

$K \rightarrow \pi$ transition form factors and the transverse spin density

Hyeon-Dong Son

Inha University, Korea

@ Progress of J-PARC Hadron Physics,
Nov. 30. - Dec. 02. 2014.

in collaboration with
Seung-il Nam
Hyun-Chul Kim

Contents

- Motivation
- Kaon Tensor Transition Form Factors from the Nonlocal Chiral Quark Model
- Generalised Parton Distributions & Transverse Densities for the Kaon transition
- Numerical Results
- Summary

Motivation

- Semileptonic decays (K_{l3})
: CKM matrix elements $|V_{us}|$
- **Tensor** transition provides non-standard type interaction

[I. Baum et al, Phys. Rev. D 84, 074503 (2011)]

- **Transverse charge & spin densities of the transition**

[C. Carlson and M. Vanderhaeghen,
Phys. Rev. Lett. 100, 032004 (2008)]

Kaon l3 decay

The decay amplitude

$$T_{K \rightarrow l \mu \pi} = \frac{G_F}{\sqrt{2}} \sin \theta_c [W^\mu(p_l, p_\nu) F_\mu(p_\pi, p_K)]$$

$$G_F = 1.116 \times 10^{-5} \text{GeV}^{-2}$$

Weak leptonic element

$$W^\mu(p_l, p_\nu) = \bar{u}(p_\nu) \Gamma^\mu \nu(p_l)$$

Hadronic matrix element

$$F_\mu(p_\pi, p_K) = c \langle \pi(p_\pi) | \Gamma_\mu | K(p_K) \rangle$$

수식, Dirac structure, flavor 더 명확하게 적어서 improve 할 것.
su가 되려면 $K^0 \rightarrow \pi^-$ 여야 함.
예를 들어 마지막 식에 $F \rightarrow \Gamma$

Hadronic Matrix Elements

❖ Vector transition

$$F_{\mu}^{K^0}(p_{\pi}, p_K) = \langle \pi^-(p_{\pi}) | \bar{s} \gamma_{\mu} u | K^0(p_K) \rangle = (p_K + p_{\pi})_{\mu} f_{l+}(t) + (p_K - p_{\pi})_{\mu} f_{l-}(t)$$

❖ Tensor transition

$$F_{\mu\nu}^{K^0}(p_{\pi}, p_K) = \langle \pi^-(p_{\pi}) | \bar{s} \sigma_{\mu\nu} u | K^0(p_K) \rangle = \frac{p_{K\mu} p_{\pi\nu} - p_{K\nu} p_{\pi\mu}}{m_K} B_T^{K\pi}(t)$$

❖ Scalar transition

$$F^{K^0}(p_{\pi}, p_K) = \langle \pi^-(p_{\pi}) | \bar{s} u | K^0(p_K) \rangle = -\frac{m_K^2 - m_{\pi}^2}{m_s - m_u} f_0(t)$$

Nonlocal Chiral Quark Model

$$S_{\text{eff}} = -N_c \text{Tr} \log \left[i\not{\partial} + i\hat{m} + i\sqrt{M(i\partial)} U^{\gamma_5} \sqrt{M(i\partial)} \right]$$

- ❖ The chiral effective action
derived from the instanton vacuum

$$U^{\gamma_5} = \exp \left[\frac{i\gamma_5}{f_\phi} (\lambda \cdot \phi) \right]$$

- ❖ No free parameter
 - Average Instanton size & separation

$$\bar{\rho} \approx \frac{1}{3} \text{ fm} \quad \bar{R} \approx 1 \text{ fm}$$

- ❖ Nonlocality
 - Momentum-dependent dynamical quark mass
- ❖ Nicely reproduces pion properties: F_{π} , EMFF
- ❖ Explicit SU(3) symmetry breaking

$$\hat{m} = \text{diag}(m_u, m_d, m_s), \quad m_u = m_d = 5 \text{ MeV}, \quad m_s = 150 \text{ MeV}$$

[D. Diakonov, Instantons at work, arXiv:hep-ph/0212026v4]

Nonlocal Chiral Quark Model

modbess ff 써서 수치계산 하는데 무슨 문제가 있었지?

- ❖ Momentum-dependent dynamical quark mass

$$\sqrt{M(i\partial)} = \sqrt{M_0 f(m) F^2(i\partial)}$$

$$F(k) = \frac{k}{\Lambda} \left[I_0\left(\frac{k}{2\Lambda}\right) K_1\left(\frac{k}{2\Lambda}\right) - I_1\left(\frac{k}{2\Lambda}\right) K_0\left(\frac{k}{2\Lambda}\right) - \frac{2\Lambda}{k} I_1\left(\frac{k}{2\Lambda}\right) K_1\left(\frac{k}{2\Lambda}\right) \right]$$

$$F_N(k) = \left(\frac{2N\Lambda^2}{2N\Lambda^2 + k^2} \right)^N \quad \Lambda = 1/\bar{\rho} = 600 \text{ MeV}$$

- ❖ Current quark mass correction $f(m)$

$$f(m) = \sqrt{1 + \frac{m^2}{d^2}} - \frac{m}{d}, d \approx 198 \text{ MeV}$$

- ❖ Zero-momentum dynamical quark mass M_0

$$M_0 \approx 350 \text{ MeV}$$

[M. Musakhanov Eur.Phys.J.C9,235(1999)]

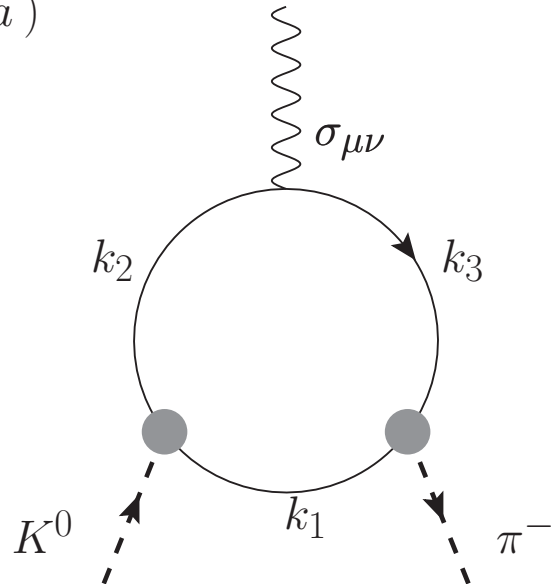
$$\text{Pion Decay Constant } f_\pi \approx 93 \text{ MeV}$$

