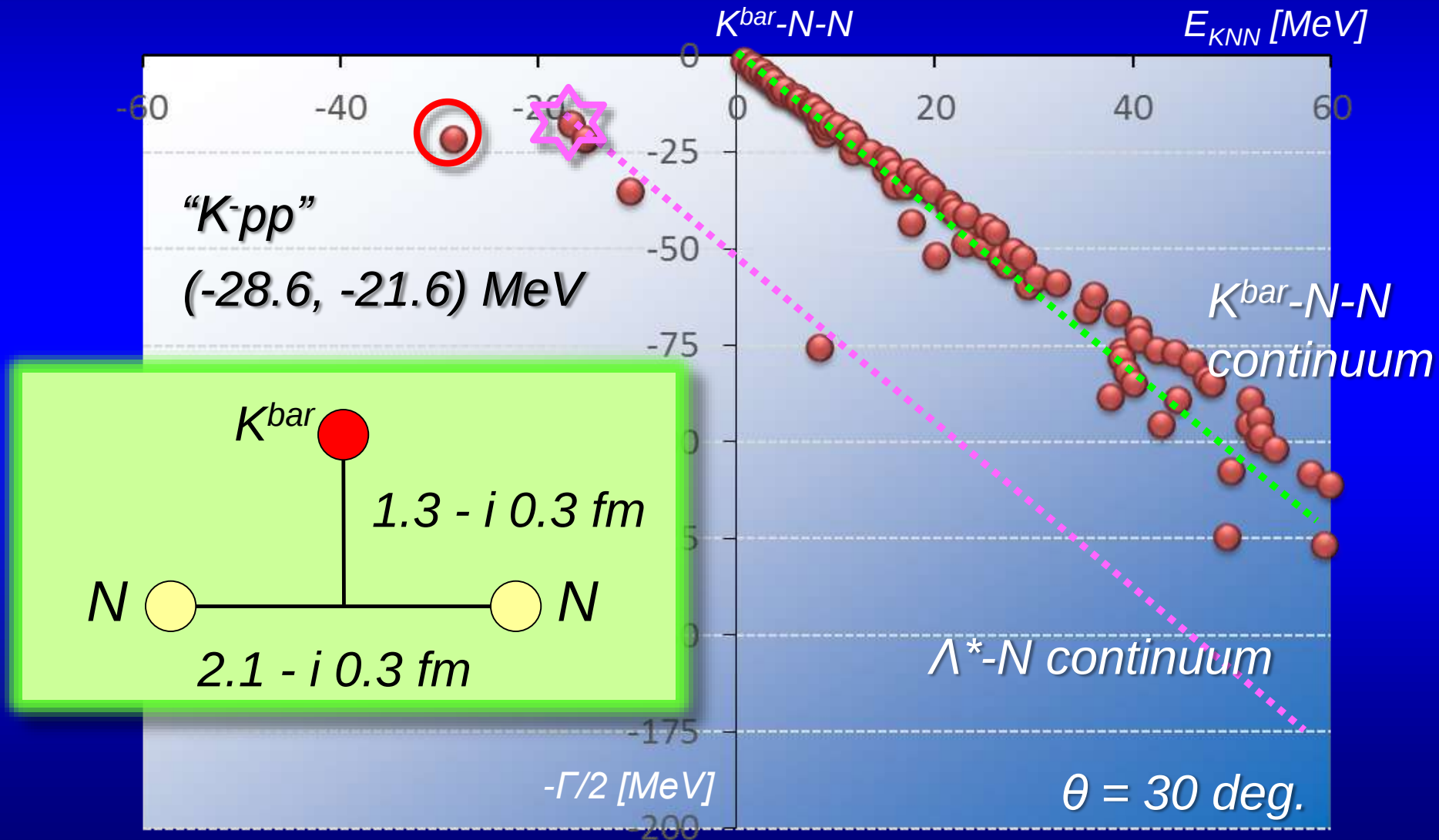
2. Define $K^{\text{bar}}N$ energy in two ways

$$E_{KN} = M_N + \omega = \begin{cases} M_N + m_K - B(K) & : \text{Field picture} \\ M_N + m_K - B(K)/2 & : \text{Particle picture} \end{cases}$$

