

Influence of threshold effects induced by charmed meson rescattering

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For Workshop on Progress on J-PARC Hadron Physics

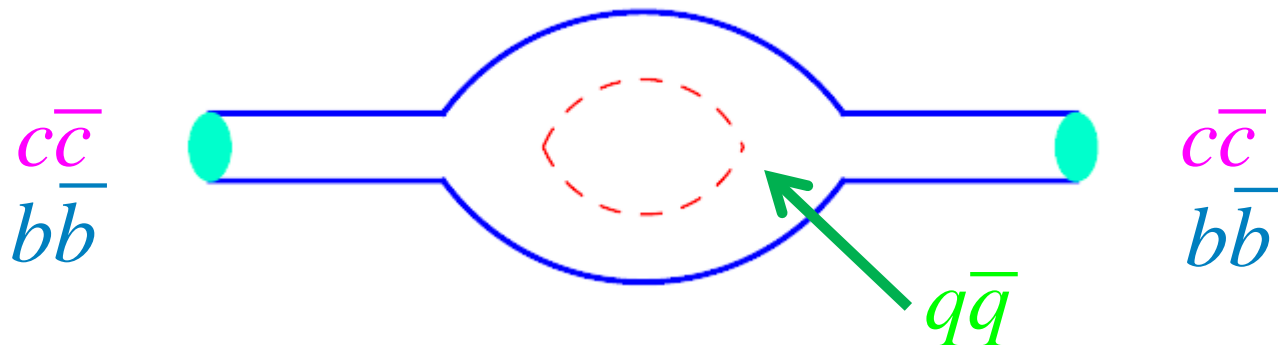
Tokai, 2014.12.02

Outline

- Motivation
- Coupled-channel effects and threshold enhancement phenomena
 - ✓ Dipion transitions $e^+e^- \rightarrow J/\psi\pi\pi, \psi'\pi\pi, h_c\pi\pi$
 - ✓ Zb production, hunting for some XYZ particles
 - ✓ Anomalous threshold singularity
- Summary

Motivation

- ✓ Observations of XYZ particles;
- ✓ Discrepancy between conventional quenched quark model predictions and experimental data;
- ✓ Coupled-channel effects will largely affect the mass and decay properties of heavy quarkonia;



Coupled-channel effects

- ✓ E. Eichten et al, PRD17,3090(1978), PRD73,104014(2006);
- ✓ M.R. Pennington, D.J. Wilson, PRD76,077502(2008);
- ✓ T. Barnes, E.S. Swanson, PRC77,055206(2008);
- ✓ B.Q. Li, C.Meng, K.T. Chao, PRD80,014012(2009);
- ✓ F.K. Guo et al; PRD83,034013(2011);
- ✓ Z.Y. Zhou, Z. Xiao, arXiv:1309.1949;
- ✓

Unconventional states in heavy quarkonium region

| State | m (MeV) | Γ (MeV) | J^{PC} | Process (mode) | Experiment ($\# \sigma$) | Year | Status |
|------------------|------------------------|-----------------------|-----------------|--|--|------|--------|
| $X(3872)$ | 3871.68 ± 0.17 | < 1.2 | $1^{++}/2^{-+}$ | $B \rightarrow K(\pi^+\pi^-J/\psi)$ $p\bar{p} \rightarrow (\pi^+\pi^-J/\psi) + \dots$ $B \rightarrow K(\omega J/\psi)$ $B \rightarrow K(D^{*0}\bar{D}^0)$ $B \rightarrow K(\gamma J/\psi)$ $B \rightarrow K(\gamma\psi(2S))$ $pp \rightarrow (\pi^+\pi^-J/\psi) + \dots$ | Belle [36, 37] (12.8), BABAR [38] (8.6) CDF [39–41] (np), DØ [42] (5.2) Belle [43] (4.3), BABAR [23] (4.0) Belle [44, 45] (6.4), BABAR [46] (4.9) Belle [47] (4.0), BABAR [48, 49] (3.6) BABAR [49] (3.5), Belle [47] (0.4) LHCb [50] (np) | 2003 | OK |
| $X(3915)$ | 3917.4 ± 2.7 | 28_{-9}^{+10} | $0/2^{?+}$ | $B \rightarrow K(\omega J/\psi)$ $e^+e^- \rightarrow e^+e^-(\omega J/\psi)$ | Belle [51] (8.1), BABAR [52] (19) Belle [53] (7.7), BABAR [23] (np) | 2004 | OK |
| $X(3940)$ | 3942_{-8}^{+9} | 37_{-17}^{+27} | $?^{?+}$ | $e^+e^- \rightarrow J/\psi(D\bar{D}^*)$ $e^+e^- \rightarrow J/\psi(\dots)$ | Belle [54] (6.0) Belle [20] (5.0) | 2007 | NC! |
| $G(3900)$ | 3943 ± 21 | 52 ± 11 | 1^{--} | $e^+e^- \rightarrow \gamma(D\bar{D})$ | BABAR [55] (np), Belle [56] (np) | 2007 | OK |
| $Y(4008)$ | 4008_{-49}^{+121} | 226 ± 97 | 1^{--} | $e^+e^- \rightarrow \gamma(\pi^+\pi^-J/\psi)$ | Belle [57] (7.4) | 2007 | NC! |
| $Z_1(4050)^+$ | 4051_{-43}^{+24} | 82_{-55}^{+51} | $?$ | $B \rightarrow K(\pi^+\chi_{c1}(1P))$ | Belle [58] (5.0), BABAR [59] (1.1) | 2008 | NC! |
| $\Upsilon(4140)$ | 4143.4 ± 3.0 | 15_{-7}^{+11} | $?^{?+}$ | $B \rightarrow K(\phi J/\psi)$ | CDF [60, 61] (5.0) | 2009 | NC! |
| $\Upsilon(4160)$ | 4156_{-25}^{+29} | 139_{-65}^{+113} | $?^{?+}$ | $e^+e^- \rightarrow J/\psi(D\bar{D}^*)$ | Belle [54] (5.5) | 2007 | NC! |
| $Z_2(4250)^+$ | 4248_{-45}^{+185} | 177_{-72}^{+321} | $?$ | $B \rightarrow K(\pi^+\chi_{c1}(1P))$ | Belle [58] (5.0), BABAR [59] (2.0) | 2008 | NC! |
| $Y(4260)$ | 4263_{-9}^{+8} | 95 ± 14 | 1^{--} | $e^+e^- \rightarrow \gamma(\pi^+\pi^-J/\psi)$ $e^+e^- \rightarrow (\pi^+\pi^-J/\psi)$ $e^+e^- \rightarrow (\pi^0\pi^0J/\psi)$ | BABAR [62, 63] (8.0) CLEO [64] (5.4), Belle [57] (15) CLEO [65] (11) CLEO [65] (5.1) | 2005 | OK |
| $Y(4274)$ | $4274.4_{-6.7}^{+8.4}$ | 32_{-15}^{+22} | $?^{?+}$ | $B \rightarrow K(\phi J/\psi)$ | CDF [61] (3.1) | 2010 | NC! |
| $\Upsilon(4350)$ | $4350.6_{-4.6}^{+4.6}$ | $19.9_{-1.4}^{+18.4}$ | $0/2^{++}$ | $e^+e^- \rightarrow e^+e^-(\psi(4S))$ | Belle [66] (3.2) | 2009 | NC! |
| $Y(4360)$ | 4361 ± 13 | 74 ± 18 | 1^{--} | $e^+e^- \rightarrow \gamma(\pi^+\pi^-\psi(2S))$ | BABAR [67] (np), Belle [68] (8.0) | 2007 | OK |
| $Z(4430)^+$ | 4443_{-18}^{+24} | 107_{-71}^{+113} | $?$ | $B \rightarrow K(\pi^+\psi(2S))$ | Belle [69, 70] (6.4), BABAR [71] (2.4) | 2007 | NC! |
| $X(4630)$ | 4634_{-11}^{+9} | 92_{-32}^{+41} | 1^{--} | $e^+e^- \rightarrow \gamma(\Lambda_c^+\Lambda_c^-)$ | Belle [72] (8.2) | 2007 | NC! |
| $Y(4660)$ | 4664 ± 12 | 48 ± 15 | 1^{--} | $e^+e^- \rightarrow \gamma(\pi^+\pi^-\psi(2S))$ | Belle [68] (5.8) | 2007 | NC! |
| $Z_b(10610)^+$ | 10607.2 ± 2.0 | 18.4 ± 2.4 | 1^+ | $\Upsilon(5S) \rightarrow \pi^-(\pi^+[b\bar{b}])$ | Belle [73, 74] (16) | 2011 | NC! |
| $Z_b(10650)^+$ | 10652.2 ± 1.5 | 11.5 ± 2.2 | 1^+ | $\Upsilon(5S) \rightarrow \pi^-(\pi^+[b\bar{b}])$ | Belle [73, 74] (16) | 2011 | NC! |
| $Y_b(10888)$ | 10888.4 ± 3.0 | $30.7_{-7.7}^{+8.9}$ | 1^{--} | $e^+e^- \rightarrow (\pi^+\pi^-\Upsilon(nS))$ | Belle [75, 76] (2.0) | 2010 | NC! |