Calculation

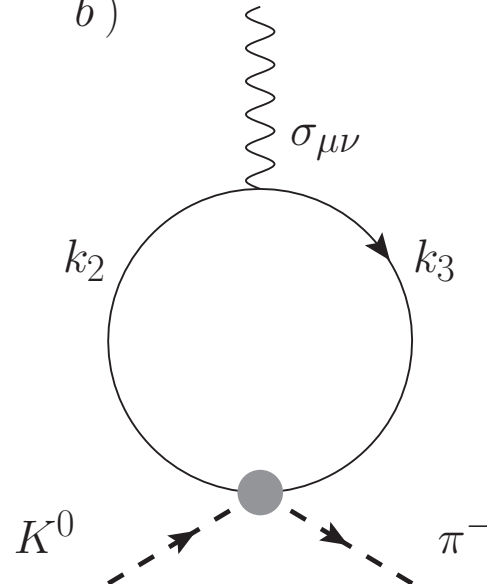
$$\langle \pi^-(p_f) | \bar{s} \sigma_{\mu\nu} u | K^0(p_i) \rangle =$$

$$\frac{8N_c}{f_\pi f_K} \int \frac{d^4 l}{(2\pi)^4} \left[\frac{M_{2d} \sqrt{M_{1u}} \sqrt{M_{3s}}}{G_{u1} G_{d2} G_{s3}} \left(k_{1\mu} (k_{2\nu} \bar{M}_{3s} - k_{3\nu} \bar{M}_{2d}) + k_{2\mu} (k_{3\nu} \bar{M}_{1u} - k_{1\nu} \bar{M}_{3s}) \right. \right. \\ \left. \left. + k_{3\mu} (k_{1\nu} \bar{M}_{2d} - k_{2\nu} \bar{M}_{1u}) \right) - \frac{\sqrt{M_{1u}} \sqrt{M_{3s}}}{2G_{u1} G_{s3}} (k_{3\mu} k_{1\nu} - k_{3\nu} k_{1\mu}) \right]$$

a)



b)



$$G_{fi} = k_i^2 + \bar{M}_{if}^2$$

$$\bar{M}_{if} = m_f + M(k_i, m_f)$$

$$k_1 = l - \frac{p_i}{2} - \frac{q}{2} \quad k_2 = l + \frac{p_i}{2} - \frac{q}{2}$$

$$k_3 = l + \frac{p_i}{2} + \frac{q}{2}$$

Experimental Values

$$f_\pi = 93 \text{ MeV}, \quad f_K = 113 \text{ MeV}$$

$$m_\pi = 140 \text{ MeV}, \quad m_K = 495 \text{ MeV}$$

QCD RG Evolution

♣ At the leading order

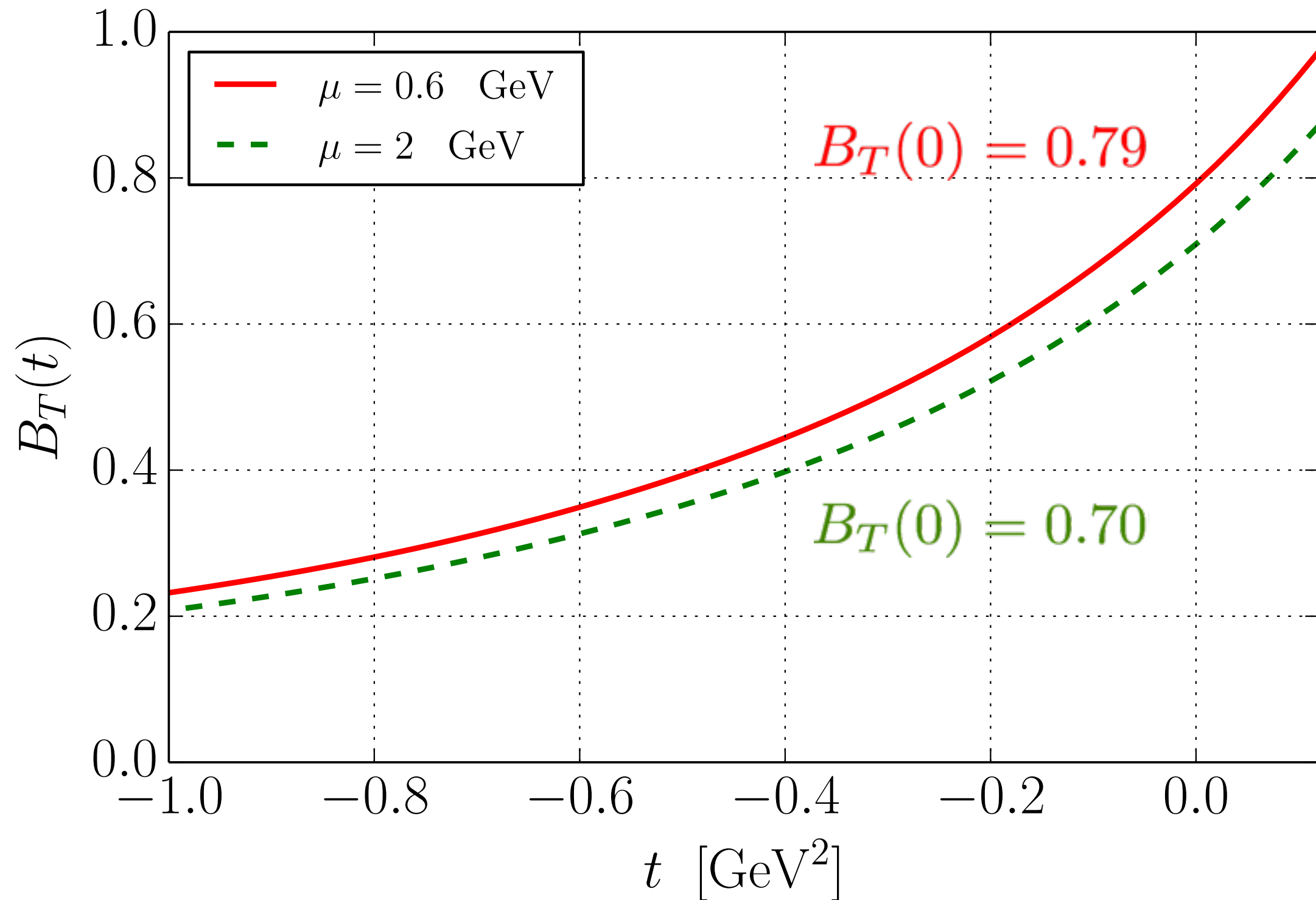
$$B_T(t, \mu) = B_T(t, \mu_0) \left[\frac{\alpha(\mu)}{\alpha(\mu_0)} \right]^{\gamma/(2\beta_0)}$$

$$\gamma = \frac{8}{3}, \quad \beta_0 = 11\frac{N_c}{3} - 2\frac{N_f}{3} \quad (N_c = 3, N_f = 3)$$

$$\Lambda_{\text{QCD}} = 0.25 \text{ GeV} \quad [\text{Glück et al. Zeits.Für.Phys. C.67(3),} \\ \text{Barone et al. Phys.Repts., 359(1-2), 1-168.}]$$

Numerical Results

Display the form factor only in the space like region.
Discuss the slope parametrisation and the values.