Result: KSW-NRv2 potential ($f_\pi=110\text{MeV}$)

NN potential : Av18 (Central + spin-spin)

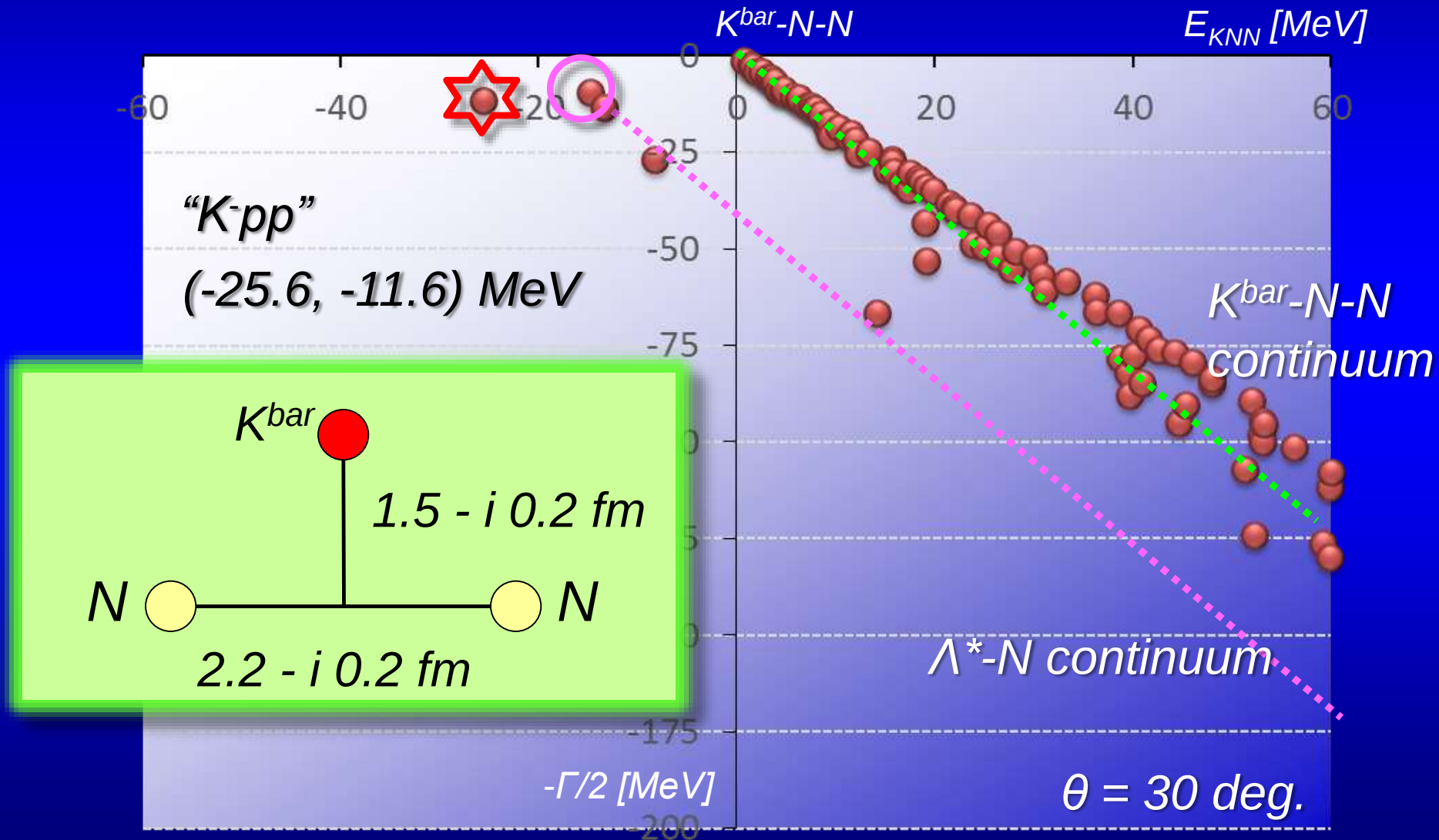
$K^{\text{bar}}N$ energy is fixed at Λ^* . (self-consistent at Λ^*)



