

# Lattice QCD analysis of partial restoration of chiral symmetry in flux-tube structure between quarks

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KEK

ヘビーコークハドロンと原子核のスペクトルと構造

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Refs. T. I., G. Cossu, and S. Hashimoto,  
PoS (LATTICE 2013) 376, arXiv:1311.0218 [hep-lat];  
PoS (Hadron 2013) 159, arXiv:1401.4293 [hep-lat].

1 Interquark Potential and Flux-tube Structure

2 Chiral Condensate in Flux-tube

- Chiral Condensate in Quark-Antiquark System
- Chiral Condensate in 3-Quark System
- Chiral Condensate in Quark-Antiquark + Quark-Antiquark System

3 Summary

## 1 Interquark Potential and Flux-tube Structure

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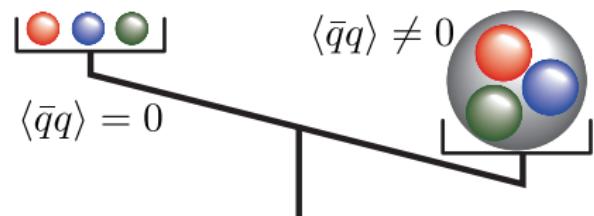
## 3 Summary

# Non-perturbative QCD Phenomena

$$\mathcal{L}_{\text{QCD}} = -\frac{1}{2} \text{Tr } G_{\mu\nu}G^{\mu\nu} + \bar{q}(i\cancel{D} - m)q$$

## Chiral Symmetry Breaking

$$\text{SU}(N_f)_L \times \text{SU}(N_f)_R \rightarrow \text{SU}(N_f)_V$$



Chiral symmetry is restored

- High temperature
  - Quark Gluon Plasma
- Finite density
  - neutron star
  - nuclear matter (partially ?)

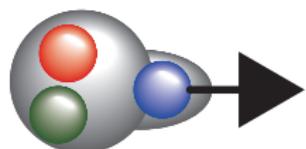
## Confinement

There are no isolated quarks.

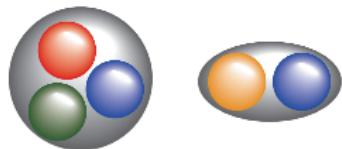
1.



2.



3.



# Quarkonia Spectroscopy and Interquark Potential

## Charmonium spectra

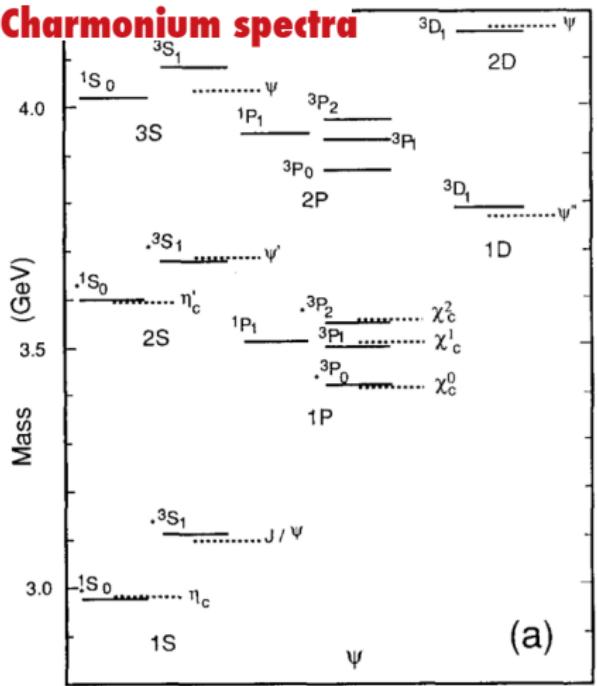
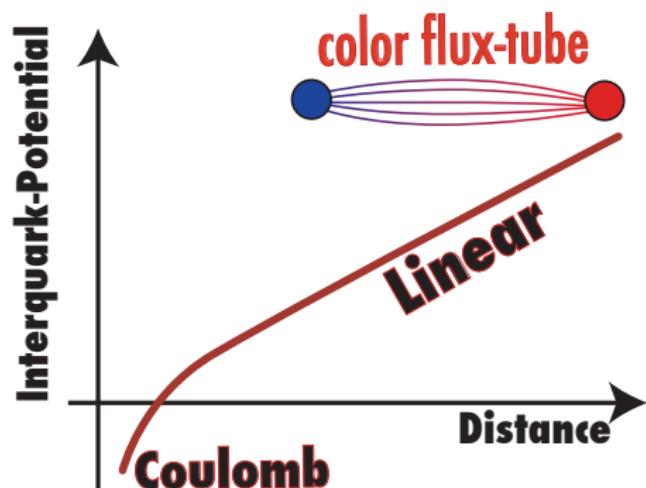


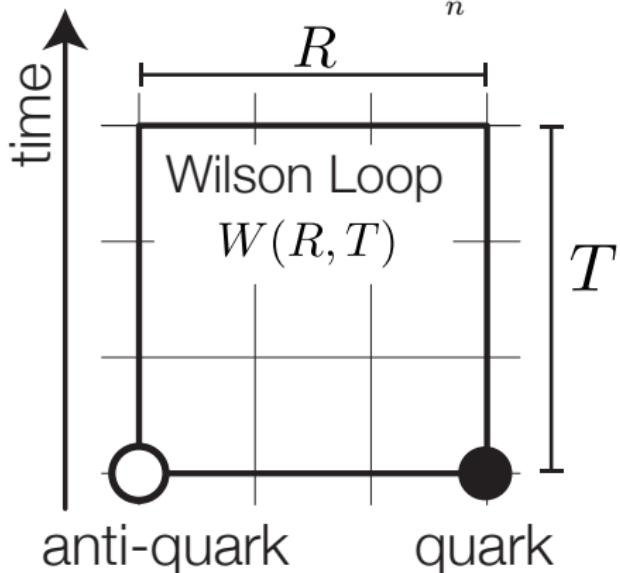
Fig. Mukherjee, et al., '93

■ Coulomb + linear potential

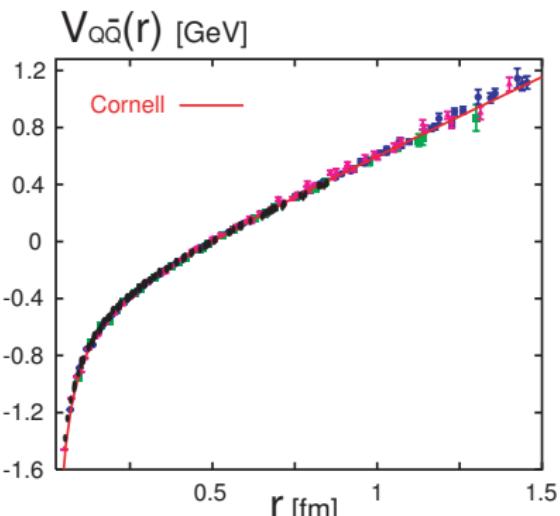


# Wilson Loop and Interquark Potential in Lattice QCD

$$\begin{aligned} W(R, T) &\equiv \text{Tr} \prod_{\text{loop}} U_\mu \\ &= \langle Q\bar{Q} | e^{-HT} | Q\bar{Q} \rangle \\ &= \sum_n C_n e^{-E_n(R)T} \sim \exp(-V_{Q\bar{Q}}(R)T) \end{aligned}$$



Coulomb + Linear potential



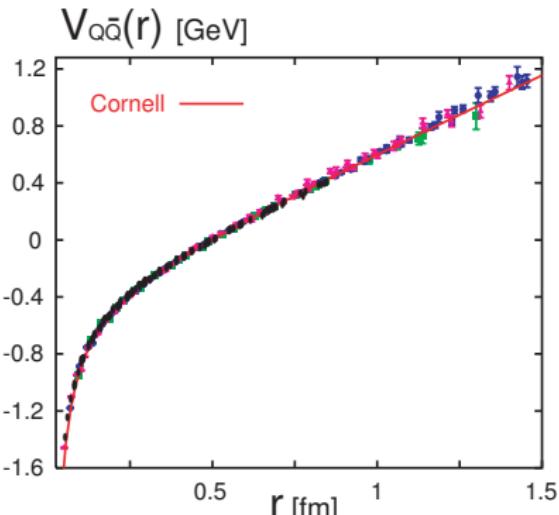
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linear potential  
color flux-tube



Coulomb + Linear potential



# Flux-tube Measurement in Lattice QCD

■ action density  $\rho(x)$  around  
Wilson loop, i.e.,  $m_Q \rightarrow \infty$   
(static) Quark-Antiquark pair

$$\langle \rho(x) \rangle_W \equiv \frac{\langle \rho(x) W(R, T) \rangle}{\langle W(R, T) \rangle} - \langle \rho \rangle$$

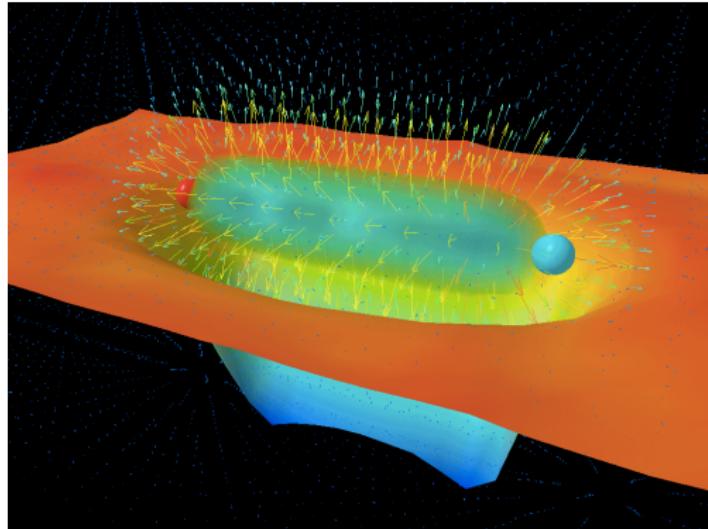
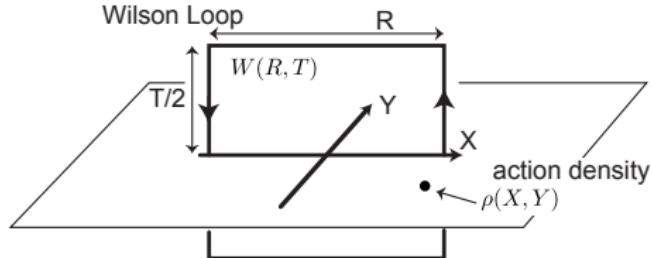
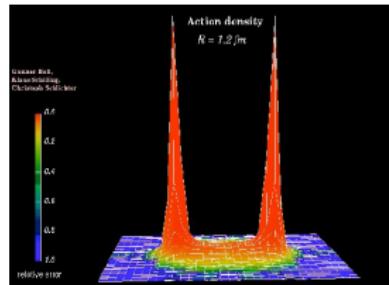