Lattice Results

📌 Present work

$$f_T^{K\pi}(t) = \frac{m_K + m_\pi}{2m_K} B_T^{K\pi}(t)$$

$$B_T^{K\pi}(0) = 0.70 \rightarrow f_T^{K\pi}(0) = 0.45$$

📌 I. Baum et al, $\mu = 2$ GeV

$$\langle \pi^0 | \bar{s} \sigma^{\mu\nu} d | K^0 \rangle = (p_\pi^\mu p_K^\nu - p_\pi^\nu p_K^\mu) \frac{\sqrt{2} f_T^{K\pi}(q^2)}{M_K + M_\pi}$$

$$f_T^{K\pi}(0) = 0.417 (14_{\text{stat}}) (5_{\text{syst}}) = 0.417 (15)$$

Extrapolated to Physical meson masses

📌 I. V. Ajinenko et al. Physics Letters B 574, 14 (2003).

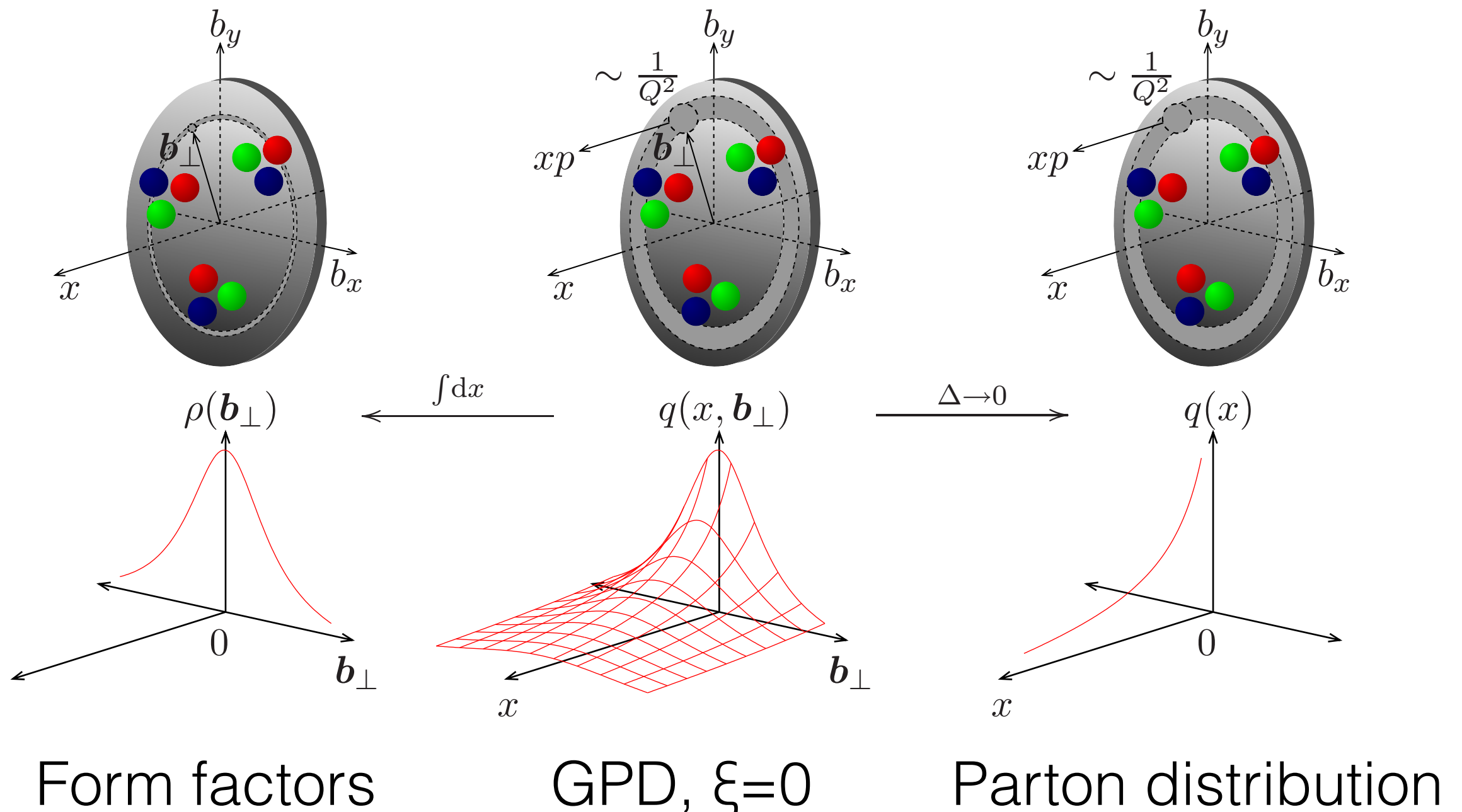
The limits on the possible tensor and scalar couplings are derived:

$$f_T/f_+(0) = 0.021_{-0.075}^{+0.064}(\text{stat}) \pm 0.026(\text{syst}),$$

$$f_S/f_+(0) = 0.002_{-0.022}^{+0.020}(\text{stat}) \pm 0.003(\text{syst}).$$

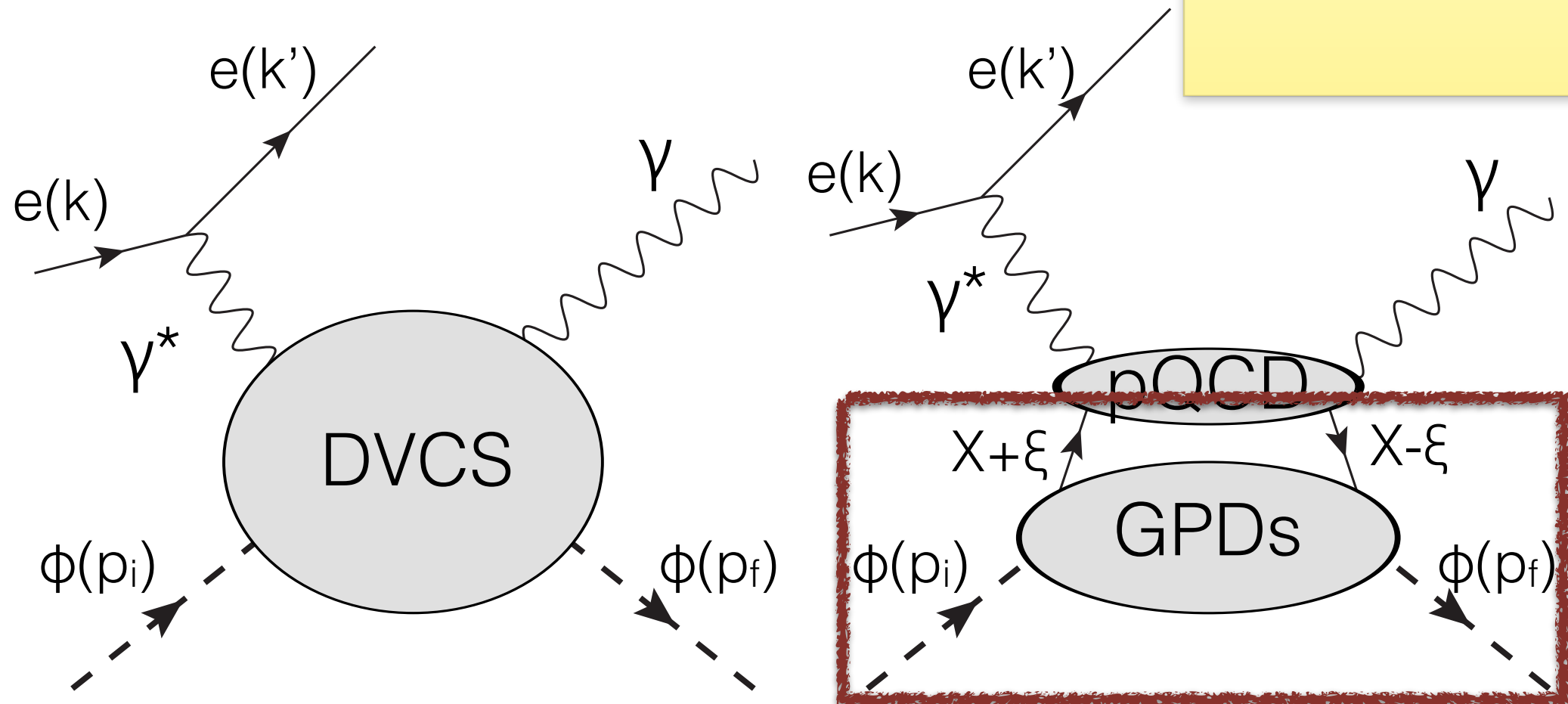
Generalised Parton Distributions

[D. Brömmel, Pion Structure From the Lattice, Regensburg Univ., Thesis]



DVCS & GPDs

How about the corresponding processes to the weak transitions? are they measurable from an experiment?



Factorization : [Collins, J., & Freund, A. (1999).Phys. Rev. D, 59(7), 074009.]

- Light cone frame

$$P^\mu = \frac{1}{2} (p_i^\mu + p_f^\mu), \quad q^\mu = p_f^\mu - p_i^\mu, \quad t = q^2 \quad n_\pm = \frac{1}{\sqrt{2}} (1, 0, 0, \pm 1)$$

$$v^\pm = \frac{1}{\sqrt{2}} (v^0 \pm v^3), \quad v_\perp = (v^1, v^2) \quad \xi = \frac{p_i^+ - p_f^+}{p_i^+ + p_f^+}$$

Generalised Parton Distribution

lightcone gauge: $n \cdot A = 0$

- Generalised parton distributions (GPDs)

$$2P^+ H_\phi^q(X, \xi, t) = \int \frac{dz^-}{2\pi} e^{iX P^+ z^-} \langle \phi(p_f) | \bar{\psi}^q(-\frac{z^-}{2}) \gamma^+ \left[-\frac{z^-}{2}, \frac{z^-}{2} \right] \psi^q(\frac{z^-}{2}) | \phi(p_i) \rangle \Big|_{z^+ = z_\perp = 0}$$

$$\frac{P^{[+t^j]}}{m_\phi} E_\phi^q(X, \xi, t) = \int \frac{dz^-}{2\pi} e^{iX P^+ z^-} \langle \phi(p_f) | \bar{\psi}^q(-\frac{z^-}{2}) i\sigma^{+j} \left[-\frac{z^-}{2}, \frac{z^-}{2} \right] \psi^q(\frac{z^-}{2}) | \phi(p_i) \rangle \Big|_{z^+ = z_\perp = 0}$$

- Kaon Transition GPDs for $K^0 \rightarrow \pi^-$

$$2P^+ H_\phi^{K\pi}(X, \xi, t) = \int \frac{dz^-}{2\pi} e^{iX P^+ z^-} \langle \pi^-(p_f) | \bar{s}(-\frac{z^-}{2}) \gamma^+ \left[-\frac{z^-}{2}, \frac{z^-}{2} \right] u(\frac{z^-}{2}) | K^0(p_i) \rangle \Big|_{z^+ = z_\perp = 0}$$

$$\frac{P^{[+t^j]}}{m_K} E_\phi^{K\pi}(X, \xi, t) = \int \frac{dz^-}{2\pi} e^{iX P^+ z^-} \langle \pi^-(p_f) | \bar{s}(-\frac{z^-}{2}) i\sigma^{+j} \left[-\frac{z^-}{2}, \frac{z^-}{2} \right] u(\frac{z^-}{2}) | K^0(p_i) \rangle \Big|_{z^+ = z_\perp = 0}$$

Expanding the nonlocal operators \rightarrow Generalised Form Factors

Generalised Form Factors

- Generalised form factors

Vector

$$\begin{aligned} & \langle \phi^a(p_f) | \psi^\dagger(0) \gamma_{\{\mu} i \overleftrightarrow{D}_{\mu_1} \dots i \overleftrightarrow{D}_{\mu_{n-1}} \} \psi(0) | \phi^b(p_i) \rangle \\ &= 2P_{\{\mu} P_{\mu_1} \dots P_{\mu_{n-1}} \} A_{n0}(t) \\ &+ \sum_{k=2 \text{ even}}^n q_{\{\mu} q_{\mu_1} \dots q_{\mu_{k-1}} P_{\mu_k} P_{\mu_{n-1}} \} 2^{-k} A_{nk}(t) \end{aligned}$$

Tensor

$$\begin{aligned} & \langle \phi(p_f) | \psi^\dagger(0) \sigma_{[\mu\nu} i \overleftrightarrow{D}_{\mu_1} \dots i \overleftrightarrow{D}_{\mu_{n-1}]} \psi(0) | \phi(p_i) \rangle \\ &= \frac{p_{[\mu} q_{\nu]} - q_{\mu} p_{\nu}}{m_\phi} \sum_{i=even}^{n-1} q_{\mu_1} \dots q_{\mu_i} P_{\mu_{i+1}} P_{\mu_{n-1}} B_{ni}(t) \end{aligned}$$

Polynomiality & Transverse charge density

- Polynomiality: Mellin moments of the GPDs

$$\int dX X^{n-1} H(X, \xi, t) \Rightarrow \text{Finite polynomial of } \xi$$

$$\int dX H(X, \xi, t) = A_{10}(t) \quad \rightarrow \text{Vector FF}$$

$$\int dX X H(X, \xi, t) = A_{20}(t) + \xi^2 A_{22}(t) \quad \rightarrow \text{EMT FFs}$$

- Transverse densities:

2D Fourier transformation into the impact parameter b_{\perp} at $\xi=0$

$$\begin{aligned} F(b_{\perp}^2) &= \int dq_{\perp}^2 e^{-ib_{\perp} \cdot q_{\perp}} \int dX H(X, q_{\perp}^2, \xi = 0) \\ &= \int dq_{\perp}^2 e^{-ib_{\perp} \cdot q_{\perp}} F(q_{\perp}^2) \end{aligned}$$

Probability distribution of the partons inside the hadrons
in the transverse impact parameter plane