Coupled-channel effects with P-wave states involved

- ✓ Combinations of S- and P-wave charmed mesons are very close to some conventional higher charmonia($\psi(4160)$, $\psi(4415)$) and $Y(4260)$, $Y(4360)$, $Z(4430)$ );
- ✓ The coupling with the parity-odd charmonia could be S-wave, supposed to be strong;

| | $D_0 D^*$ | $D_1' D$ | $D_1' D^*$ | $D_1 D$ | $D_1 D^*$ | $D_2 D$ | $D_2 D^*$ | $D_{s0} D_s^*$ | $D_{s1} D_s$ |
|-----------------|-----------|----------|------------|---------|-----------|---------|-----------|----------------|--------------|
| Threshold [MeV] | 4325 | 4292 | 4434 | 4286 | 4428 | 4327 | 4470 | 4430 | 4424 |

Connections between coupled channel effects and XYZ?

Model based on HHChPT

Doublets with light degrees of freedom $j^P=1/2^-,1/2^+,3/2^+$

$$\begin{aligned} H_a &= \frac{1 + \not{v}}{2} [\mathcal{D}_{a\mu}^* \gamma^\mu - \mathcal{D}_a \gamma_5] , \\ S_a &= \frac{1 + \not{v}}{2} [\mathcal{D}_{1a}^{\prime\mu} \gamma_\mu \gamma_5 - \mathcal{D}_{0a}^*] , \\ T_a^\mu &= \frac{1 + \not{v}}{2} \left\{ \mathcal{D}_{2a}^{\mu\nu} \gamma_\nu \right. \\ &\quad \left. - \sqrt{\frac{3}{2}} \mathcal{D}_{1a\nu} \gamma_5 \left[g^{\mu\nu} - \frac{1}{3} \gamma^\nu (\gamma^\mu - v^\mu) \right] \right\} , \end{aligned}$$

HQSS allowed coupling (LDOF will also be conserved)

| | HH | SH | TH |
|------------|--------|--------|--------|
| $\psi(nS)$ | P-wave | S-wave | D-wave |
| $\psi(nD)$ | P-wave | D-wave | S-wave |

Leading Order Effective Lagrangian

According to HHChPT power counting

$$\mathcal{L}_1 = \frac{g_T}{\sqrt{2}} < J^{\mu\nu} \bar{H}_a^\dagger \gamma_\nu \bar{T}_{a\mu} - J^{\mu\nu} \bar{T}_{a\mu}^\dagger \gamma_\nu \bar{H}_a >$$

$$+ ig_H < J^{\mu\nu} \bar{H}_a^\dagger \gamma_\mu \overleftrightarrow{\partial}_\nu \bar{H}_a >$$

$$+ g_S < J \bar{S}_a^\dagger \bar{H}_a + J \bar{H}_a^\dagger \bar{S}_a >$$

$$+ C_S < J \bar{H}_b^\dagger \gamma_\mu \gamma_5 \bar{H}_a \mathcal{A}_{ba}^\mu >$$

$$+ iC_P < J^\mu \bar{H}_b^\dagger \sigma_{\mu\nu} \gamma_5 \bar{H}_a \mathcal{A}_{ba}^\nu > + h.c.,$$