Fig. Leinweber et al. '03

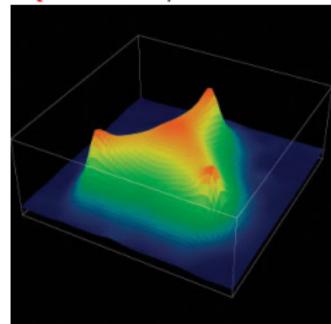
# Flux-tube Structures for Multi-Quark Systems

interquark potential is proportional to **the flux-tube length**

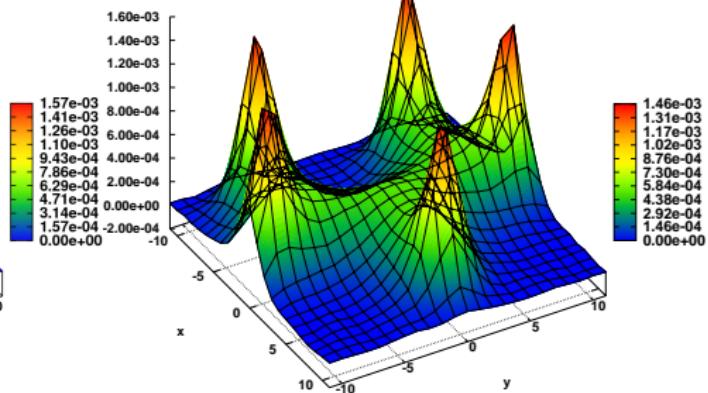
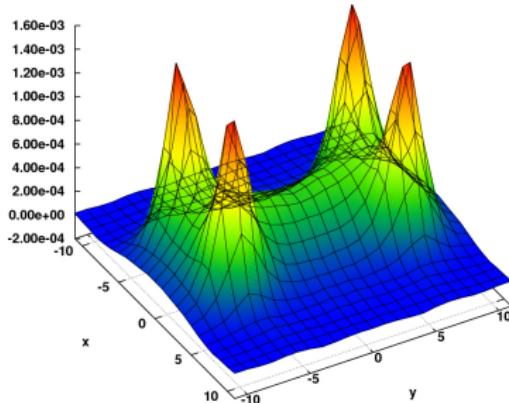
QQ-bar : Bali-Schlichter-Schilling '95



3Q : Ichie, et al. '03

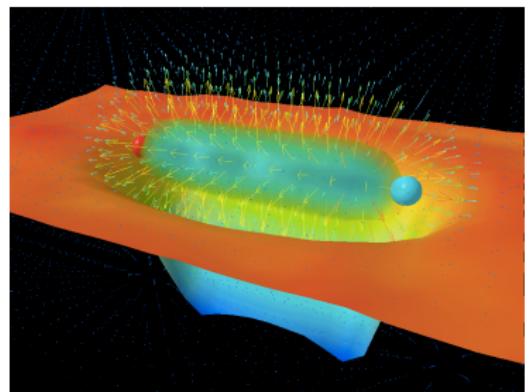
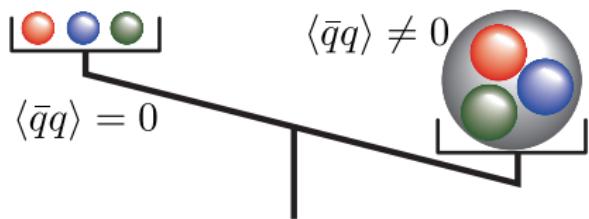


4Q and 5Q : Bicudo-Cardoso-Cardoso '11



# Question

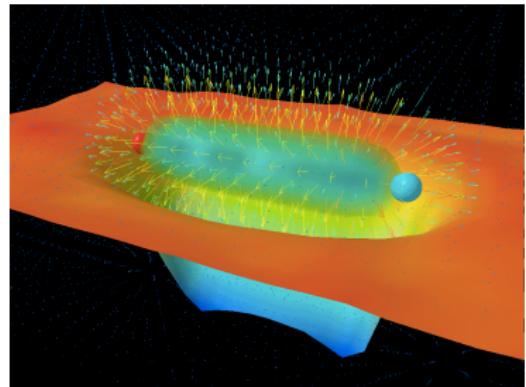
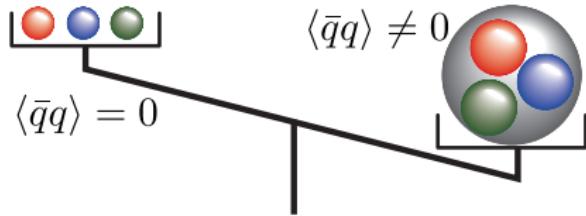
Chiral Symmetry Breaking  $\times$  Quark Confinement = ?



# Question

Chiral Symmetry Breaking  $\times$  Quark Confinement = ?

- Quarks are confined in **hadrons** by color flux-tube.
- How about “**chiral symmetry**” in color flux-tube, i.e., “hadron” ?
- From **lattice QCD**, we analyze **chiral condensate** inside color flux-tube.



## ① Interquark Potential and Flux-tube Structure

## ② Chiral Condensate in Flux-tube

- Chiral Condensate in Quark-Antiquark System
- Chiral Condensate in 3-Quark System
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③ Summary

# Chiral Condensate and Dirac Eigenmode

- QCD Lagrangian

$$\mathcal{L}_{\text{QCD}} = -\frac{1}{2} \text{Tr } G_{\mu\nu}G_{\mu\nu} + \bar{q}(\not{D} + m)q$$

- chiral condensate  $\langle \bar{q}q \rangle \Leftarrow \text{eigenvalue of Dirac operator } \not{D}\psi_\lambda = \lambda\psi_\lambda$

$$-\langle \bar{q}q \rangle = \langle \text{Tr } S_q \rangle = \lim_{m \rightarrow 0} \lim_{V \rightarrow \infty} \langle \text{Tr} \frac{1}{\not{D} + m} \rangle = \lim_{m \rightarrow 0} \lim_{V \rightarrow \infty} \frac{1}{V} \sum_{\lambda} \frac{1}{\lambda + m}$$

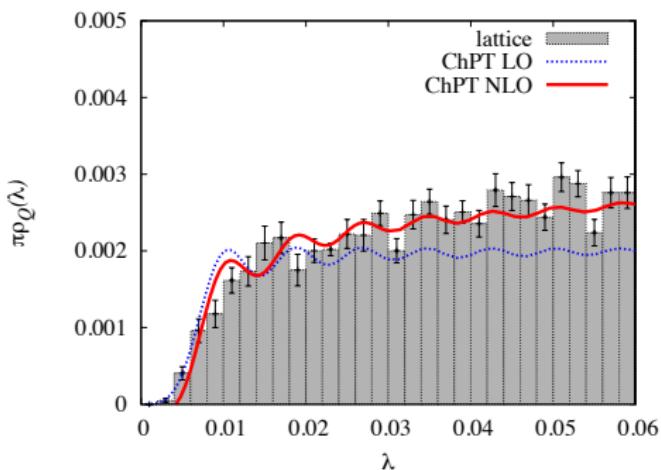
## Banks–Casher Relation

$$\langle \bar{q}q \rangle = -\pi \langle \rho(0) \rangle$$

near-zero Dirac mode

⇒ chiral symmetry breaking

$\rho(\lambda)$ : eigenmode density



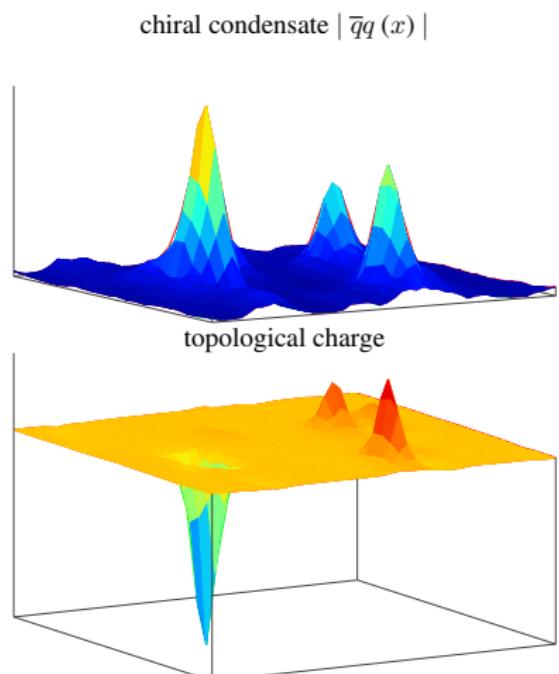
# Local Structure of Chiral Condensate in QCD Vacuum

*a snapshot of QCD vacuum*

$$\begin{aligned}\langle \bar{q}q \rangle &= -\text{Tr} \frac{1}{\not{D} + m} = -\frac{1}{V} \sum_{\lambda} \frac{1}{\lambda + m} \\ &= -\frac{1}{V} \sum_{\lambda} \sum_x \frac{\psi_{\lambda}^{\dagger}(x) \psi_{\lambda}(x)}{\lambda + m} \\ &= \frac{1}{V} \sum_x \bar{q}q(x)\end{aligned}$$

with Dirac eigenfunction  $\not{D}\psi_{\lambda} = \lambda\psi_{\lambda}$

- a local chiral condensate  $\bar{q}q(x)$  forms clusters  $\Rightarrow$  instantons