Result: KSW-NRv2 potential ($f_\pi=110\text{MeV}$)

NN potential : Av18 (Central + spin-spin)

$K^{\text{bar}}N$ energy is self-consistent at K -pp. (Ansatz 1)



4. Issues on kaonic nuclei

i) $K^{\text{bar}}N$ interaction, $\Lambda(1405)$

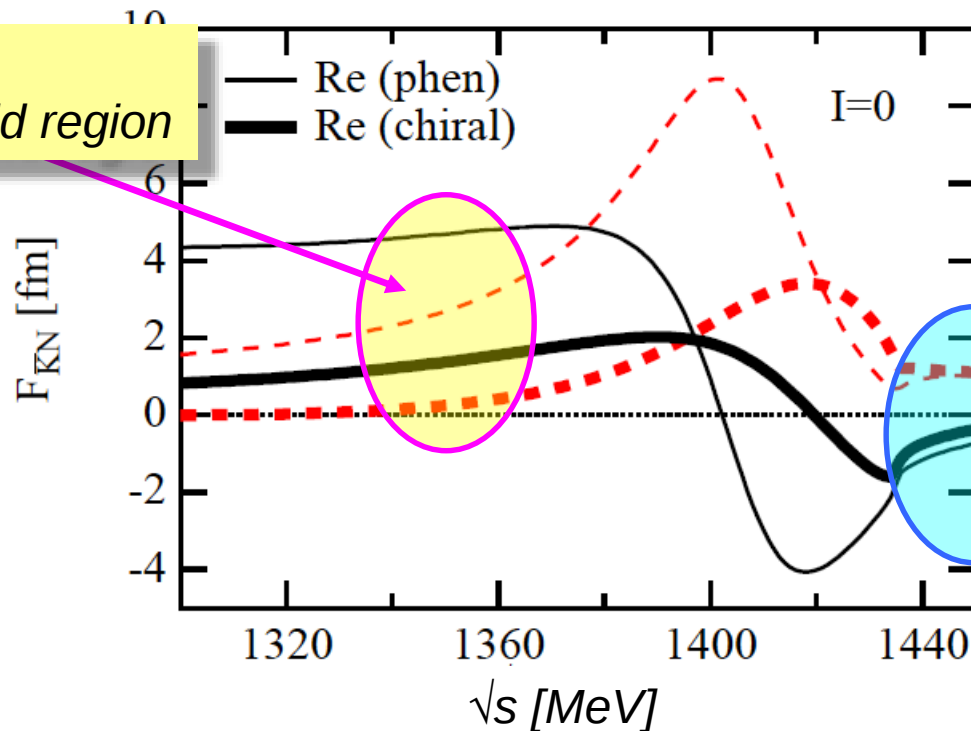
“ $K^{\text{bar}}N$ potential is shallow or deep?”

Phenomenological potential = Deep
(Akaishi-Yamazaki, E-indep.)

Chiral SU(3)-based potential = Shallow
(chiral unitary model, E-dep.)

$I=0$ $K^{\text{bar}}N$ full scattering amplitude

Very different
in sub-threshold region



i) $K^{\text{bar}}N$ interaction, $\Lambda(1405)$

“ $K^{\text{bar}}N$ potential is shallow or deep?”

Phenomenological potential = Deep
(Akaishi-Yamazaki, E-indep.)

Chiral SU(3)-based potential = Shallow
(chiral unitary model, E-dep.)

- $K^{\text{bar}}N$ scattering amplitude

- Almost the same, above $K^{\text{bar}}N$ threshold
- Resonance structure ... AY: 1405 MeV, Chiral: 1420 MeV
- AY pot. is ~4 times more attractive than Chiral pot. far below the resonance.