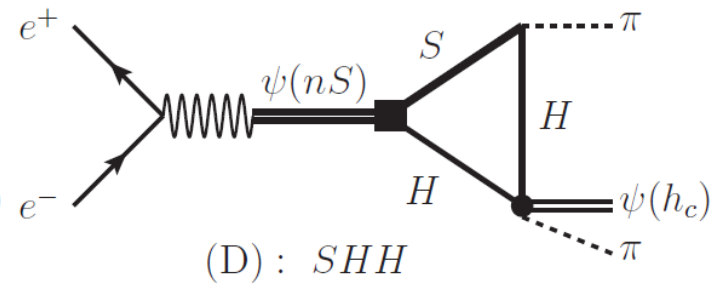
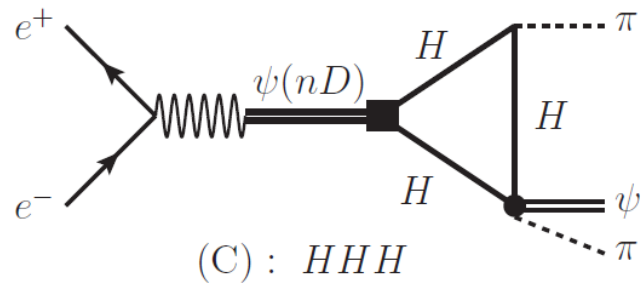
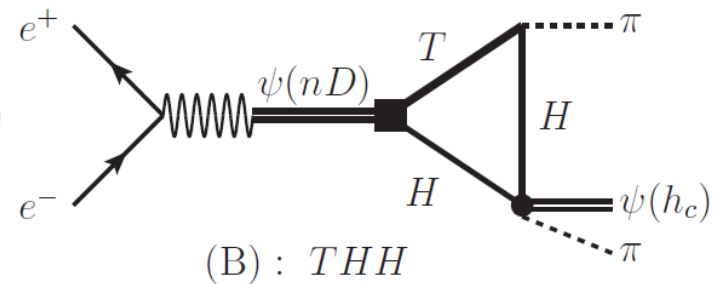
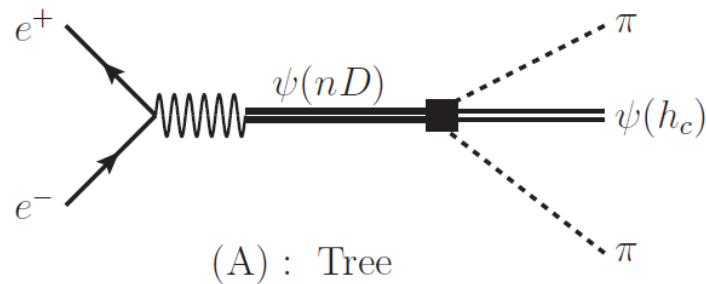
$$\mathcal{L}_2 = i \frac{h'}{\Lambda_\chi} < \bar{H}_a T_b^\mu \gamma^\nu \gamma_5 (D_\mu \mathcal{A}_\nu + D_\nu \mathcal{A}_\mu)_{ba} >$$

$$+ ih < \bar{H}_a S_b \gamma_\mu \gamma_5 \mathcal{A}_{ba}^\mu > + ig < H_b \gamma_\mu \gamma_5 \mathcal{A}_{ba}^\mu \bar{H}_a > .$$

$$J = \frac{1 + \not{v}}{2} [\psi(nS)^\mu \gamma_\mu] \frac{1 - \not{v}}{2}, J^\mu = \frac{1 + \not{v}}{2} [h_c(nP)^\mu \gamma_5] \frac{1 - \not{v}}{2}$$

$$J^{\mu\nu} = \frac{1 + \not{v}}{2} \left\{ \psi(nD)_\alpha \left[\frac{1}{2} \sqrt{\frac{3}{5}} [(\gamma^\mu - v^\mu) g^{\alpha\nu} + (\gamma^\nu - v^\nu) g^{\alpha\mu}] - \sqrt{\frac{1}{15}} (g^{\mu\nu} - v^\mu v^\nu) \gamma^\alpha \right] \right\} \frac{1 - \not{v}}{2}$$

Dipion Transitions



$\psi(4160)$ as the input $\psi(nD)$:

- ✓ Widely accepted as a conventional 2^3D_1 charmonia
- ✓ Couple to TH via S-wave, respect HQSS
- ✓ Close to $Y(4260)$

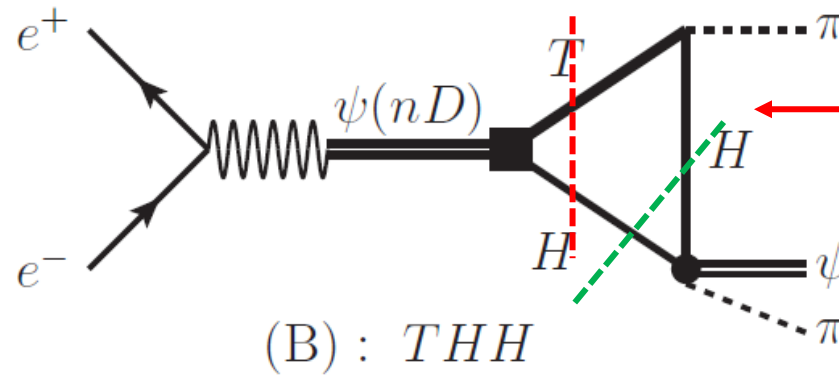
$M=4153\pm 3$ MeV, $\Gamma=103\pm 8$ MeV **PDG averaged**

$M=4191.7\pm 6.5$ MeV, $\Gamma=71.8\pm 12.3$ MeV **BES, PLB660,315(2008)**

$M=4193\pm 7$ MeV, $\Gamma=79\pm 14$ MeV **X.H. Mo et al, PRD82,077501(2010)**

$\Gamma_{ee}=0.83\pm 0.07$ KeV, not small

THH Loop



Triangle singularity (TS)
may occur under special
kinematic configurations

$$\text{I) } \{D_1 D [D^*]\},$$

$$\text{II) } \{D_1 D^* [D^*]\},$$

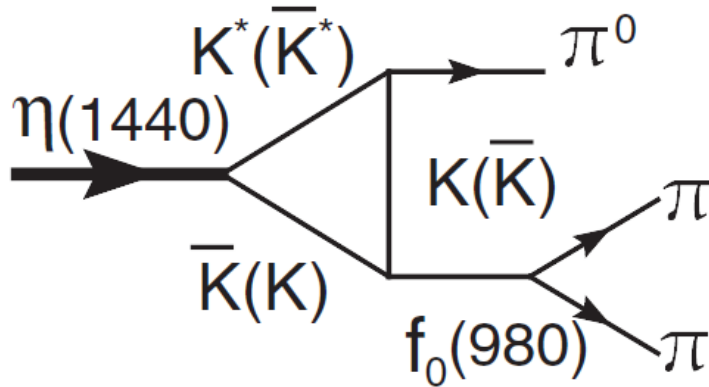
$$\text{III) } \{D_2 D^* [D]\},$$

$$\text{IV) } \{D_2 D^* [D^*]\},$$

In the heavy quark limit

$$\mathcal{M}^I : \mathcal{M}^{II} : \mathcal{M}^{III} : \mathcal{M}^{IV} = 1 : \frac{1}{2} : -\frac{1}{5} : \frac{3}{10} .$$