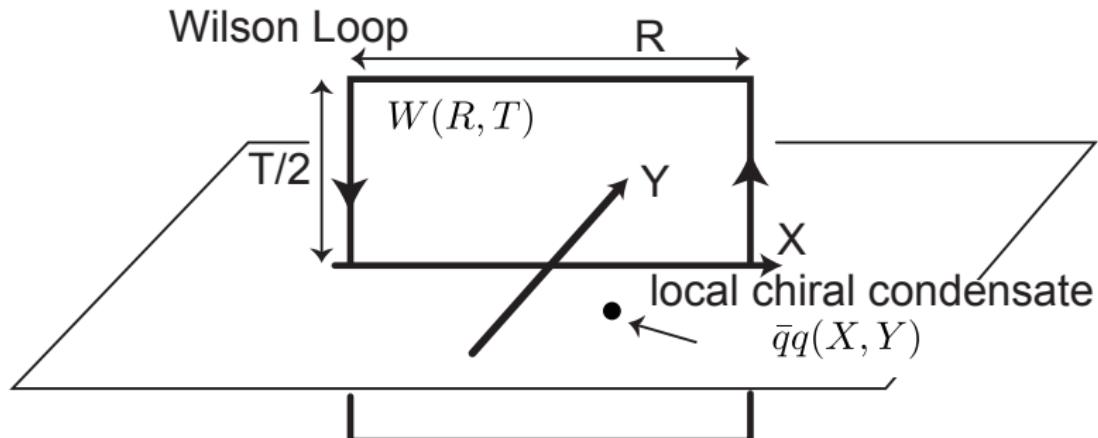
# Chiral Condensate inside Color Flux-tube

local chiral condensate is given by

$$\bar{q}q(x) = - \sum_{\lambda} \frac{\psi_{\lambda}^{\dagger}(x)\psi_{\lambda}(x)}{\lambda + m} \quad \text{with Dirac eigenmode } D\psi_{\lambda} = \lambda\psi_{\lambda}$$

change of the chiral condensate around the Wilson loop

$$\langle \bar{q}q(x) \rangle_W \equiv \frac{\langle \bar{q}q(x)W(R, T) \rangle}{\langle W(R, T) \rangle} - \langle \bar{q}q \rangle$$



## About Lattice QCD Setup

In this study, we mainly use

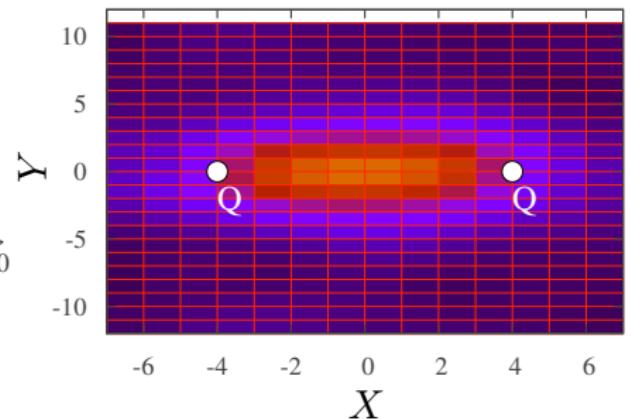
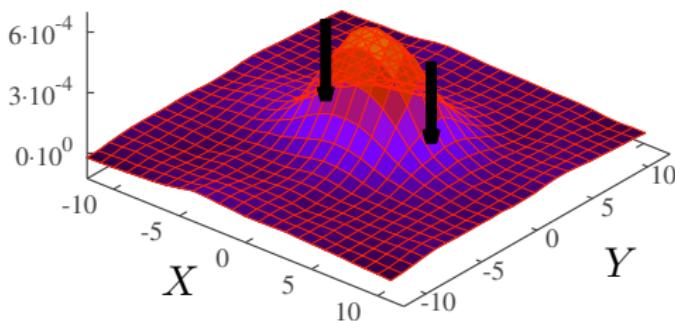
- 2+1 flavor dynamical **overlap-fermion** configuration by JLQCD Coll.
  - ⇒ overlap-fermion keeps “**exact chiral symmetry**” on lattice
  - ⇒ overlap-fermion Dirac eigenmodes are ideal probe  
to analyze chiral properties in lattice QCD
- pion mass  $m_\pi \sim 300$  MeV, kaon mass  $m_K \sim 500$  MeV
- lattice volume  $24^3 \times 48$
- lattice spacing  $a^{-1} = 1.759(10)$  GeV  $\Rightarrow a \sim 0.112$  fm

# Reduction of Chiral Condensate in the Flux-tube

difference of (truncated) chiral condensate with quark-antiquark

$$\langle \bar{q}q(x) \rangle_W^{(N)} \equiv \frac{\langle \bar{q}q^{(N)}(x)W(R, T) \rangle}{\langle W(R, T) \rangle} - \langle \bar{q}q \rangle^{(N)} \quad (N : \text{eigenmode number})$$

- $\langle \bar{q}q(x) \rangle_W > 0 \Rightarrow |\langle \bar{q}q(x) \rangle|_{\text{flux}} < |\langle \bar{q}q \rangle| \quad \because \langle \bar{q}q \rangle < 0 \text{ in vacuum}$
- partial restoration of chiral symmetry in the flux ("Bag-model"-like)

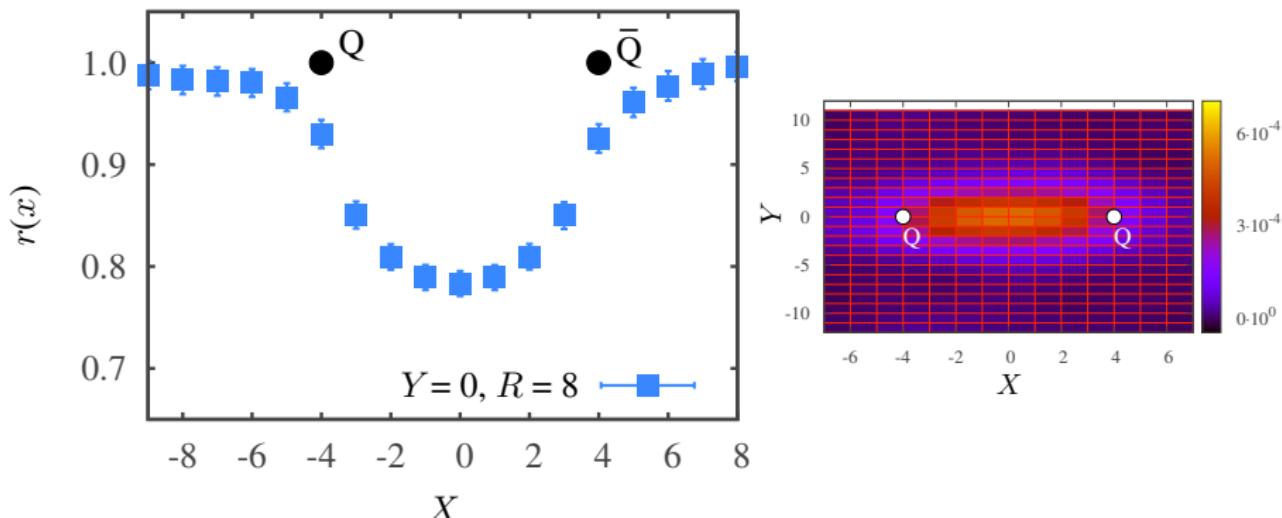


(anti-)quark at  $(4, 0)$  and  $(-4, 0)$ . lattice unit  $a \sim 0.11$  fm, pion mass  $\sim 300$  MeV

# Ratio of Chiral Condensate

$$r(x) \equiv \frac{\langle \bar{q}q^{(\text{subt})}(x)W(R,T) \rangle}{\langle \bar{q}q^{(\text{subt})} \rangle \langle W(R,T) \rangle} < 1$$

⇒ about 20 % of chiral condensate is reduced at flux core  
due to strong color-electric field between quarks  
cross section of flux-tube



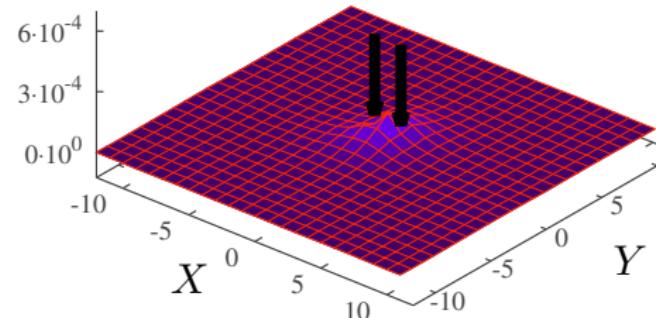
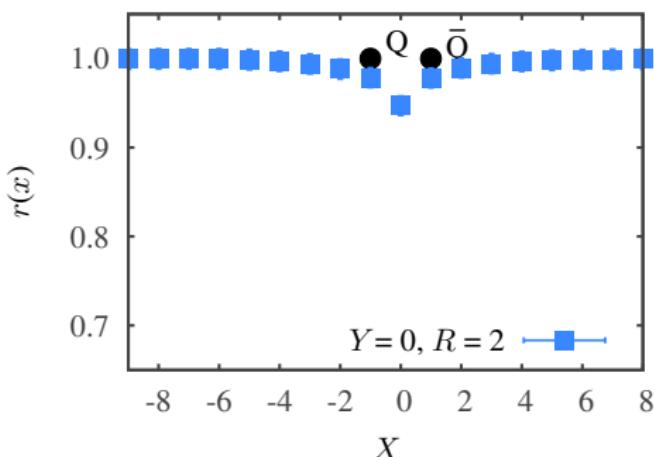
(taking continuum limit  $a \rightarrow 0$ ) subtracted condensate

$$\langle \bar{q}q \rangle^{(N)} = \langle \bar{q}q^{(\text{subt})} \rangle + c_1^{(N)} m_q/a^2 + c_2^{(N)} m_q^3$$