i) $K^{\text{bar}}N$ in

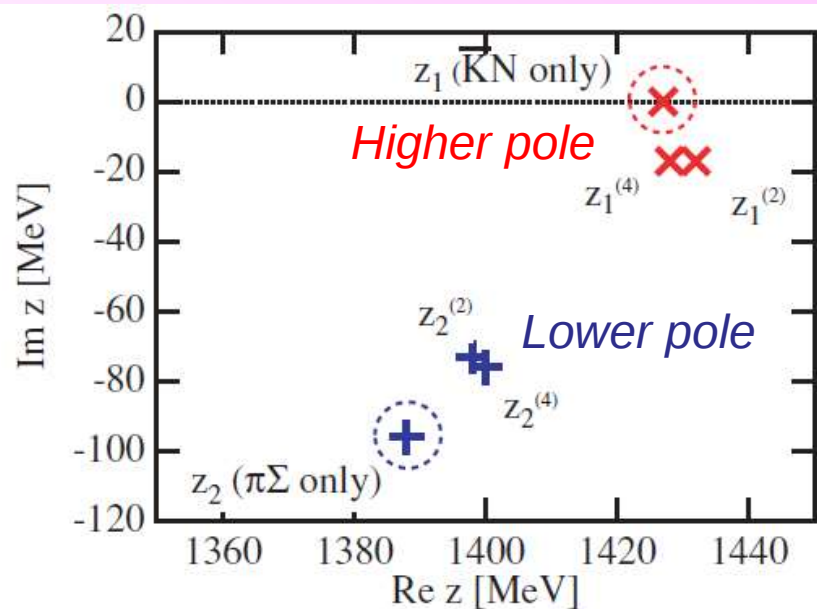
“ $K^{\text{bar}}N$ poten

Phenomenolog
(Akaishi-Yamaza

• $K^{\text{bar}}N$ scattering

- Almost
- Resona
- AY pot.

Pole positions of $I=0$ channel



T. Hyodo, W. Weise, PRC 77 (2008) 035204.

• $\Lambda(1405) = K^{\text{bar}}N$ quasi-bound state, “Building block of K^{bar} nuclei”

Double-pole nature of chiral potential

D. Jido, J.A. Oller, E. Oset, A. Ramos, U.-G. Meißner, Nucl. Phys. A 725 (2003) 181.

T. Hyodo, W. Weise, Phys. Rev. C 77 (2008) 035204.

	B and Γ	Coupling to	
Higher pole	Shallow and narrow	$K^{\text{bar}}N$	Close to $K^{\text{bar}}N$ threshold
Lower pole	Deep and broad	$\pi\Sigma$	Generated by energy dependence

= Shallow

nce.

i) $K^{\text{bar}}N$ interaction, $\Lambda(1405)$

“ $K^{\text{bar}}N$ potential is shallow or deep?”

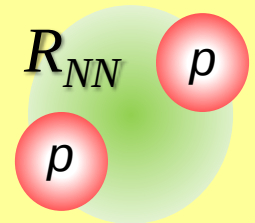
Phenomenological potential = Deep
(Akaishi-Yamazaki, E-indep.)

Chiral SU(3)-based potential = Shallow
(chiral unitary model, E-dep.)