GFFs & Transverse Densities for the Ka

K⁰ 스테이트에 (p_i)가 빠져있다 πππππ
π

- Generalised form factors for kaon transitions: n =1

$$\langle \pi^-(p_f) | \bar{s}(0) \gamma_\mu u(0) | K^0 \rangle = 2P_\mu A_{10}^{K\pi}(t) + q_\mu C_{10}^{K\pi}(t)$$

$$\langle \pi^-(p_f) | \bar{s}(0) \sigma_{\mu\nu} u(0) | K^0 \rangle = \frac{p_{i\mu} p_{f\nu} - p_{i\nu} p_{f\mu}}{m_K} B_{10}^{K\pi}(t)$$

- Quarks with definite transverse polarisation **s**

$$\frac{1}{2} \bar{\psi} [\gamma^+ - s^j i \sigma^{+j} \gamma_5] \psi \quad \vdots \quad \sigma^{\mu\nu} \gamma_5 = -\frac{1}{2} \epsilon^{\mu\nu\alpha\beta} i \sigma_{\alpha\beta}$$

- Transverse transition density, ξ=0

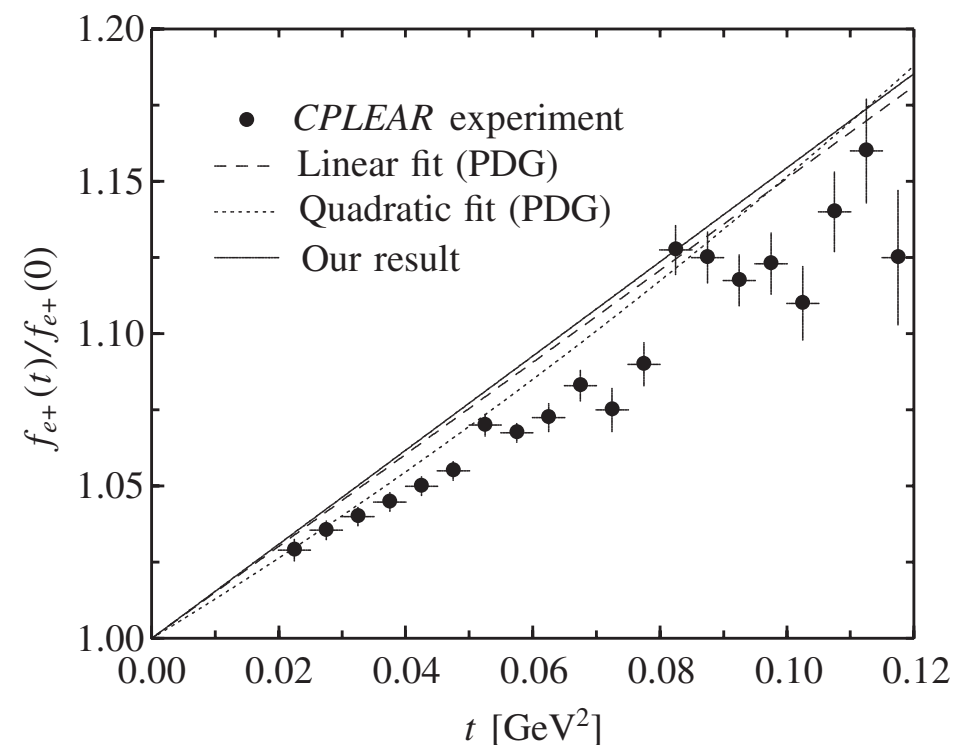
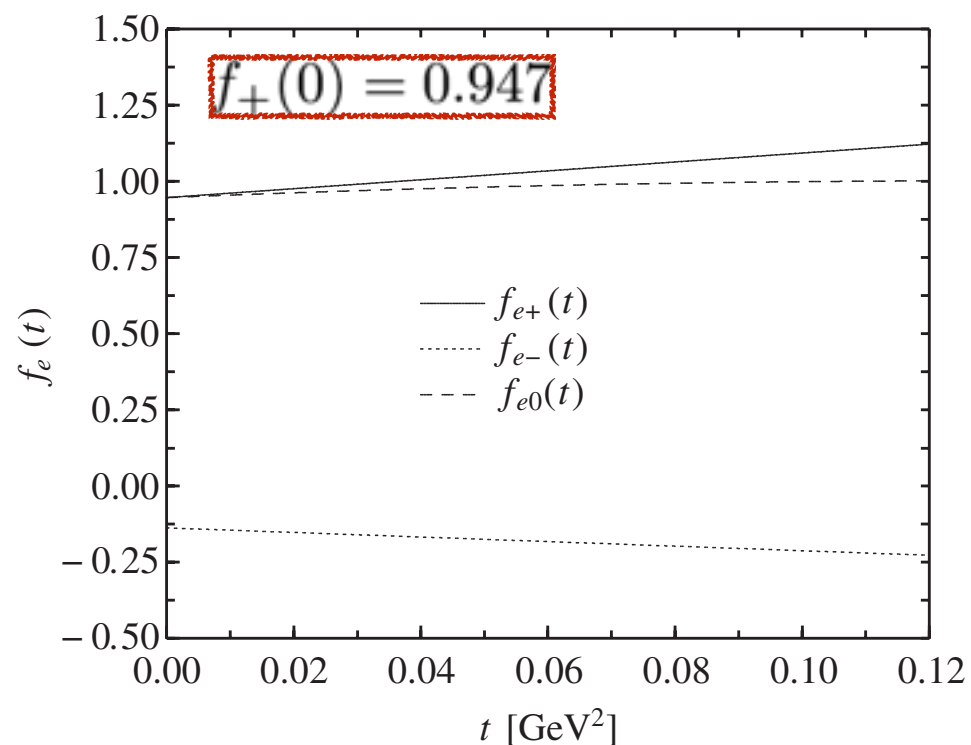
$$\rho^{K\pi}(b_\perp, s_\perp) = \int dX \rho(X, b_\perp, s_\perp) = \frac{1}{2} \left[A_{10}^{K\pi}(b_\perp^2) - \frac{s_\perp^i \epsilon^{ij} b_\perp^j}{m_K} \frac{\partial B_{10}^{K\pi}(b_\perp^2)}{\partial b_\perp^2} \right]$$

Transverse charge & spin density

How the quark spin is distributed
in the transverse plane during the K- π transition process

$$\rho^{K\pi}(b_{\perp}, s_{\perp}) = \int dX \rho(X, b_{\perp}, s_{\perp}) = \frac{1}{2} \left[A_{10}^{K\pi}(b_{\perp}^2) - \frac{s_{\perp}^i \epsilon^{ij} b_{\perp}^j}{m_K} \frac{\partial B_{10}^{K\pi}(b_{\perp}^2)}{\partial b_{\perp}^2} \right]$$

$$A_{10}^{K\pi}(b_{\perp}^2) \rightarrow f_+(b_{\perp}^2), \quad B_{10}^{K\pi}(b_{\perp}^2) \rightarrow B_T^{K\pi}(b_{\perp}^2)$$