ref. JLQCD Coll. Noaki '09

# Quark Separation Dependence of Restoration

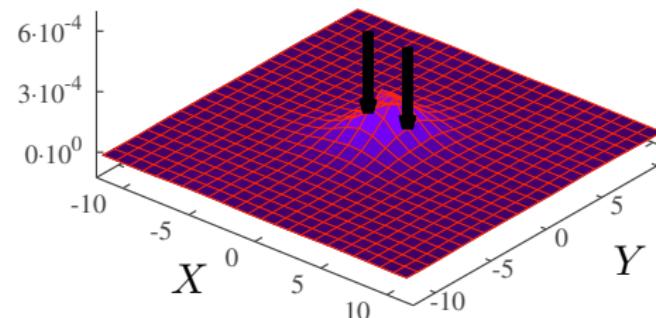
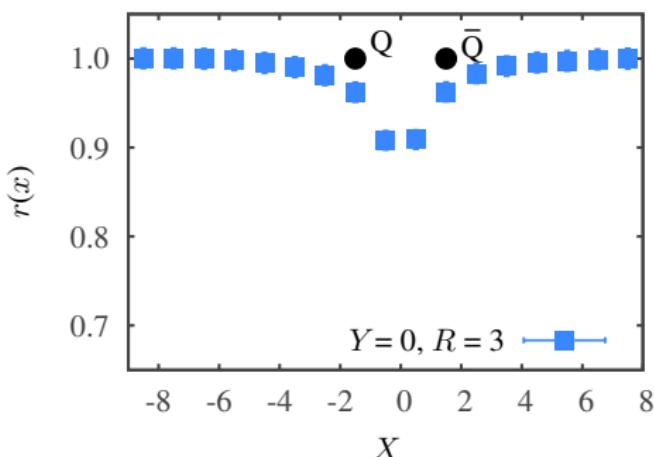
partial restoration **increases** with the interquark separation  $R$ ,  
until the string breaking occurs around  $R \sim 1$  fm



cf. lattice unit  $a \sim 0.11$  fm

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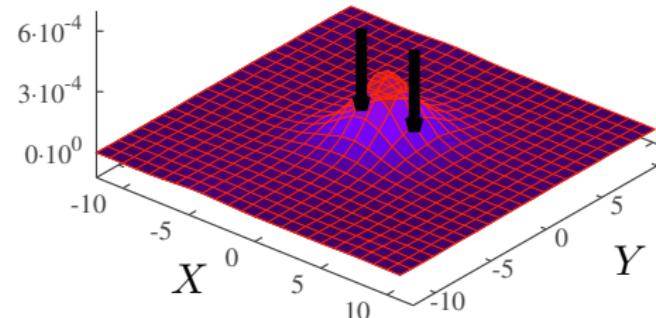
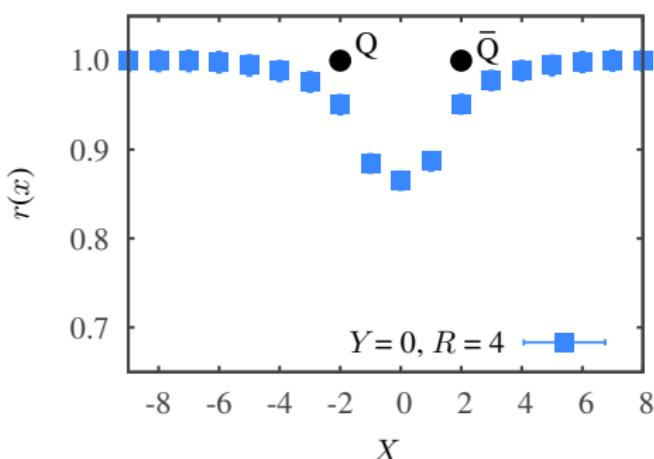
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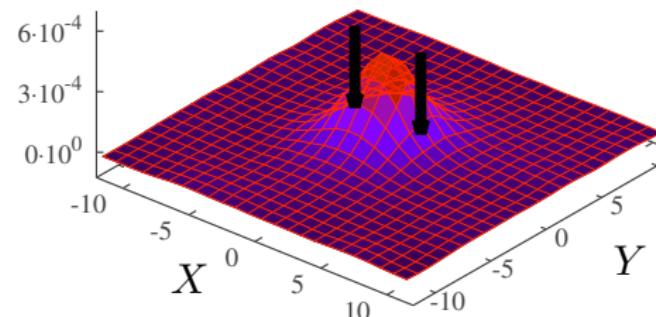
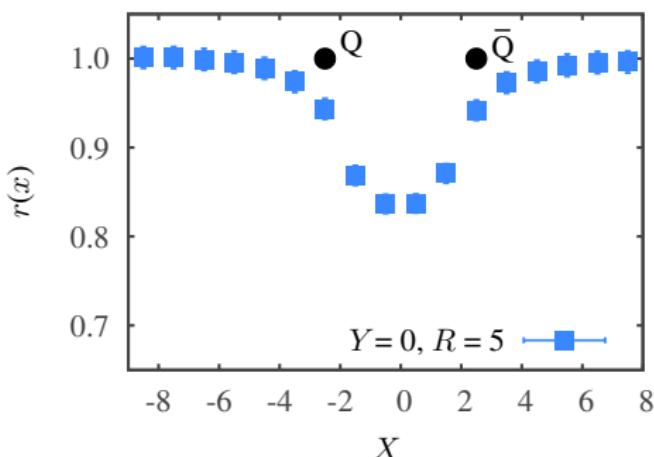
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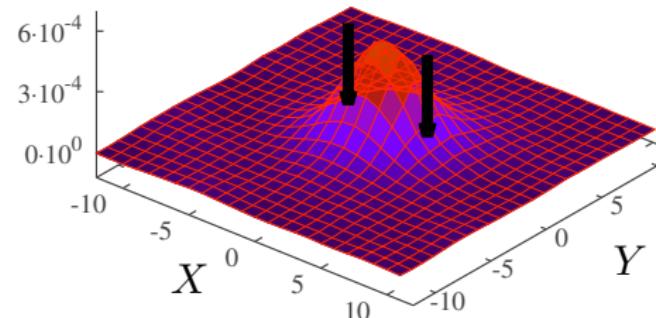
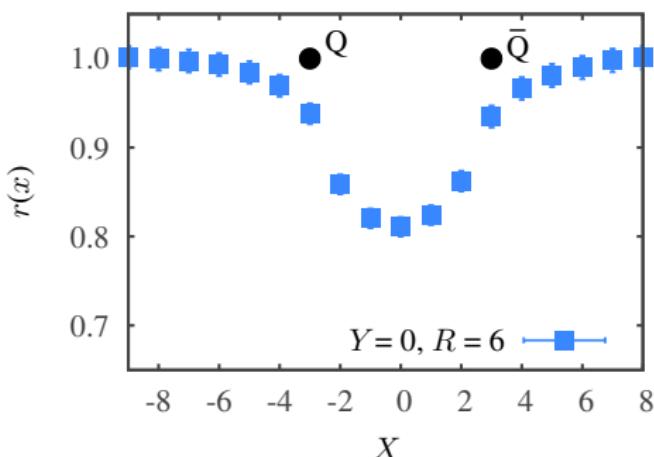
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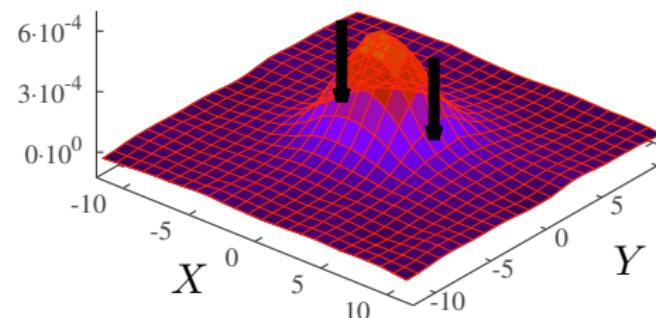
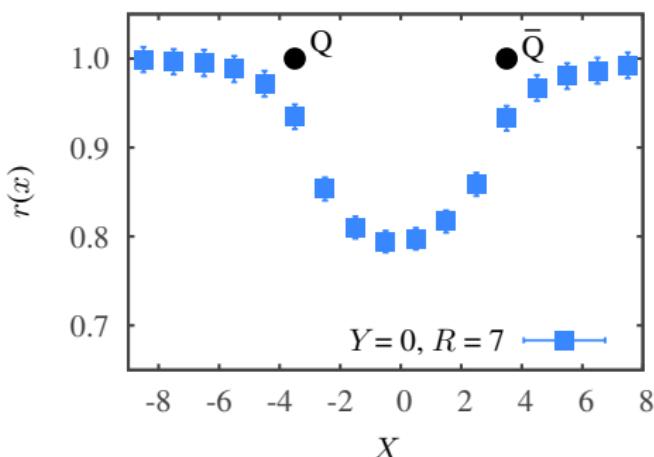
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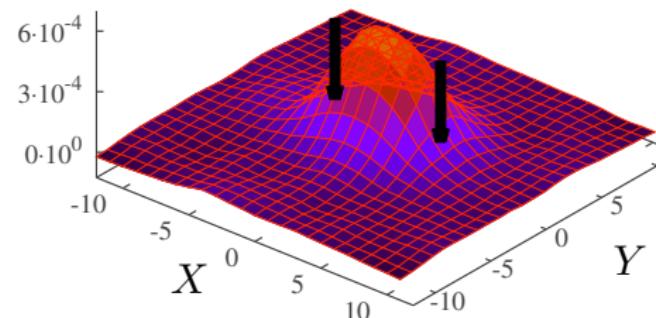
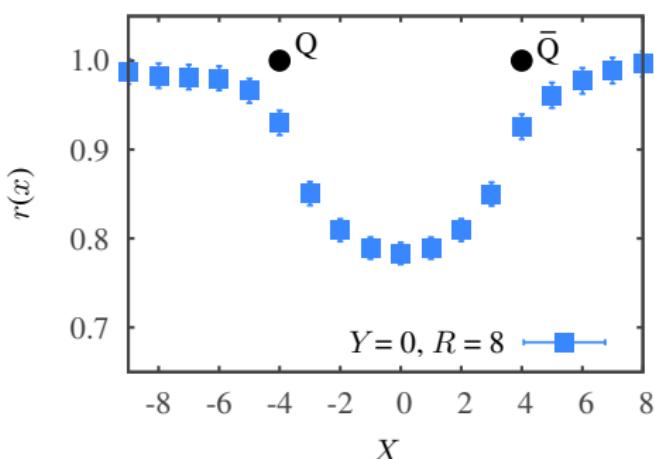
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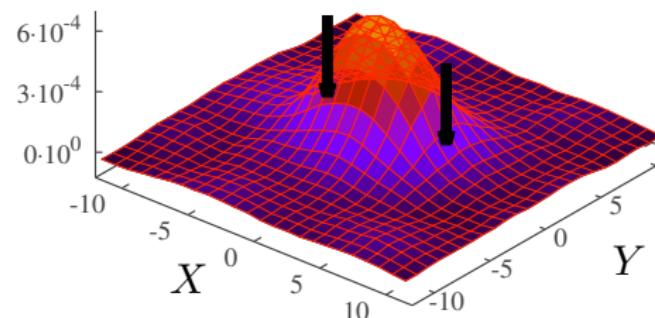
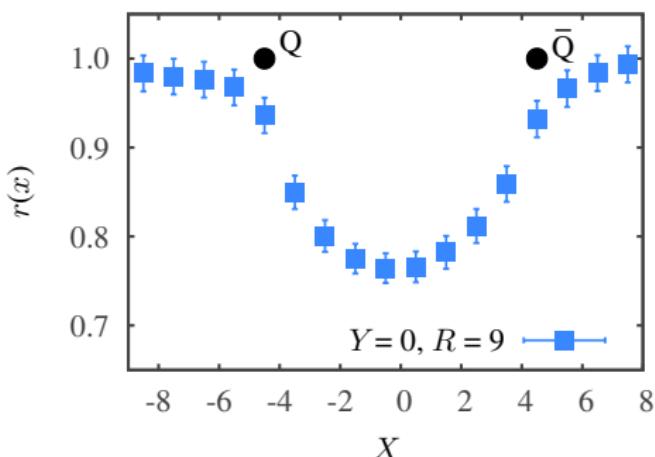
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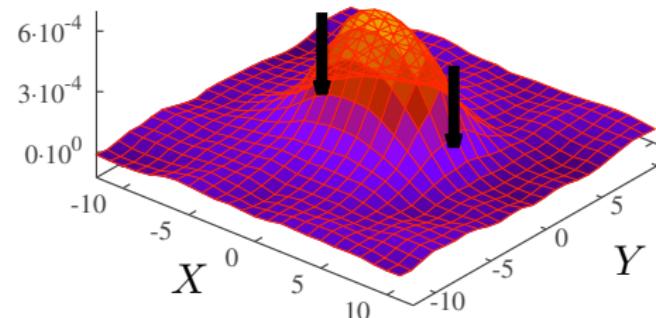
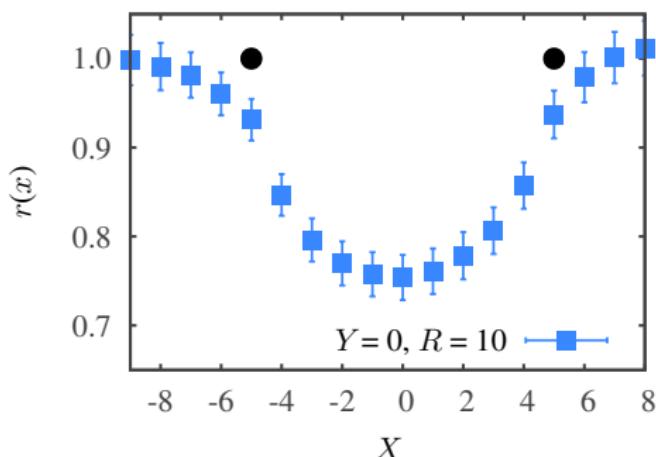
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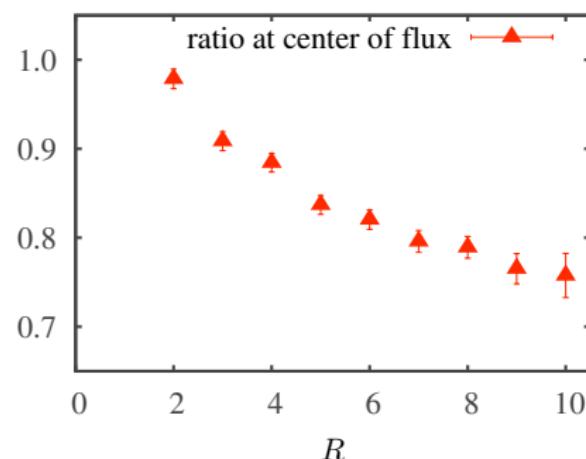
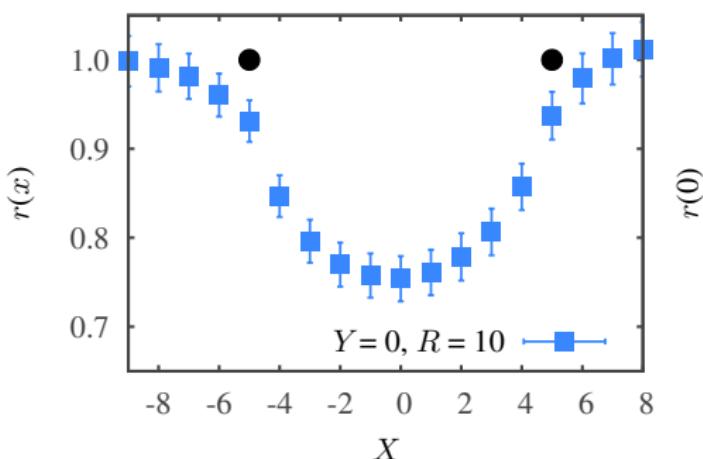
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## ② Chiral Condensate in Flux-tube