Triangle Singularity



**Largely isospin violation in $\eta(1405/1475) \rightarrow 3\pi$
Br~10%, [BESIII, PRL108,182001 (2012)]**

Wu,Liu,Zhao&Zou, PRL108,081803(2012)

References:

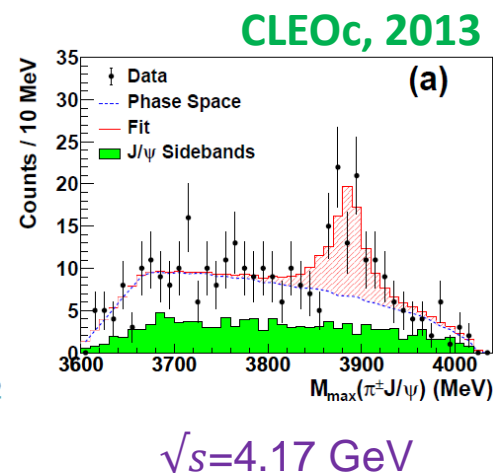
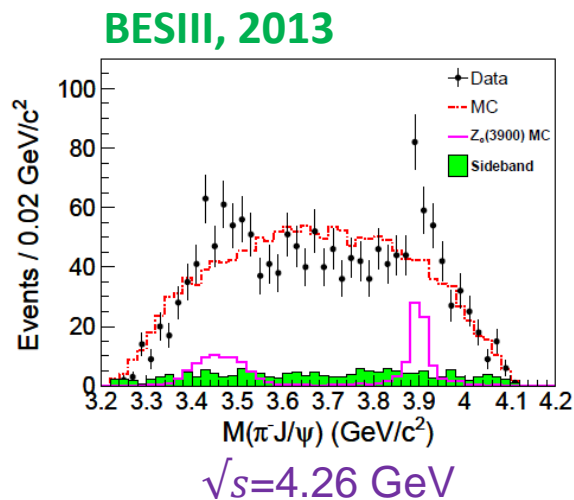
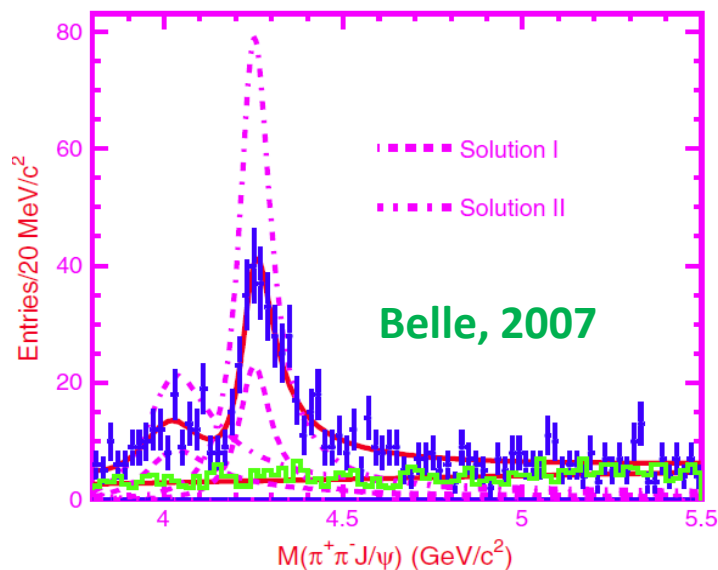
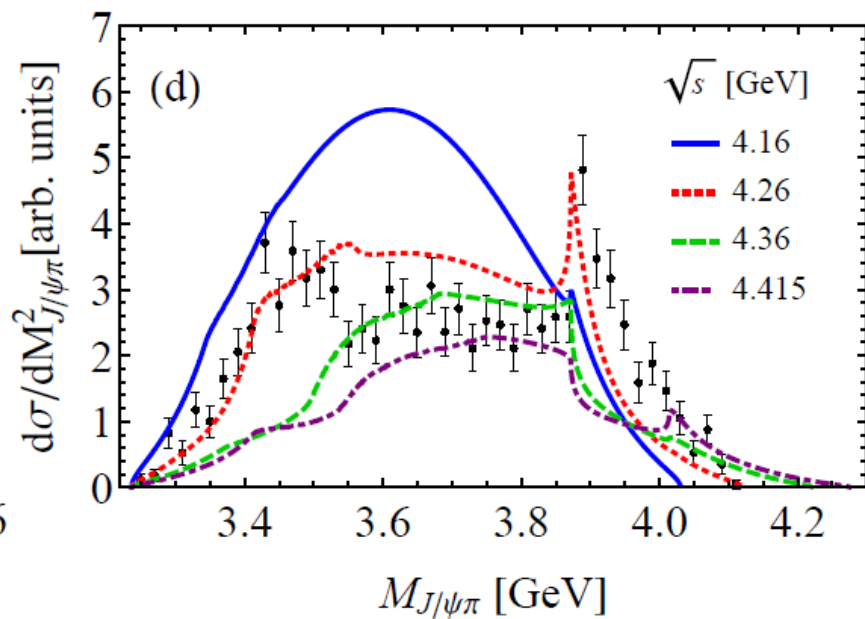
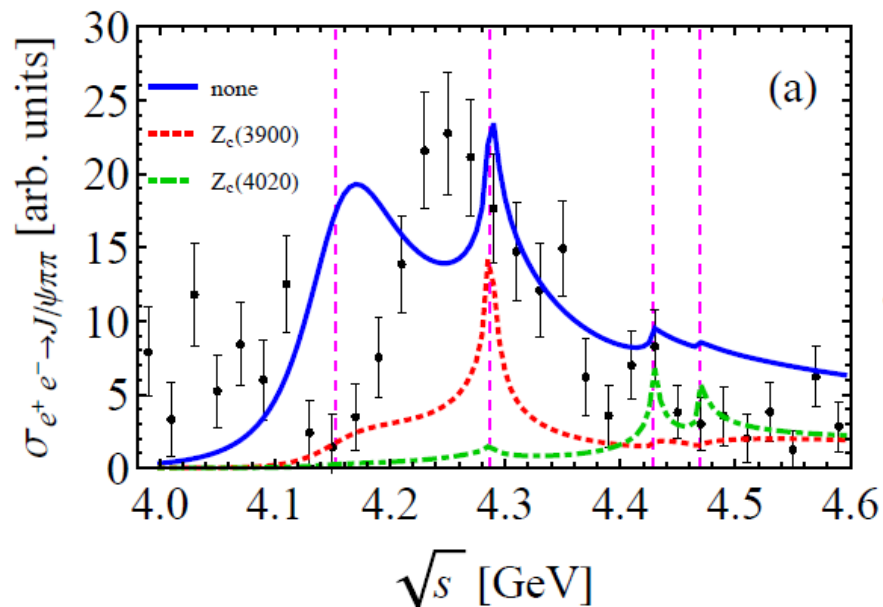
Landshoff and Treiman, Phys.Rev. 127,649(1962)