• Difference in K^-pp

	B. E. [MeV]	Γ [MeV]	R_{NN} [fm]	“Density”
AY	49	60	1.8	$\sim 2 \rho_0$
Chiral	17	47	2.2	ρ_0



i) $K^{\text{bar}}N$ interaction, $\Lambda(1405)$

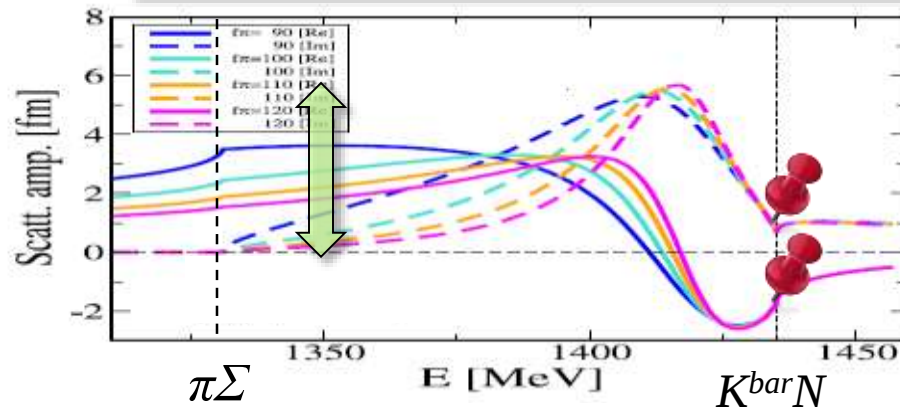
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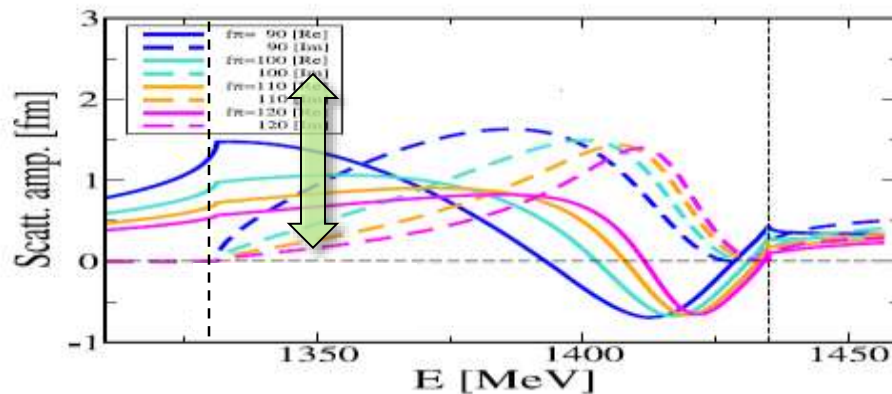
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Uncertainty in sub-threshold region

$K^{\text{bar}}N$



$\pi\Sigma$



i) $K^{\text{bar}}N$ interaction, $\Lambda(1405)$

“ $K^{\text{bar}}N$ potential is shallow or deep?”

Phenomenological potential = Deep
(Akaishi-Yamazaki, E-indep.)

Cl
(cl)

Precise measurement of Kaonic atoms

- Kaonic Hydrogen: KpX, DEAR, SIDDHATA
- Kaonic ^4He : KEK (E570), J-PARC (E17)
- Kaonic Deuteron, ^3He : SIDDHATA2

To fix uncertainty in sub-threshold region...