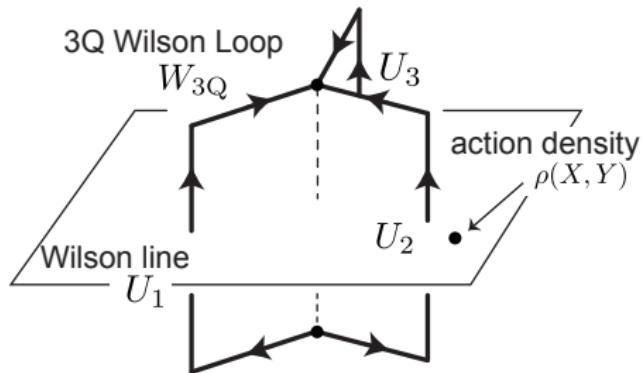
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- **Chiral Condensate in 3-Quark System**
- Chiral Condensate in Quark-Antiquark + Quark-Antiquark System

## ③ Summary

# 3-Quark System

3Q-Wilson loop

$$W_{3Q} \equiv \frac{1}{3!} \varepsilon_{abc} \varepsilon_{a'b'c'} U_1^{aa'} U_2^{bb'} U_3^{cc'}$$



with Wilson lines  $U_k$   
cf. Takahashi-Suganuma '01

Y-type flux formation in 3Q-system

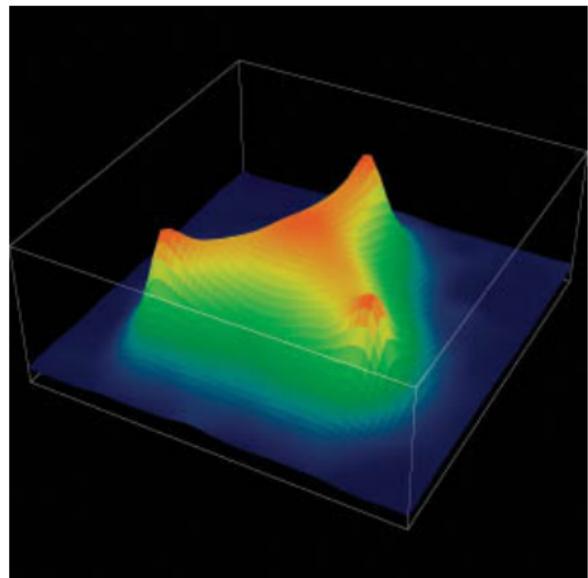


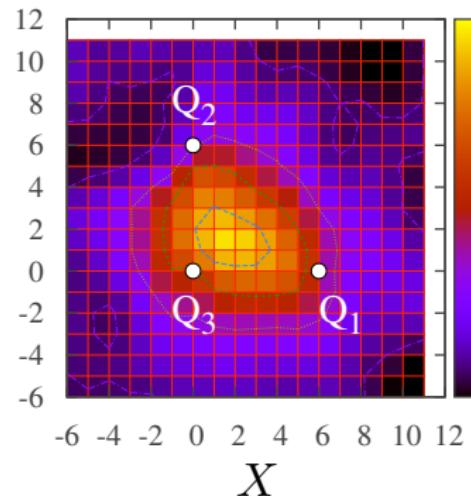
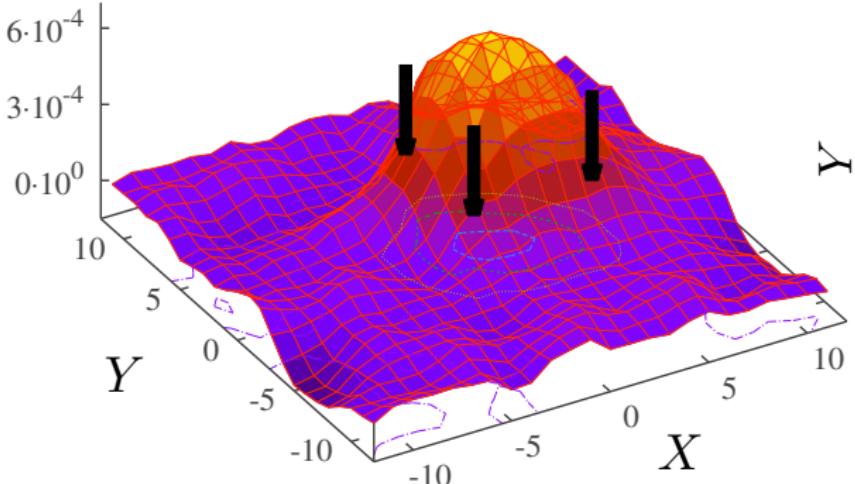
Figure : Ichie et al. '03

# Partial Restoration of Chiral Symmetry in 3Q system

■ chiral condensate in 3Q system

$$\langle \bar{q}q(x) \rangle_{W_{3Q}} \equiv \frac{\langle \bar{q}q(x) W_{3Q} \rangle}{\langle W_{3Q} \rangle} - \langle \bar{q}q \rangle > 0$$