Landshoff and Treiman, Nuovo Cimento 19,1249(1961)

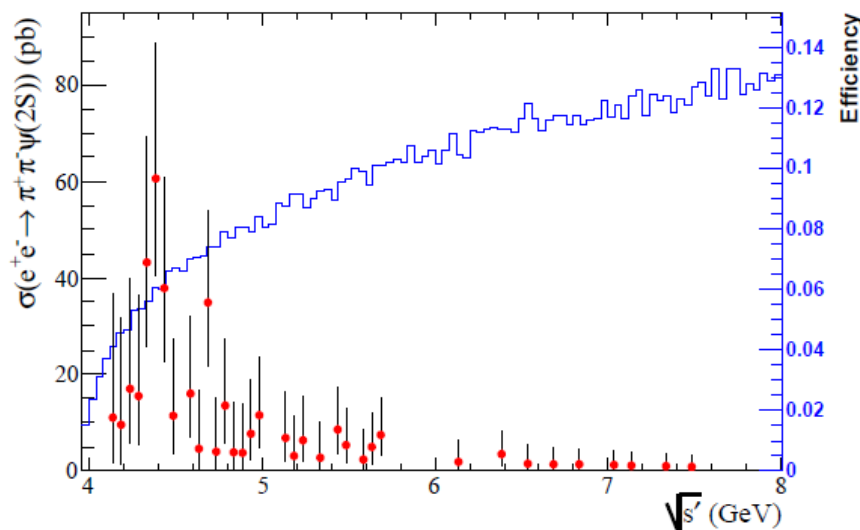
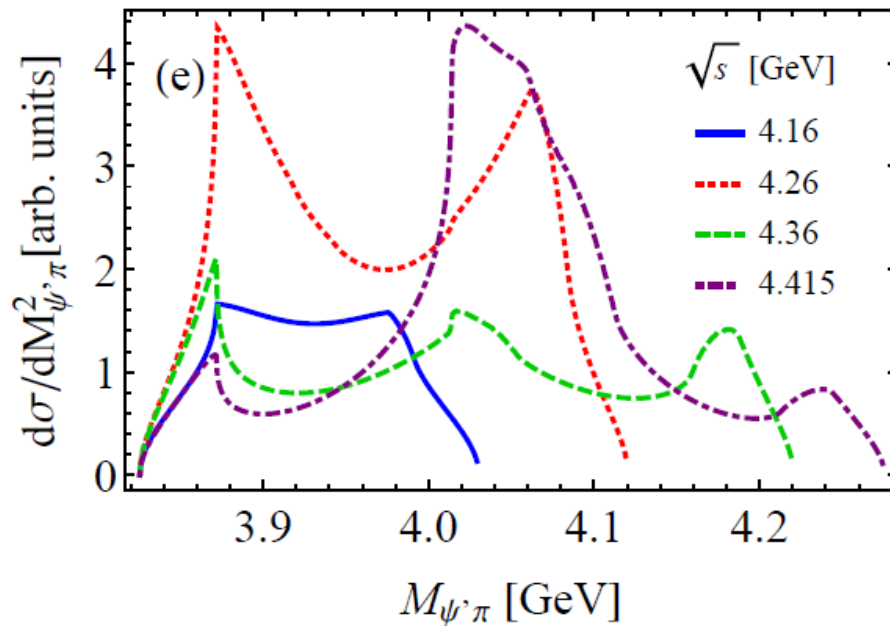
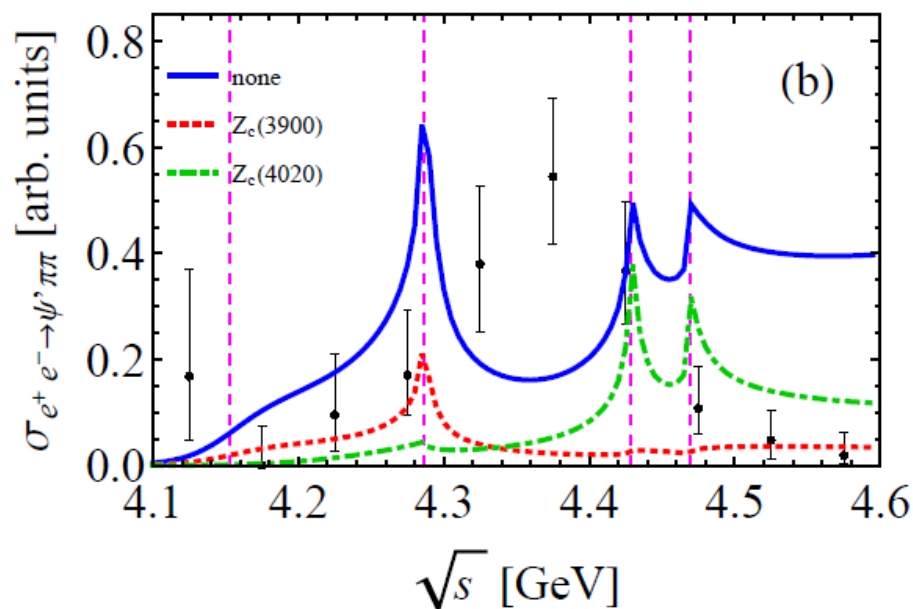
Eden et al., <<The Analytic S-Matrix>>, 1966

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$e^+e^- \rightarrow J/\psi\pi\pi$ via $\psi(4160)$ and THH loops