$K^{\text{bar}}N$

$\Lambda(1405)$ study with $d(K^-, n)$

J-PARC (E31) will be done

Scattering
data

$\pi\Sigma$ scattering length
from Λ_c decay

T. Hyodo, M. Oka,
PRC84, 035201 (2011)

$\pi\Sigma$ mass spectrum
 $\pi\Sigma^+$, $\pi^+\Sigma^-$, $\pi^0\Sigma^0$

Precise $\pi\Sigma$ line shape

CLAS/Jlab, LEPS/SPring8

$\pi\Sigma$ threshold

$K^{\text{bar}}N$ threshold

$\sqrt{s}$

$\sqrt{s}$

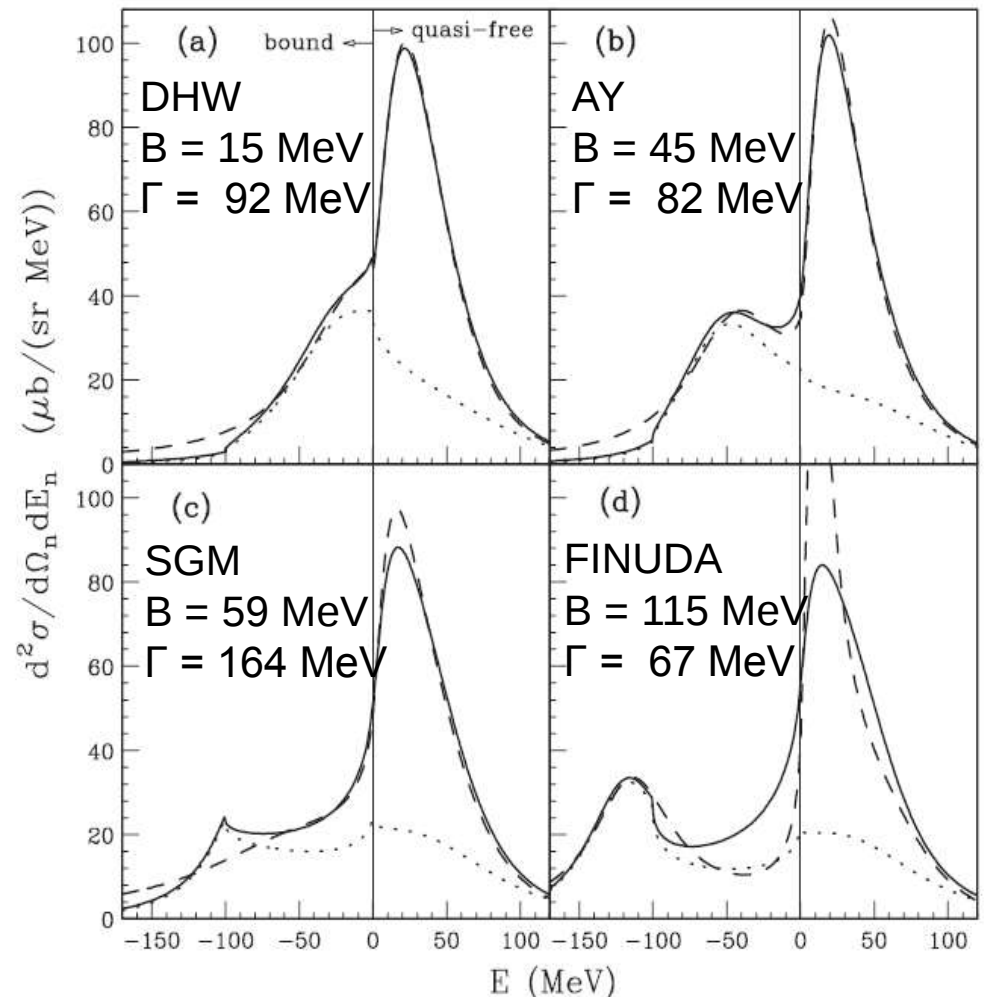
ii) Comparison between theory and experiment

Peak structure
observed in exp.

- Large decay width
- $\pi\Sigma N$ threshold

Pole position
calculated in theor.

Inclusive spectra of $^3\text{He}(K^-, \text{inflight } n)$



T. Koike and T. Harada, PRC80, 055208 (2009)

ii) Comparison between theory and experiment



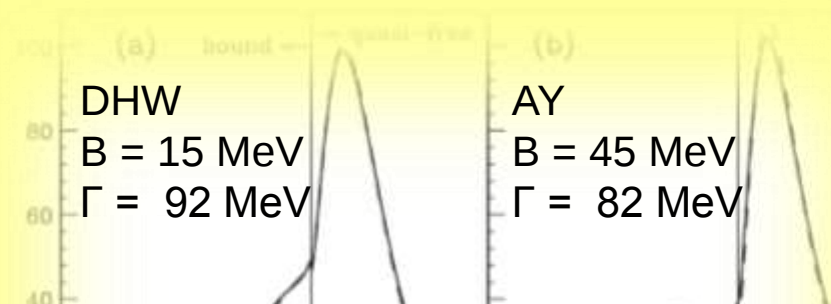
Peak structure
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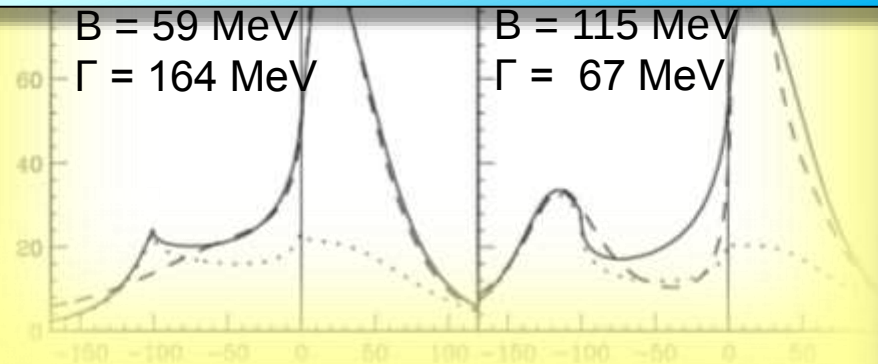
Pole position
calculated in theor.