○ Y-type flux is smeared due to small quark separation



right triangle configuration :  $Q_1 = (6, 0), Q_2 = (0, 6), Q_3 = (0, 0)$

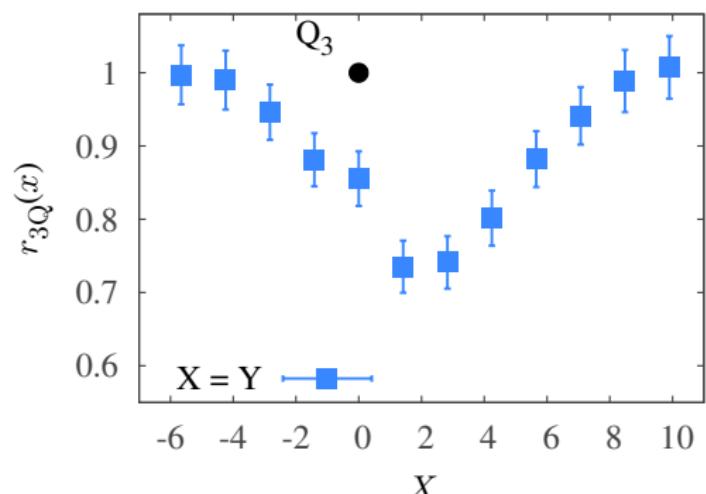
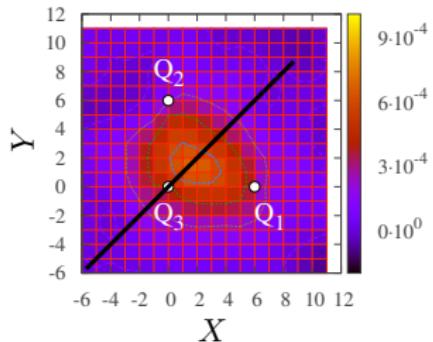
lattice unit  $a \simeq 0.11$  fm

# Ratio of Chiral Condensate in 3Q system

$$r_{3Q}(x) \equiv \frac{\langle \bar{q}q^{(\text{subt})}(x)W_{3Q} \rangle}{\langle \bar{q}q^{(\text{subt})} \rangle \langle W_{3Q} \rangle} < 1 \quad \text{inside flux}$$

- about  $20 \sim 30\%$  of chiral condensate is reduced in 3Q-system
- “partial restoration of chiral symmetry” in “Baryon”

cross-section of 3Q-flux



## ① Interquark Potential and Flux-tube Structure

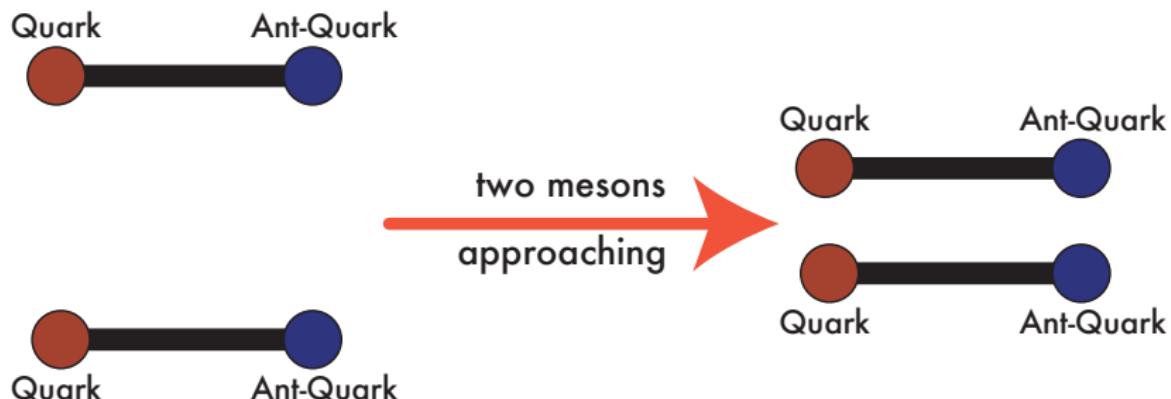
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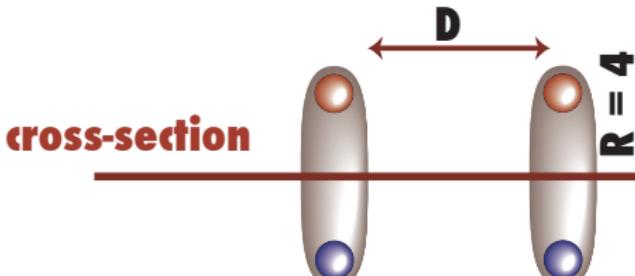
## ③ Summary

# Meson-Meson System

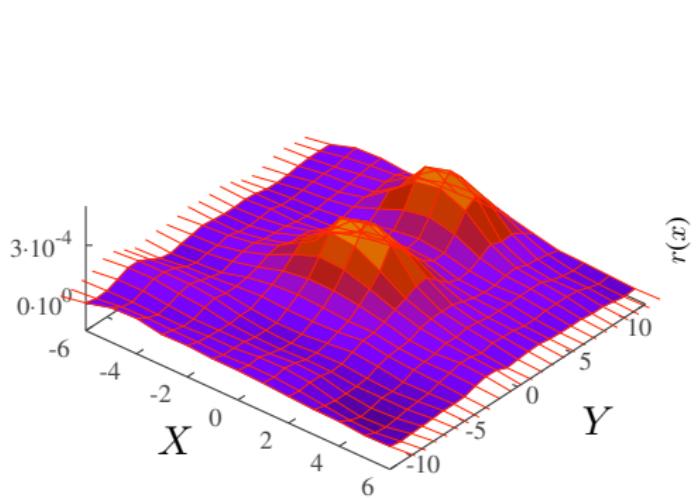
Considering two mesons are approaching,  
how about chiral symmetry restoration ?  
enhanced or not ?