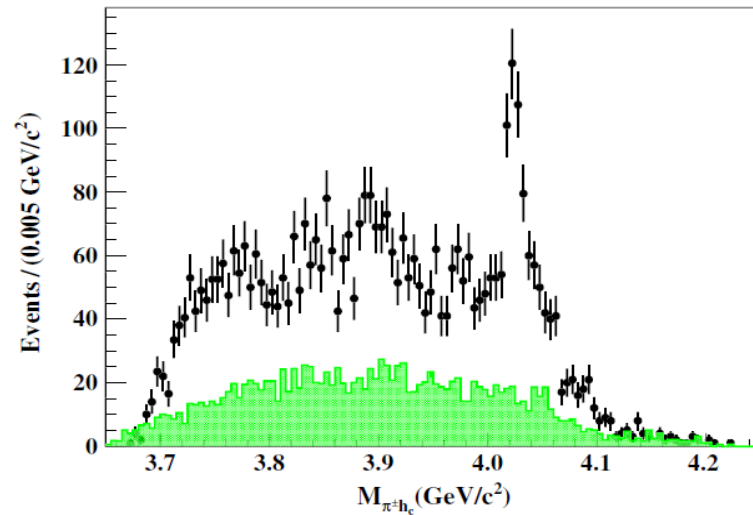
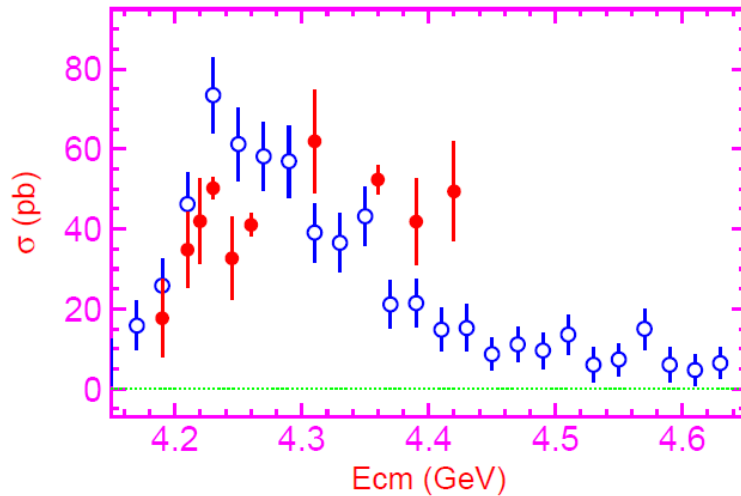
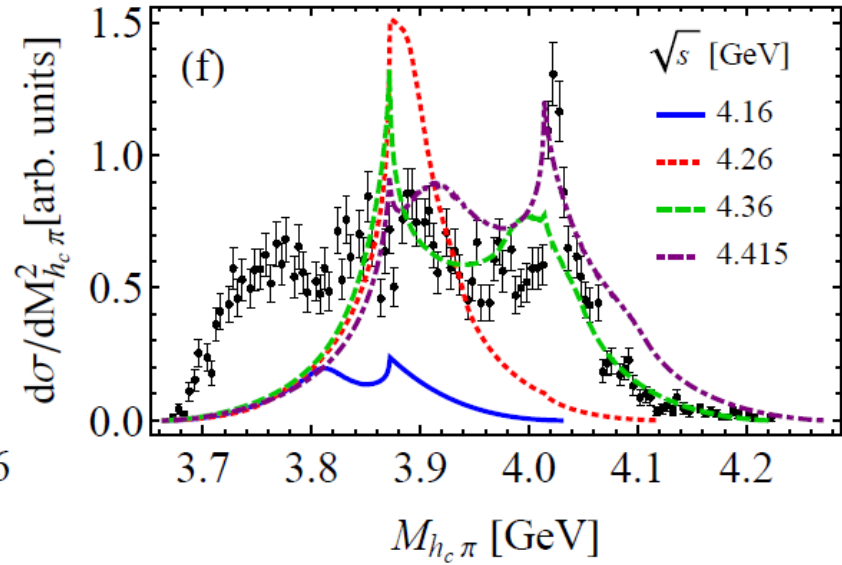
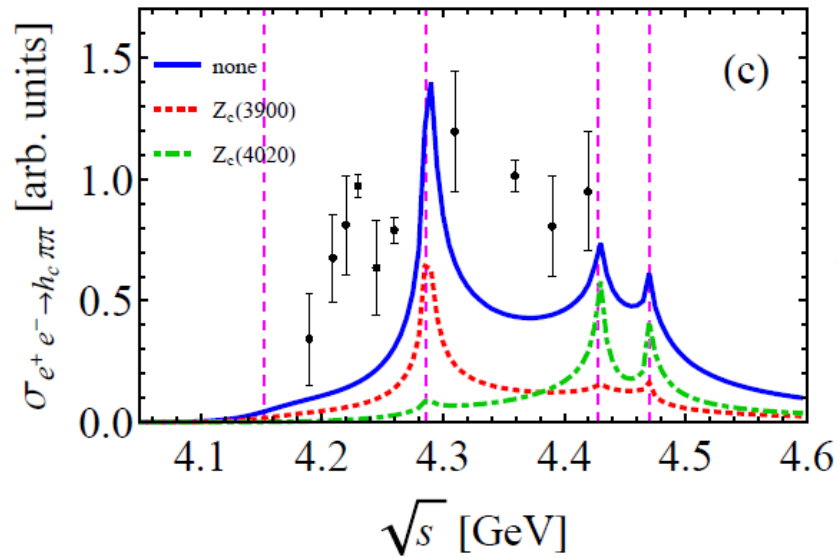


$e^+e^- \rightarrow \psi'\pi\pi$ via $\psi(4160)$ and THH loops



- ✓ Results are sensitive to kinematics
- ✓ Same dynamics, different kinematics
- ✓ Direct prediction to check the mechanism

$e^+e^- \rightarrow h_c \pi\pi$ via $\psi(4160)$ and THH loops



Sum of energy points from 3.9 to 4.42

C.Z. Yuan, arXiv:1312.6399

Belle, PRL111,242001(2013)