Inclusive spectra of $^3\text{He}(K^-, \text{inflight } n)$



Need a bridge!
... Reaction calculation
in theoretical study



iii) Relevant degree of freedom

K^- abs. on Li, C

$p + p$

$^3\text{He} (K^-, n)$

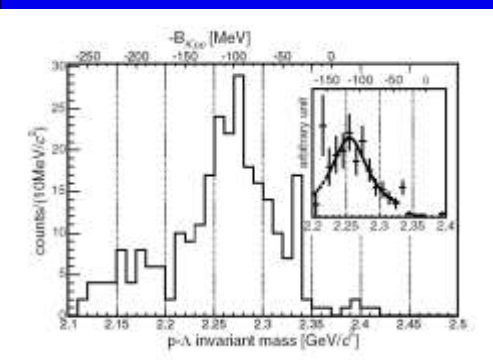
$d (\pi^+, K^+)$

K^-pp

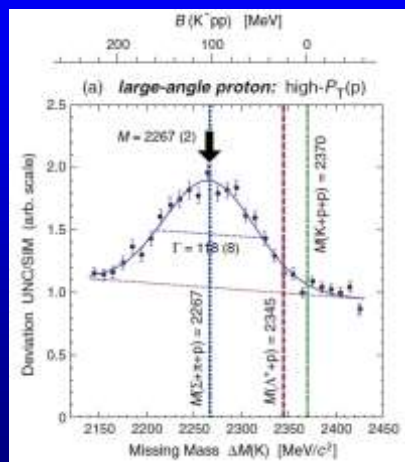
p

Λ

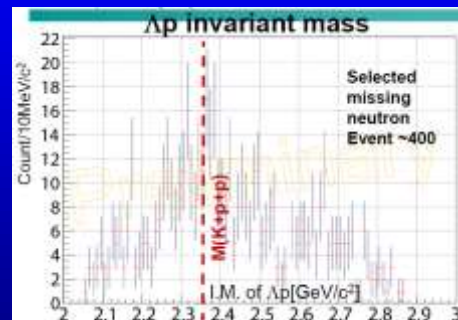
Invariant mass



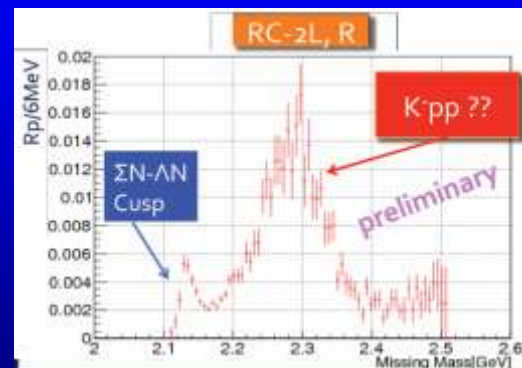
FINUDA



DISTO



J-PARC E15



J-PARC E27

iii) Relevant degree of freedom



K^- abs. on Li, C

$p + p$

**Something with
 $B=2, S=-1$, charge $+1$**

p

$^3\text{He} (K^-, n)$

$d (\pi^+, K^+)$

$K^-pp?$ $\pi\Sigma N?$
 $\Lambda^*N??$ $\Sigma^*N??$
 $qqq-qqq-q^{\text{bar}} s ???$

Λ

giant mass



- Measure various quantities (spin, parity, size, ...)
- Systematic study of various systems

5. Summary

5. Summary

Kaonic nuclei (nuclear system with anti-kaon)

... Exotic system due to strong $K^{\text{bar}}N$ attraction?

Doorway to dense matter?

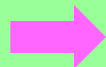
K^-pp = “Prototype of K^{bar} nuclei”

Lots of theoretical studies predict K^-pp can be bound below $K^{\text{bar}} + 2N$ threshold.

Variational calc.
Faddeev-AGS

×

Chiral pot.
Phen. pot.