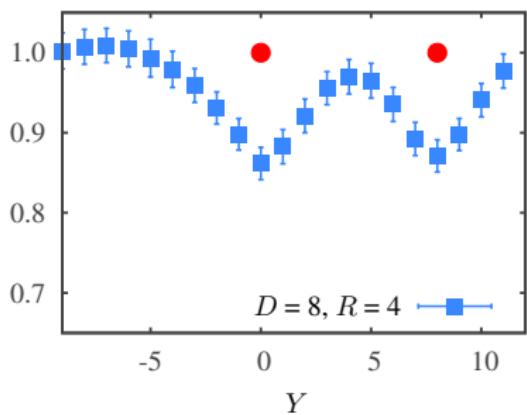
# Quark-Antiquark Pair



cross-section

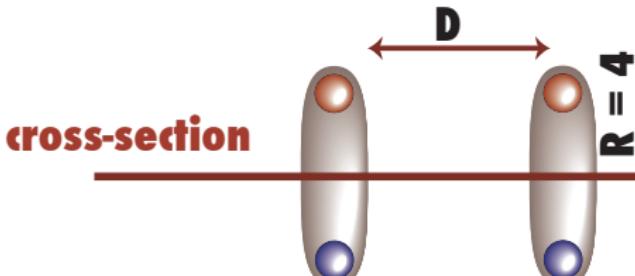


ratio of chiral condensate

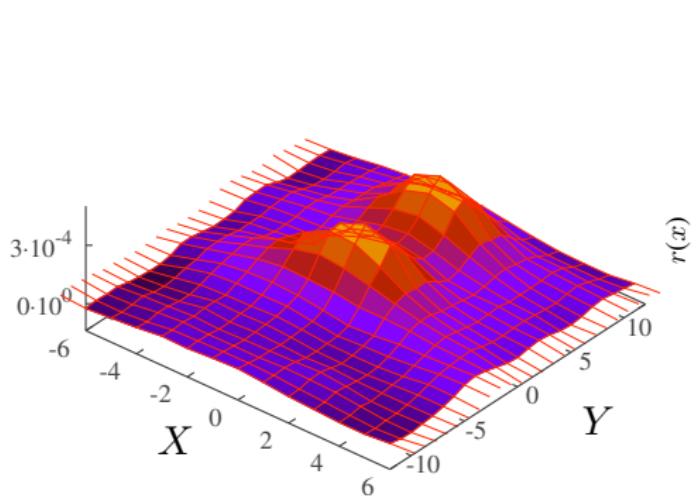


cf. lattice unit  $a \sim 0.11$  fm

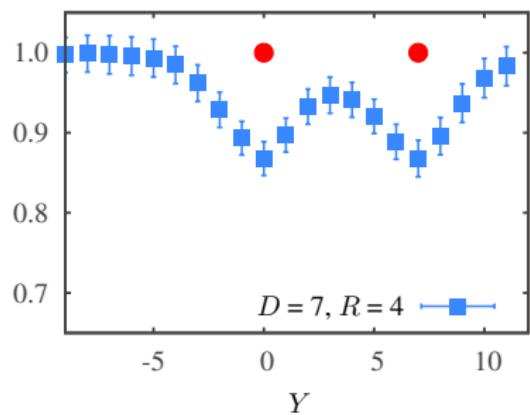
# Quark-Antiquark Pair



cross-section

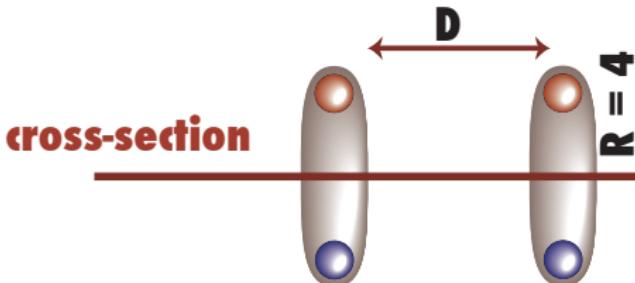


ratio of chiral condensate

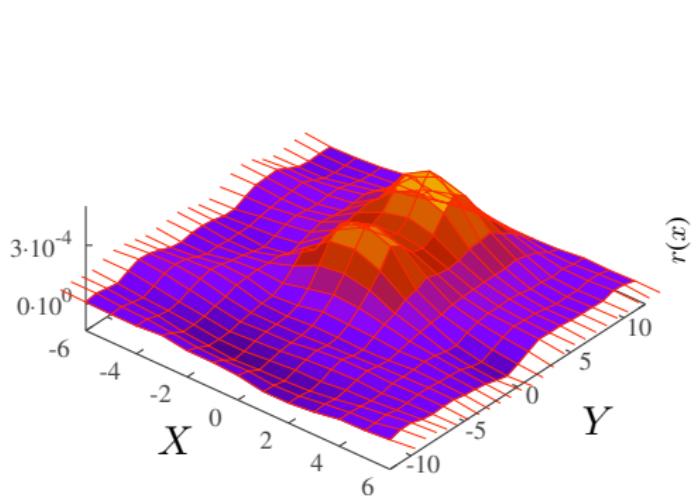


cf. lattice unit  $a \sim 0.11$  fm

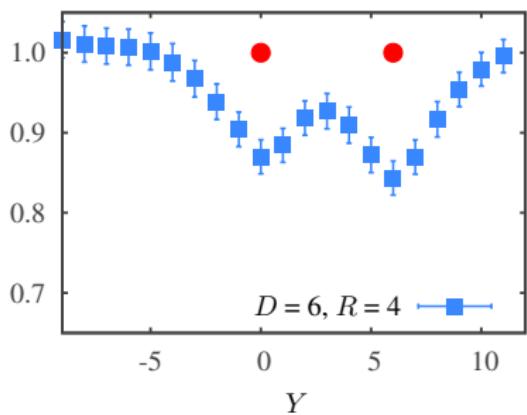
# Quark-Antiquark Pair



cross-section

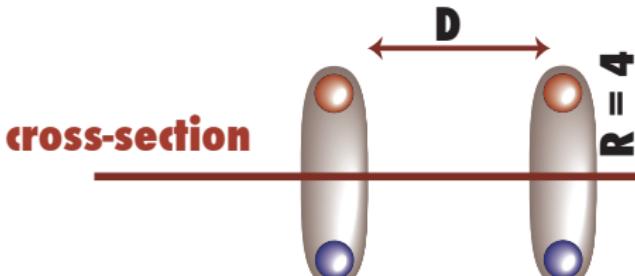


ratio of chiral condensate



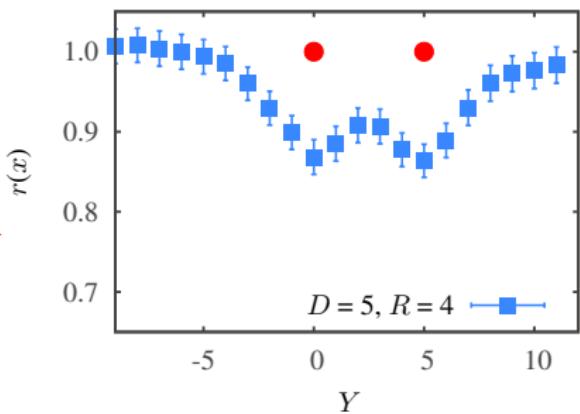
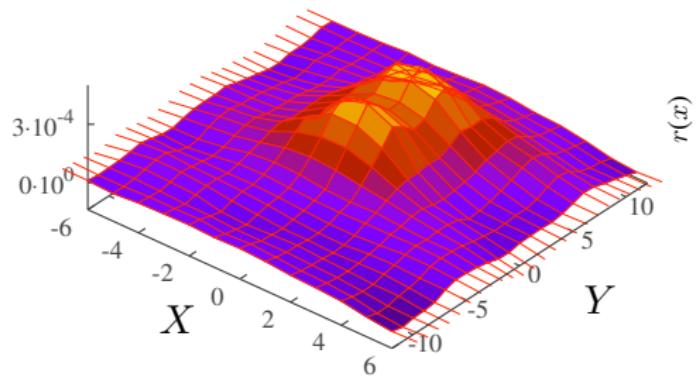
cf. lattice unit  $a \sim 0.11$  fm

# Quark-Antiquark Pair



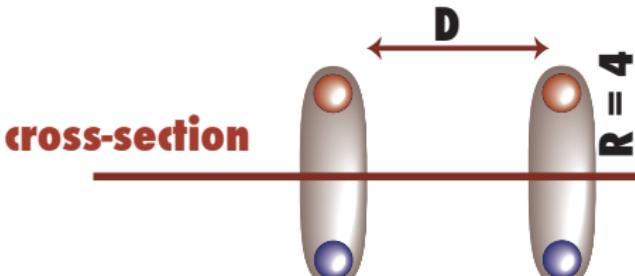
cross-section

ratio of chiral condensate

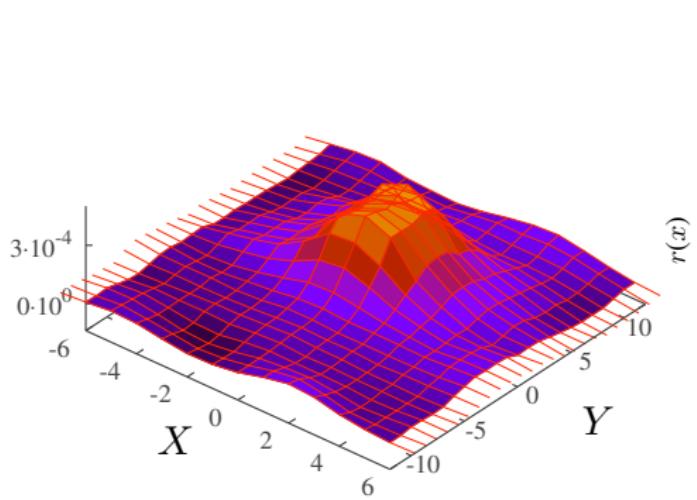


cf. lattice unit  $a \sim 0.11$  fm

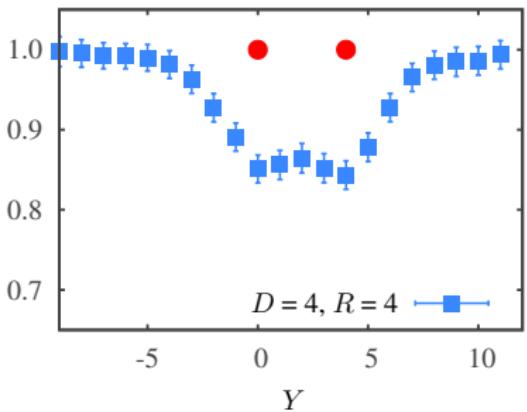
# Quark-Antiquark Pair



cross-section

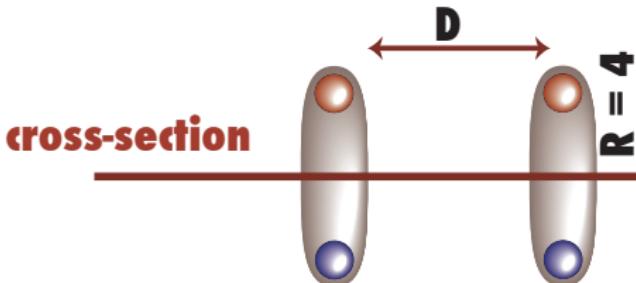


ratio of chiral condensate

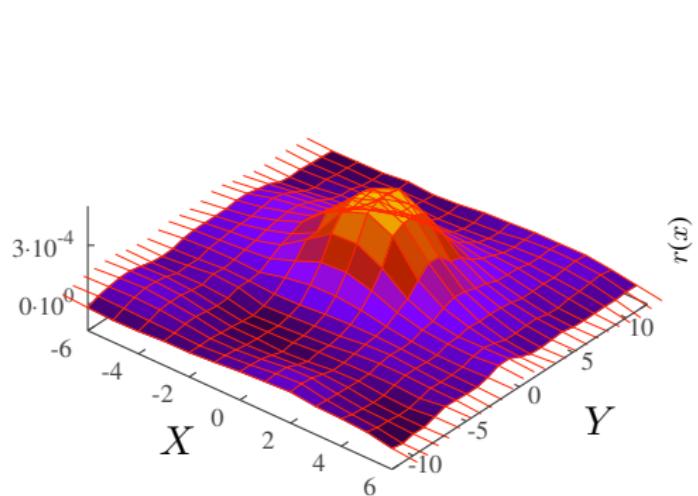


cf. lattice unit  $a \sim 0.11$  fm

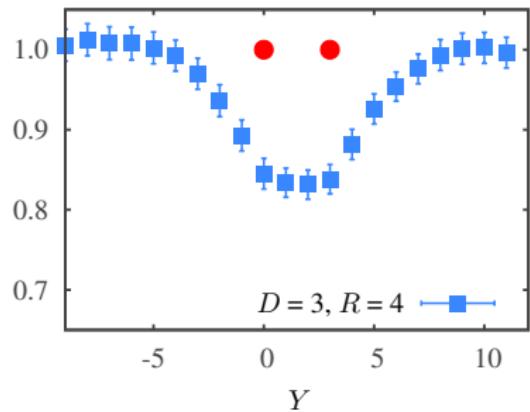
# Quark-Antiquark Pair



cross-section



ratio of chiral condensate

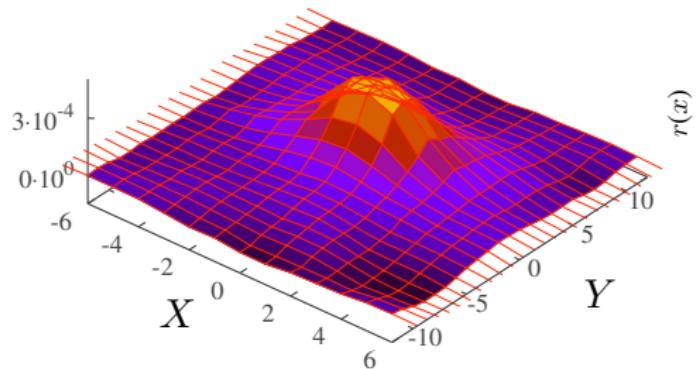
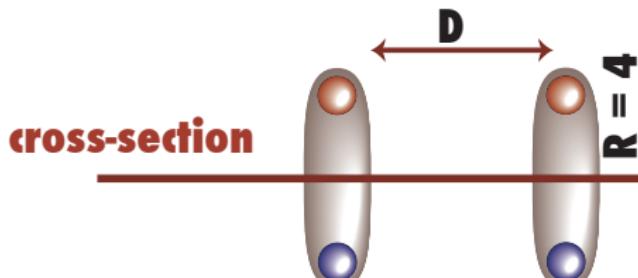


cf. lattice unit  $a \sim 0.11$  fm

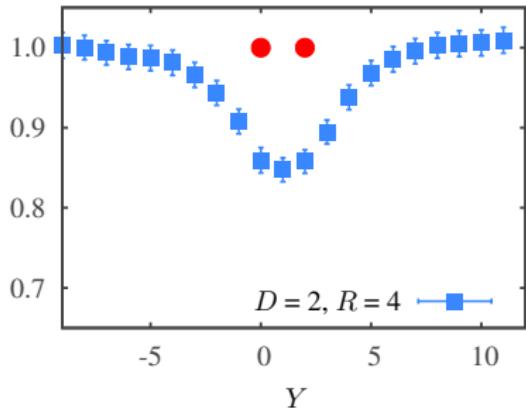
# Quark-Antiquark Pair

Partial restoration of chiral symmetry is the same as QQ-bar-system.