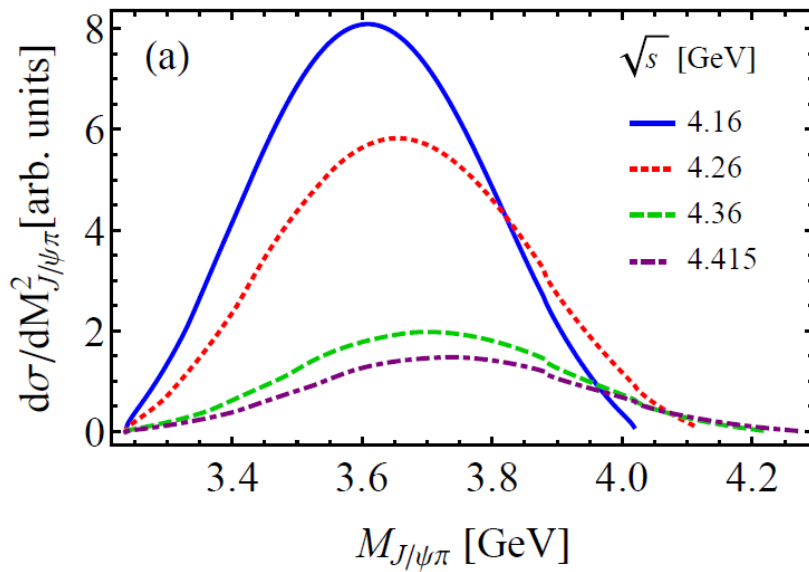
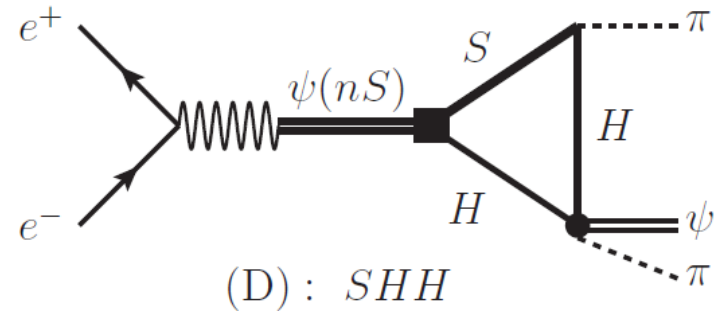
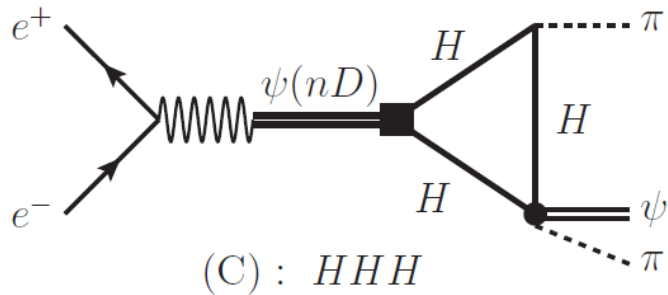
Comparison with Molecule Ansatz

- ✓ Similar points: kinematics, singularities of rescattering loops
- ✓ Different points:
 - Incorporates the D_1D , D_1D^* , D_2D^* combinations in a single Lagrangian with the relative phase and coupling strength fixed in the heavy quark limit;
 - No matter whether the molecular state exist or not, it seems to be natural to suppose the coupled channel effects should exist physically

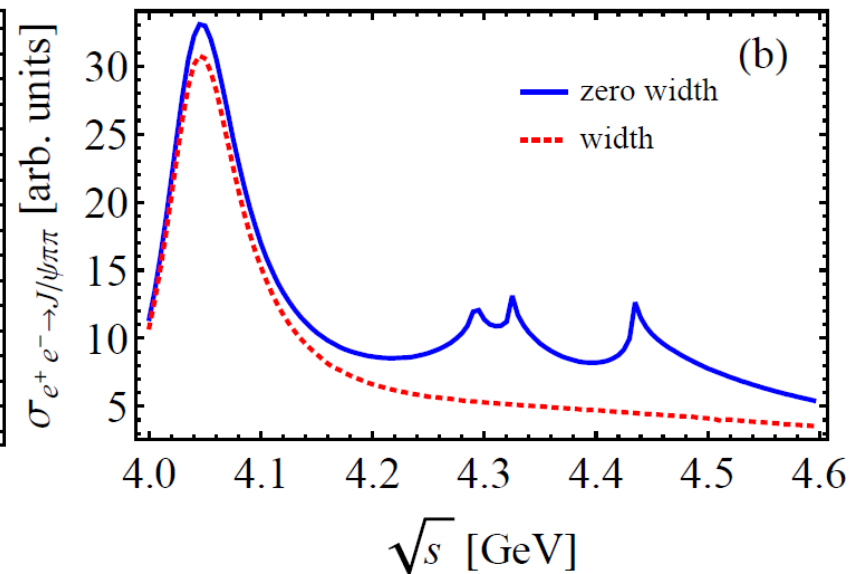
Y(4260) as molecular state:

- ✓ D_1D , D_0D^* molecular state, F. Close et al, PRL102,242003(2009), PRD81,074033(2010); Kalashnikova and Nefediev, PRD77,054025(2008);
- ✓ Potential model, G.J. Ding PRD79,014001(2009)
- ✓ Connection with Zc(3900), Wang, Zhao and Hanhart, PRL111,132003(2013); X.H. Liu and G.Li, PRD88,014013(2013)
- ✓

HHH Loop and SHH Loop

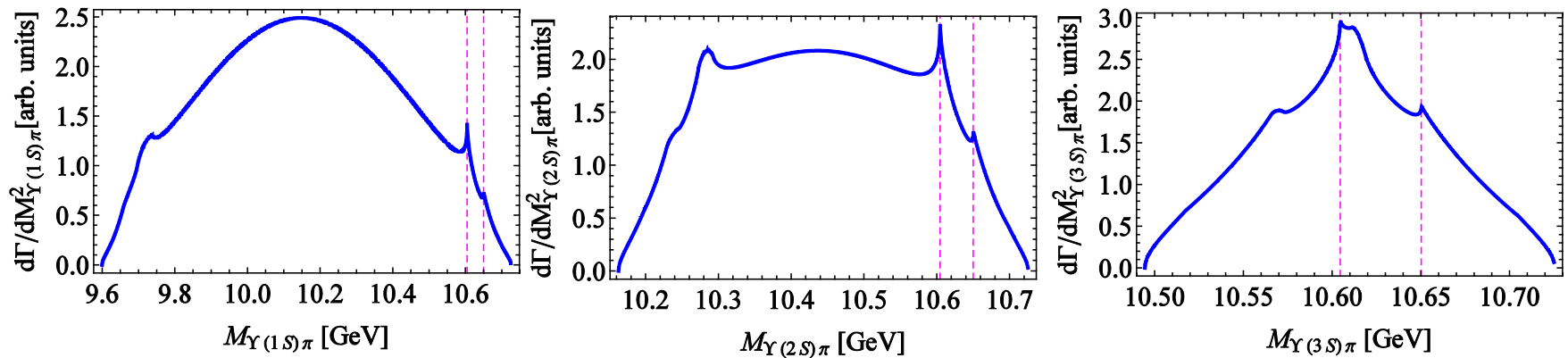
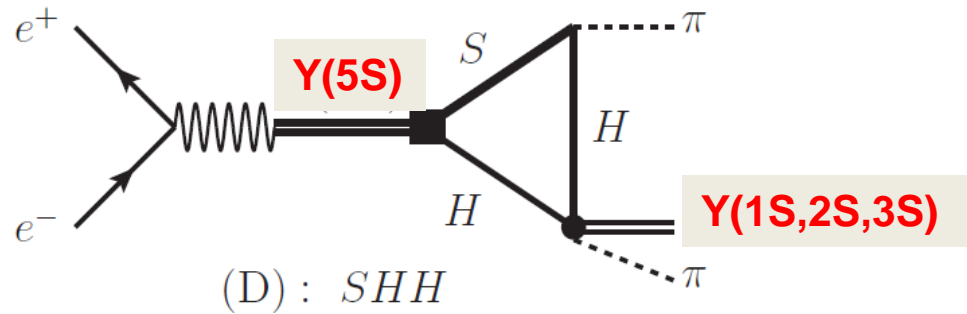