K^-pp exists as a **resonance** between $K^{\text{bar}}NN$ and $\pi\Sigma N$ thresholds.
Binding energy < 100 MeV, Large decay width (>40 MeV)

K^-pp studied with ccCSM+Feshbach method

Effective single-channel potential derived in ccCSM

Chiral SU(3)-based pot. (Gauss-type)

Self-consistency for kaon's **complex** energy

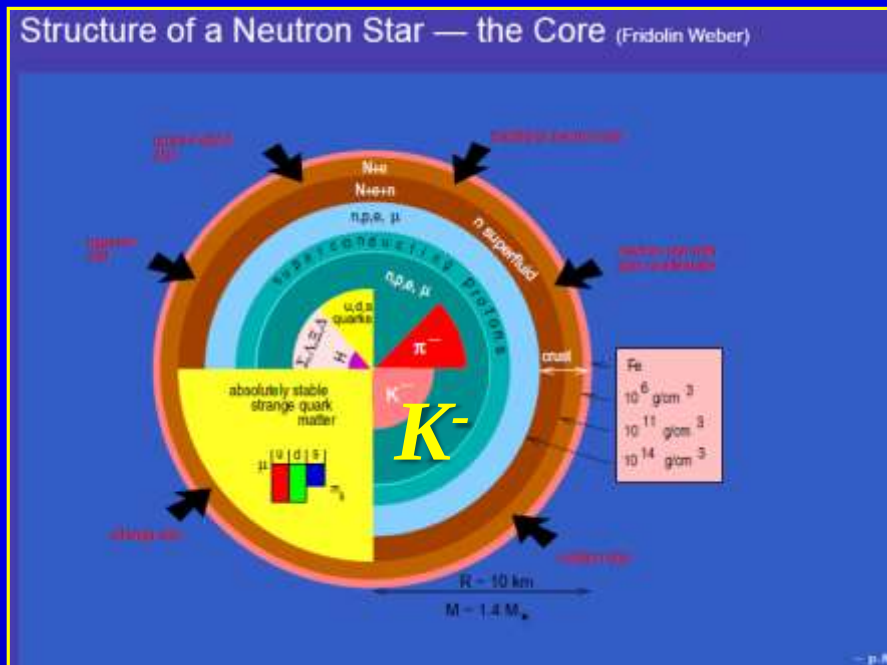


$(B, \Gamma/2) = (26, 12) \text{ MeV} : \text{Ansatz 1}$
 $(27, 19) \text{ MeV} : \text{Ansatz 2}$
Mean NN distance $\sim 2.2 \text{ fm}$

Issues on kaonic nuclei

- $K^{\text{bar}}N$ interaction “Deep or shallow?”, “Double pole of $\Lambda(1405)$ ” by chiral potential
... Solve uncertainty in $K^{\text{bar}}N$ sub-threshold region
- Observed spectra and calculated pole position
... Large decay width, threshold effect $\Rightarrow$ Help theoretical analysis of reaction study
- What is observed in invariant-mass measurement? K^-pp ? $\pi\Sigma N$? Λ^*N ? Σ^*N ? $7q-q^{\text{bar}}$ s
 $\Rightarrow$ Measurement of various quantities (spin-parity, size ...), Systematic study

- *What's the role of anti-kaon in neutron star?*
- *Compare with charm sector...*



Analogy of strangeness and charm?

$\Lambda(1405) 1/2^-$

$$I(J^P) = 0(\frac{1}{2}^-)$$

Mass $m = 1405.1^{+1.3}_{-1.0}$ MeV

Full width $\Gamma = 50 \pm 2$ MeV

Below $\bar{K}N$ threshold

Strangeness

uds

$\Lambda(1405)$ DECAY MODES

Fraction (Γ_i/Γ)

p (MeV/c)

$\Sigma \pi$

100 %

155

$\Lambda_c(2595)^+$

$$I(J^P) = 0(\frac{1}{2}^+)$$

The spin-parity follows from the fact that $\Sigma_c(2455)\pi$ decays, with little available phase space, are dominant. This assumes that $J^P = 1/2^+$ for the $\Sigma_c(2455)$.

Mass $m = 2592.25 \pm 0.28$ MeV

$m - m_{\Lambda^+} = 305.79 \pm 0.24$ MeV

Full width $\Gamma = 2.6 \pm 0.6$ MeV

Charm

udc

Thank you for your attention!