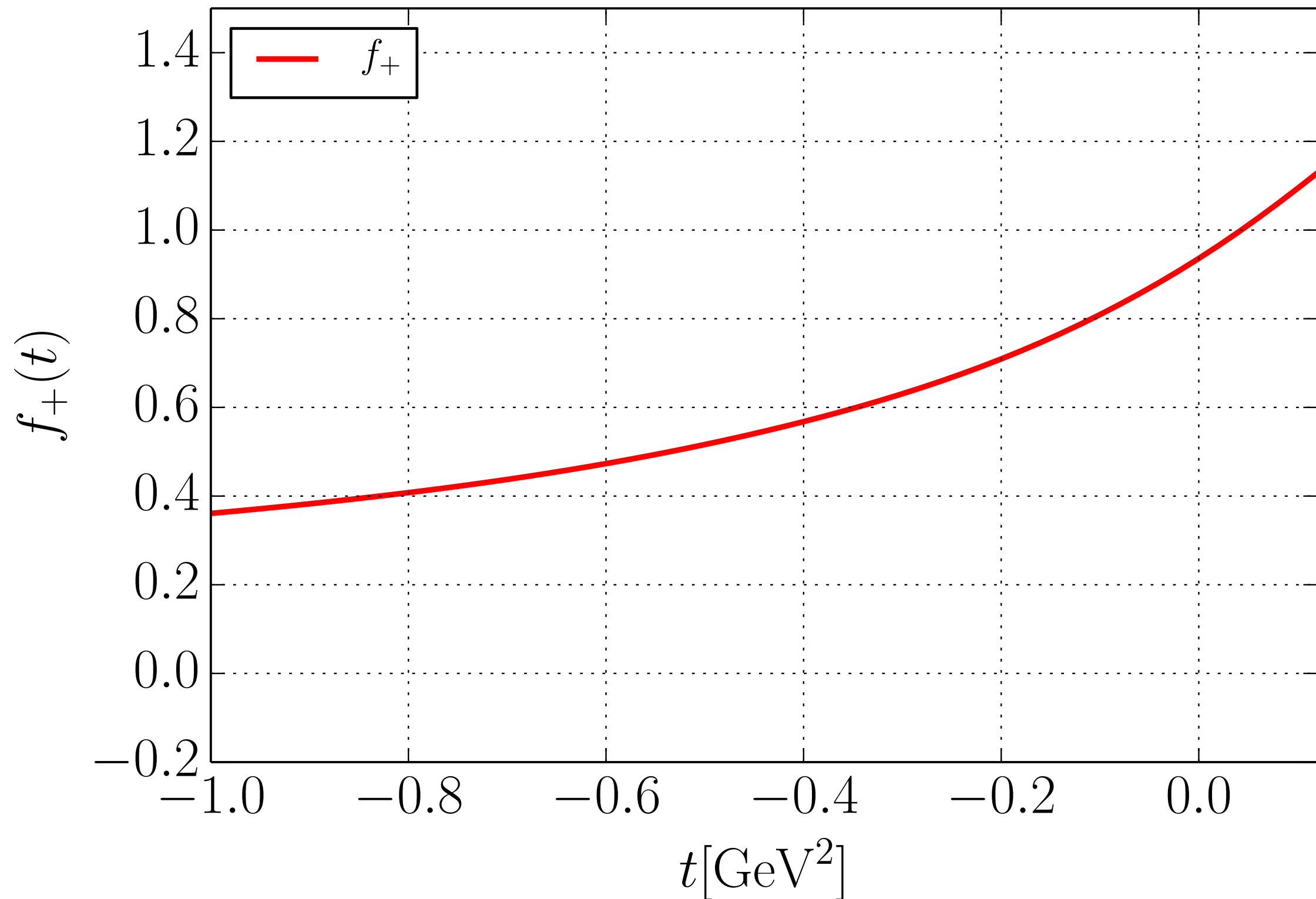


$$m_l^2 < t < (m_K - m_{\pi})^2 \approx 0.12 \text{ GeV}^2$$

[S.-i. Nam and H.-Ch. Kim, Phys. Rev. D 75, 094011 (2007).]

Vector form factor $f_+(t)$

The local & nonlocal things might be misleading. It's better to show only the total result.



P-pole parametrisation



p-pole parametrisation

is this slide necessary?

$$F(t) = \frac{F(0)}{\left(1 - \frac{t}{pM^2}\right)^p}$$

$$f_+(t)$$

$$B_T^{K\pi}(t)$$

p

1.31

2.2

M [GeV]