$$r_{Q\bar{Q}+Q\bar{Q}} = r_{Q\bar{Q}}$$



ratio of chiral condensate



cf. lattice unit  $a \sim 0.11$  fm

## ① Interquark Potential and Flux-tube Structure

## ② Chiral Condensate in Flux-tube

- Chiral Condensate in Quark-Antiquark System
- Chiral Condensate in 3-Quark System
- Chiral Condensate in Quark-Antiquark + Quark-Antiquark System

## ③ Summary

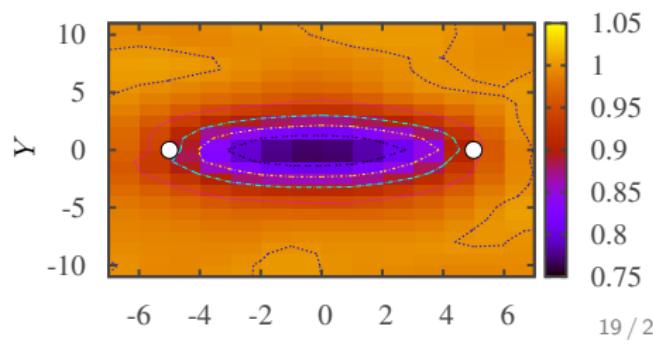
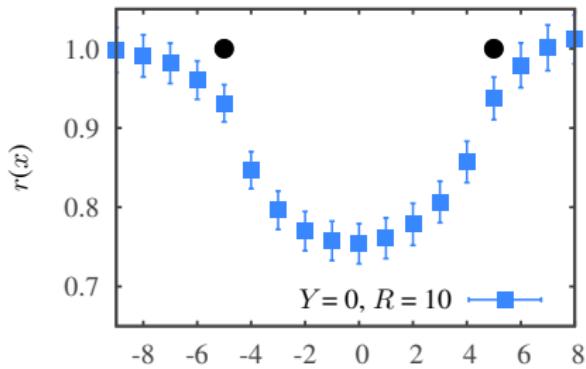
## Summary

- A color flux-tube appears between quarks, which produces linear confining potential.
- By measuring the chiral condensate  $\bar{q}q(x)$  inside flux-tube, we show partial restoration of chiral symmetry

$$\frac{\langle \bar{q}q \rangle_{\text{in flux}}}{\langle \bar{q}q \rangle_{\text{vacuum}}} = 0.7 \sim 0.8$$

⇒ effect of strong *color-electric fields* between quarks

- Partial restoration inside the flux-tube may change
  - light quark properties in QQq, QQqq ?
  - quark pair production rate ?



## Summary and Outlooks

In this work, we probe **local chiral condensate**  $\bar{q}q(x)$  inside the flux-tube.

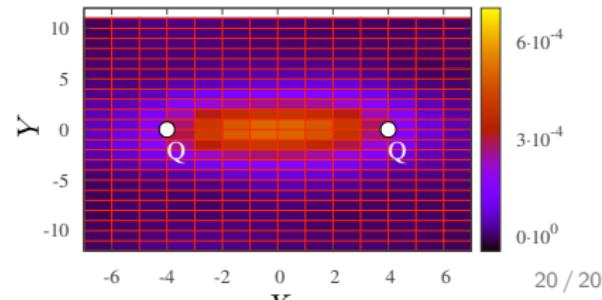
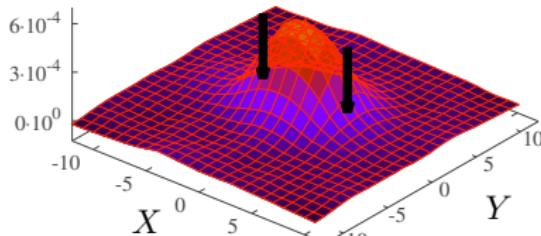
■ We can use various kinds of operators such as

- quark number density, energy density  
    ⇒ color charge screening, string breaking, hadronization
- topological charge density  
    ⇒  $U_A(1)$  anomaly inside hadron
- energy-momentum tensor  
    ⇒ energy of flux-tube, pressure of bag
- *and so on*

■ It is also possible to apply tetra-quark, meson-meson systems.

■ In addition to hadrons, color flux also appears after heavy ion collision.

■ We can also generalize this method in QGP phase at finite temperature.



4

## Appendix

# Thickness of Flux-tube and Chiral Symmetry Restoration

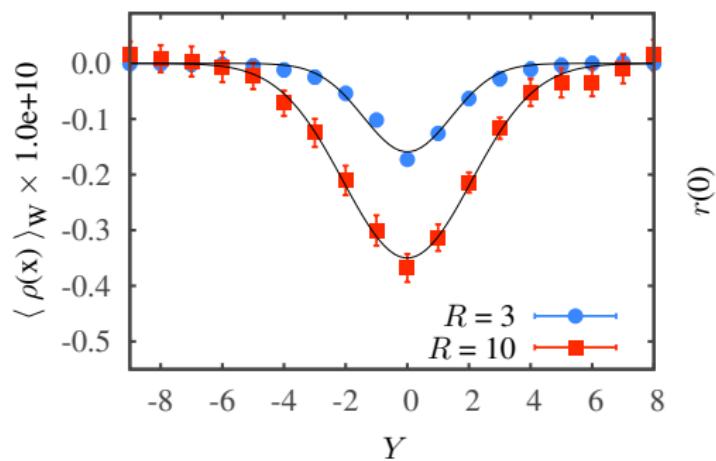
- By increasing the length of quantum flux-tube  $R$ ,  
the thickness  $\sigma$  grows as

$$\sigma \sim \ln R$$

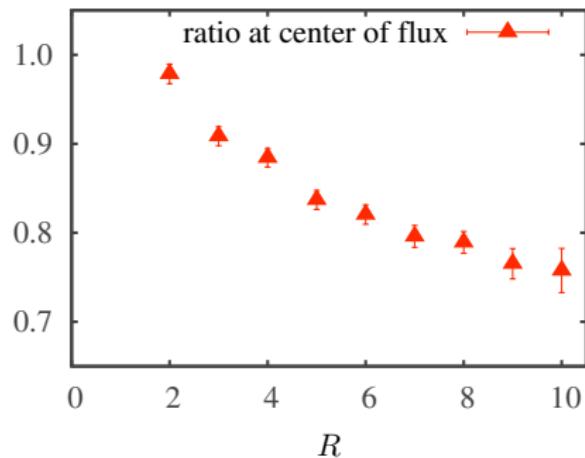
which is called as “Roughening.”<sup>1</sup>

- Partial restoration of chiral symmetry relates its thickness ?

Cross-section of flux-tube



Ratio of chiral condensate



<sup>1</sup>Hasenfratz-Hasenfratz-Hasenfratz '80, Luscher-Munster-Weisz '81

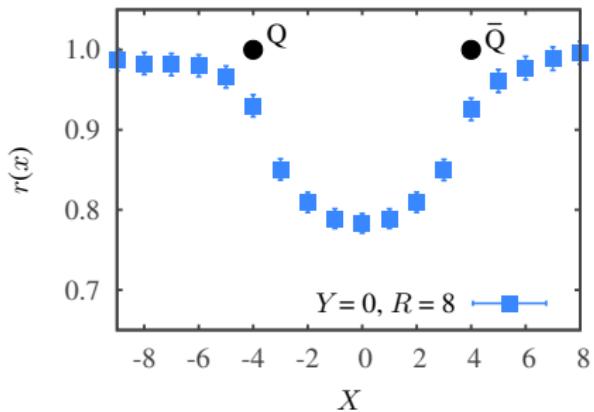
# Flux-tube and Chiral Condensate Ratio

flux-tube — action density distribution around static color sources

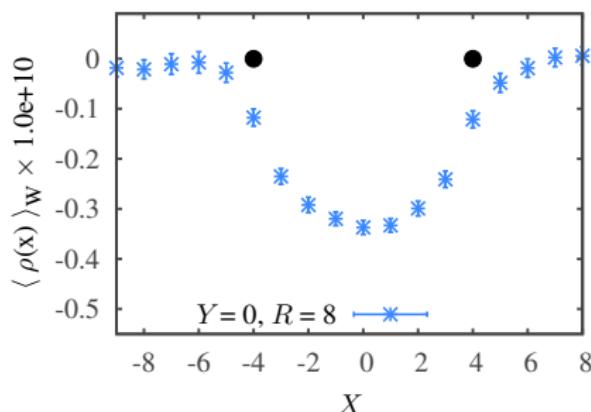
$$\langle \rho(x) \rangle_W \equiv \frac{\langle \rho(x) W(R, T) \rangle}{\langle W(R, T) \rangle} - \langle \rho \rangle$$

⇒ chiral condensate is reduced inside “flux-tube”

ratio of chiral condensate



action density distribution



# Quark Mass Dependence of Chiral Condensate Reduction

$16^3 \times 48$  lattice with low-lying 120 eigenmodes

- $m_{ud} = 0.015$  :  $m_\pi \sim 0.30$  GeV
- $m_{ud} = 0.050$  :  $m_\pi \sim 0.53$  GeV