0.85

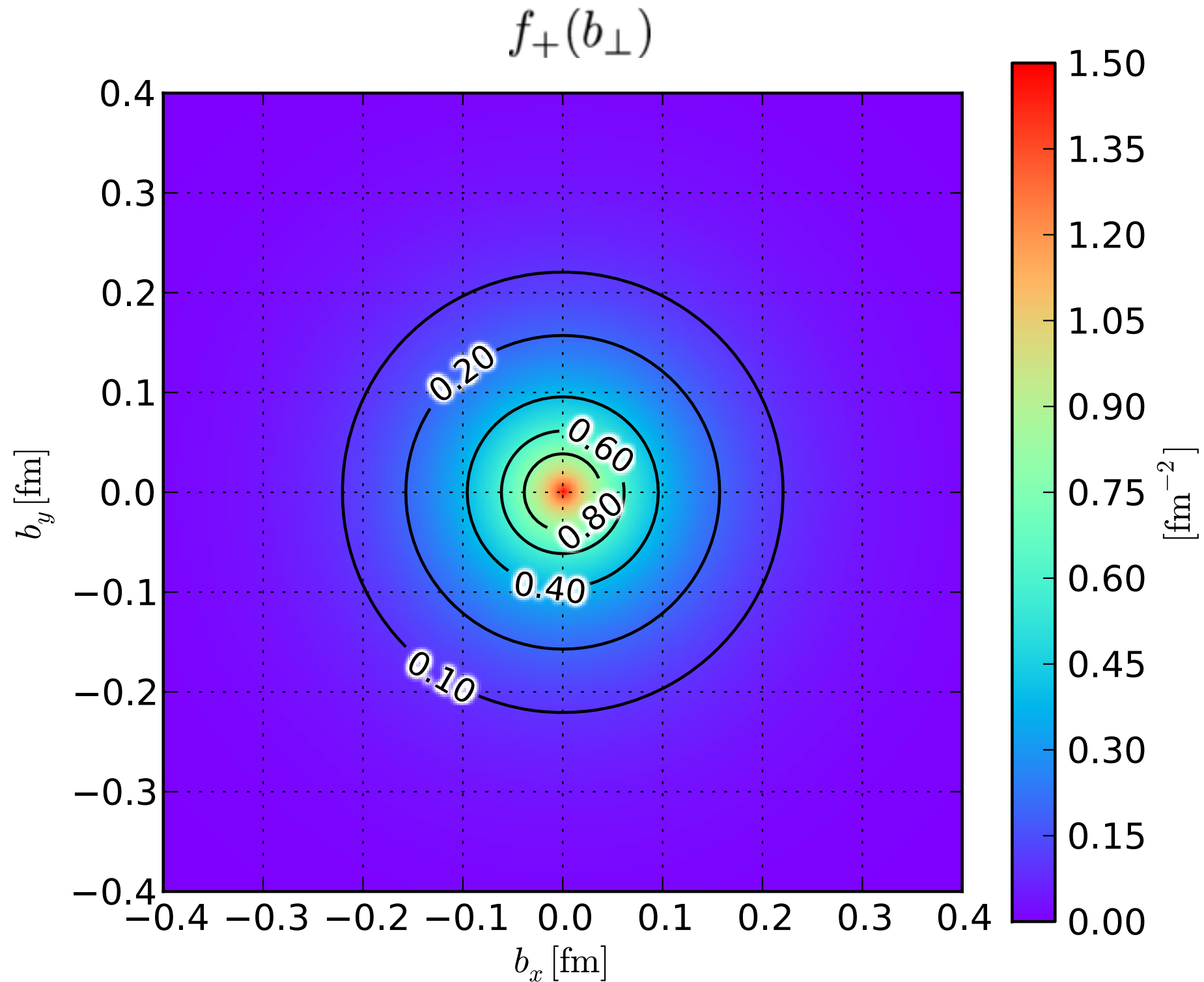
0.78

t=0

0.95

0.70

Transverse Charge Density

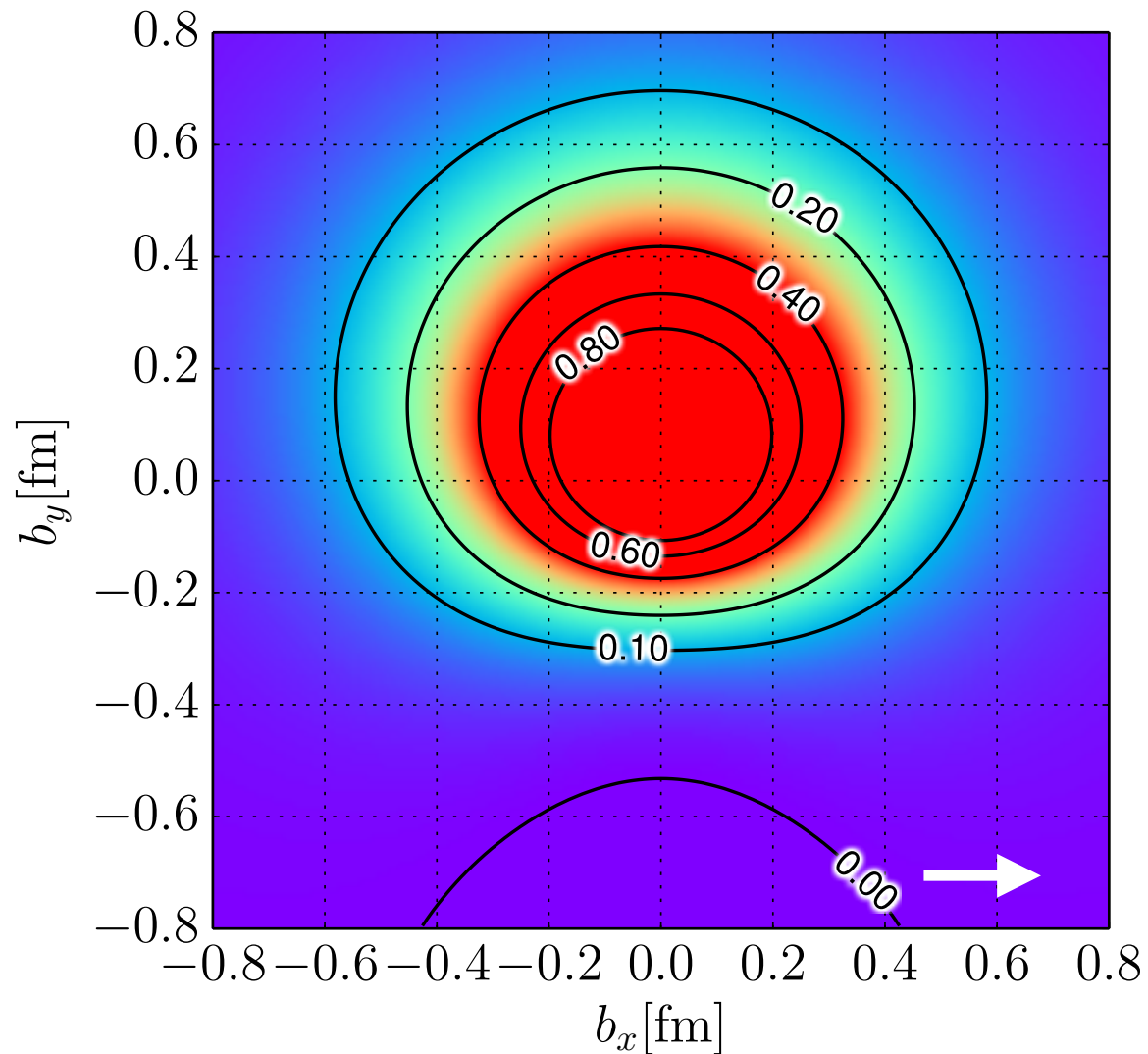


Transverse Spin Density

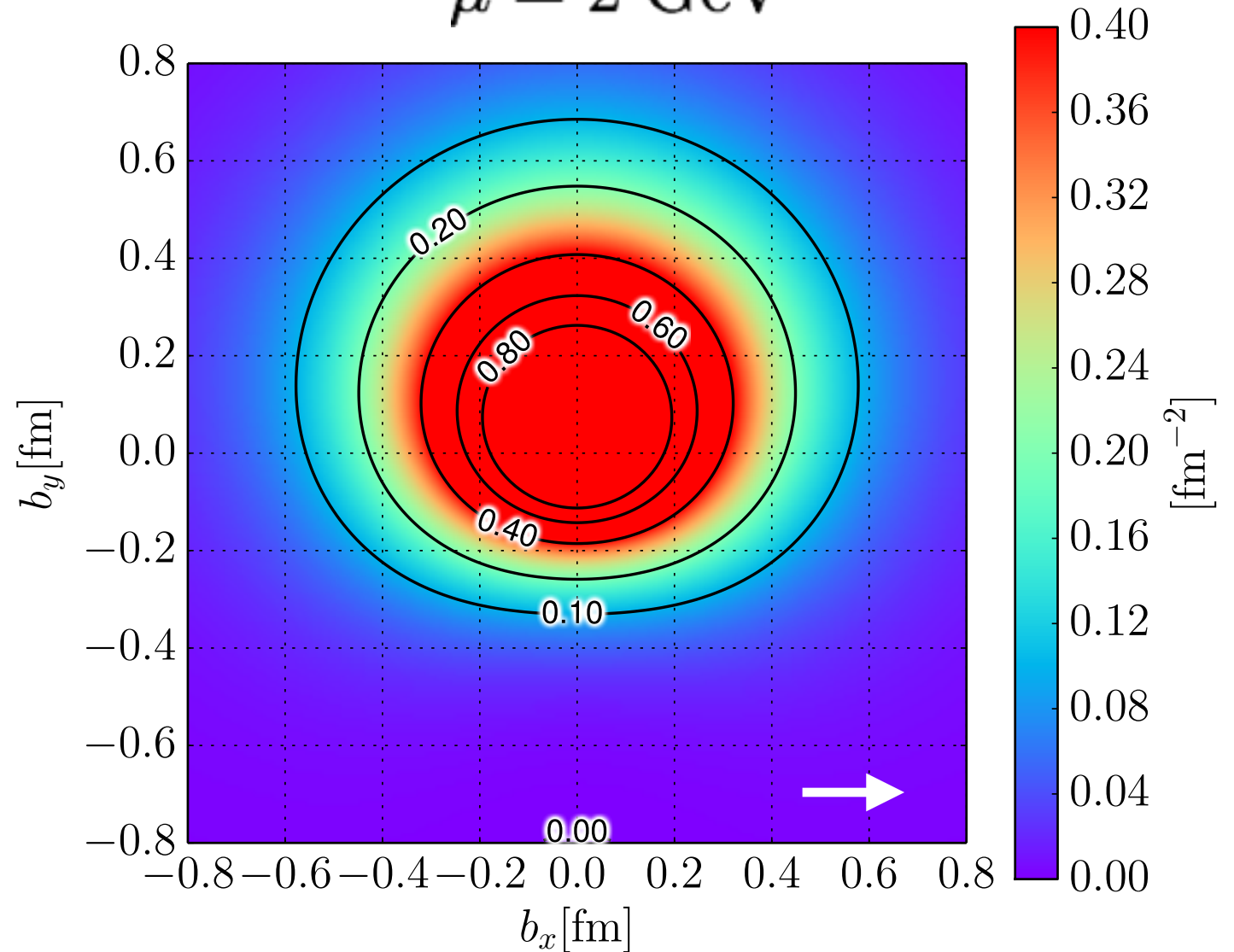
그림 업데이트할 것.
이 슬라이드랑 이 다음 슬라이드를 합칠 수 있을 것 같다.

$$\rho^{K\pi}(b_{\perp}, s_{b_x} = 1)$$

$\mu = 0.6 \text{ GeV}$

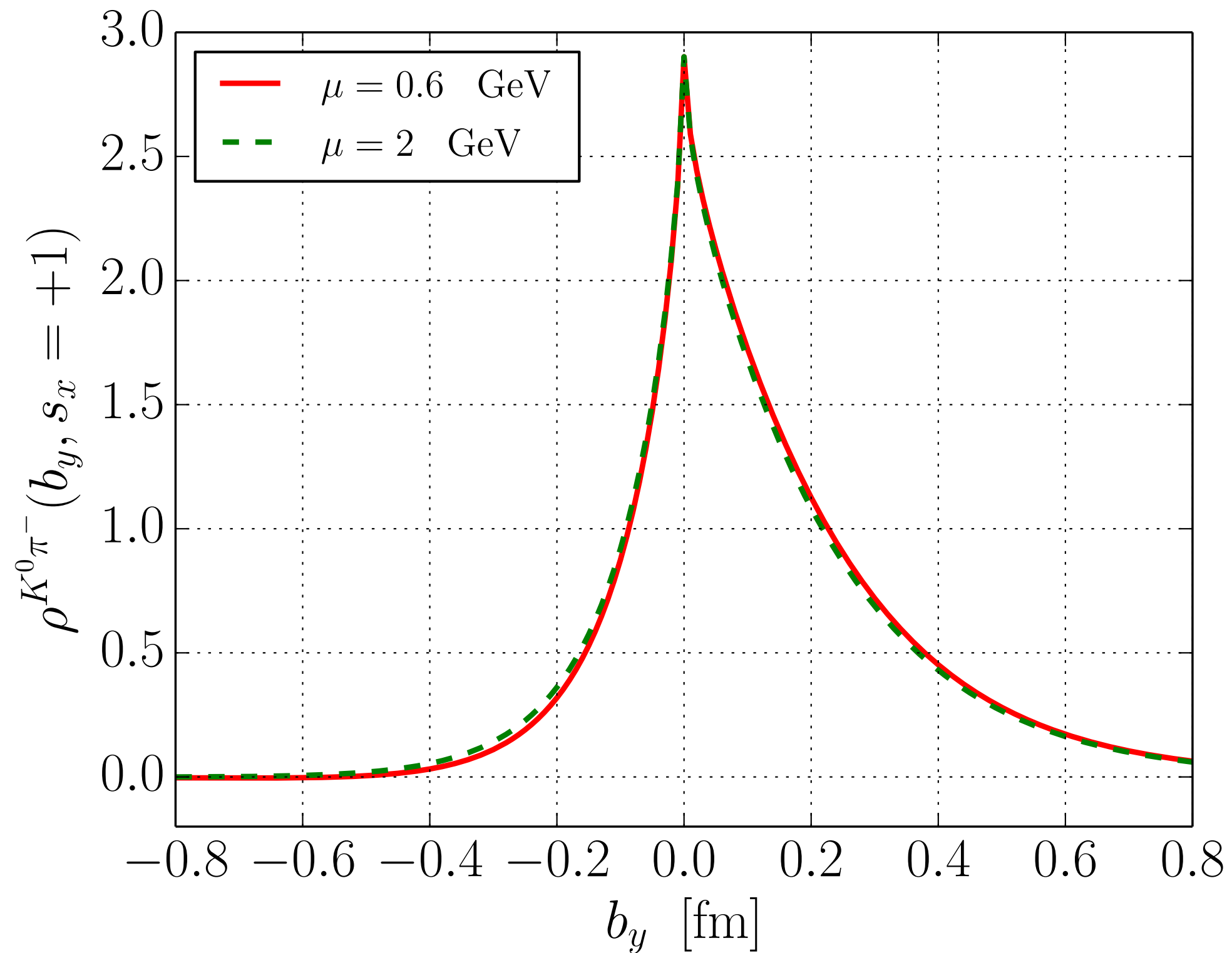


$\mu = 2 \text{ GeV}$



$$\langle b_y \rangle^{K\pi} = \frac{\int d^2 b_{\perp} b_y \rho^{K\pi}(b_{\perp}, s_{\perp})}{\int d^2 b_{\perp} \rho^{K\pi}(b_{\perp}, s_{\perp})} = \frac{1}{2m_K} \frac{B_T^{K\pi}(t=0)}{f_+(t=0)} = (0.17, 0.15) \text{ fm}$$

Transverse Spin Density



Summary & Outlook

- K to pi tensor transition form factor
- In good agreement with the lattice result
- Distorted spin structure during the transition process when the quark spin is polarized
- Transition GPDs of the pseudo scalar mesons

Thank you very much!