No obvious cusp around 3.9 GeV,
inconsistent with CLEOc result



Broad width of D_0 and D_1' will lower
the amplitude and smooth the cusps

Zb Production



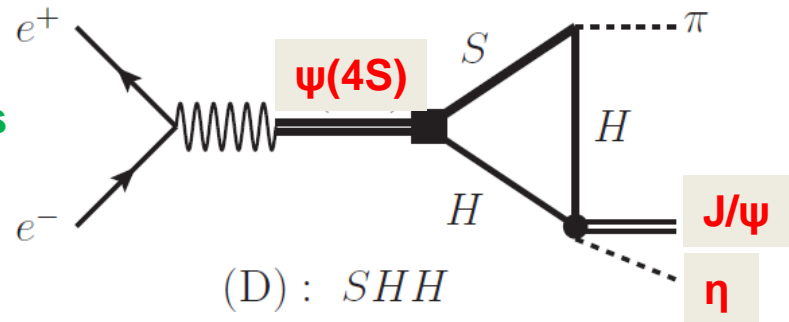
Cusps at BB^* and B^*B^* thresholds



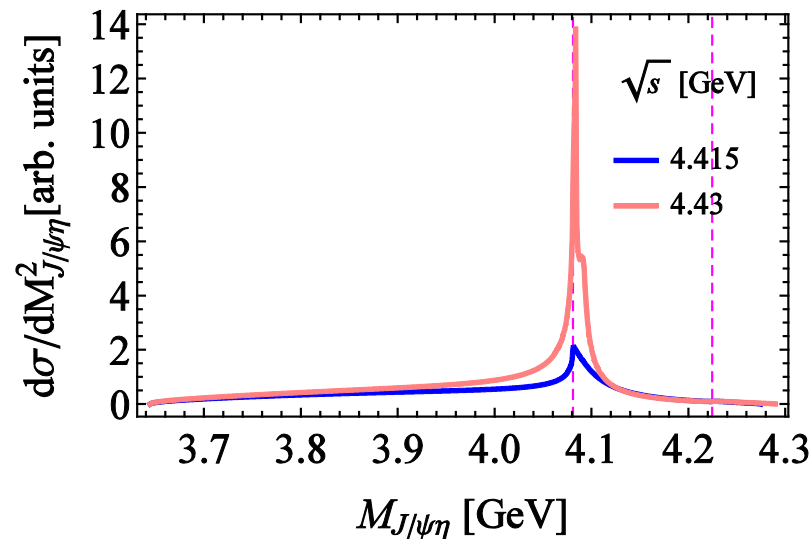
Production of $Z_b(10610)$ and $Z_b(10650)$

Hunting For Partners of Y4160 & Y4274

Isospin violation processes



Preliminary



Cusps at $D_s D_s^*$ and $D_s^* D_s$ thresholds
($C=-1$)



Production of partners of
Y(4160) and Y(4274)
(observed in $J/\psi\phi$, $C=+1$)

Anomalous Threshold Singularity

✓ Singularity in the complex space

Landau Equation

$$I_3 = \prod_{i=1}^3 \int_0^1 da_i \frac{\delta(1 - \sum_k a_k)}{D - i\epsilon}$$

$$D = \sum_{i,j} a_i a_j Y_{ij}, \quad Y_{ij} = \frac{1}{2} [m_i^2 + m_j^2 - (q_{i-1} - q_{j-1})^2]$$

Necessary conditions

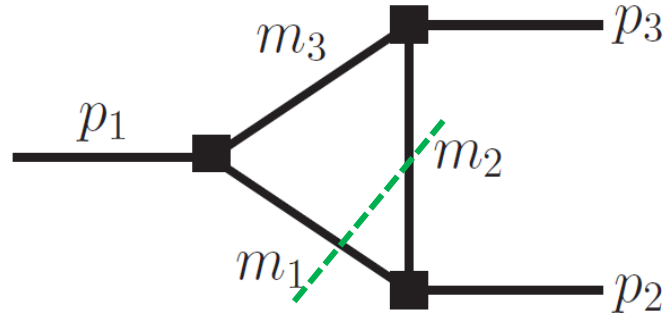
$$D = 0,$$

$$\text{either } a_j = 0 \text{ or } \frac{\partial D}{\partial a_j} = 0.$$

Leading singularity



Anomalous Threshold Singularity



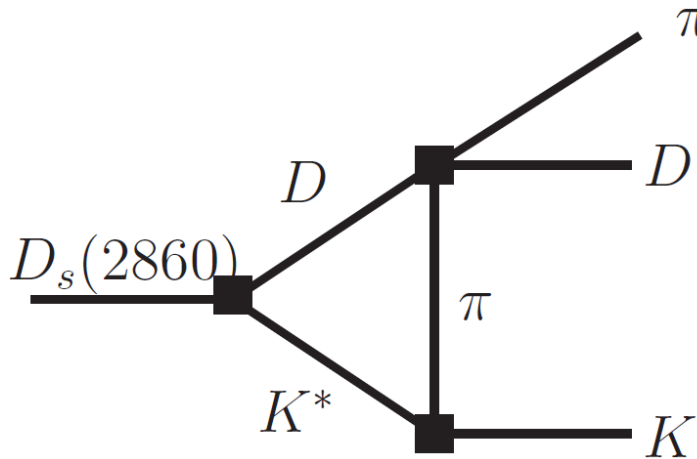
$$p_1^2 = W^2, p_2^2 = s_2, p_3^2 = m^2.$$

Normal threshold $s_n = (m_1 + m_2)^2$

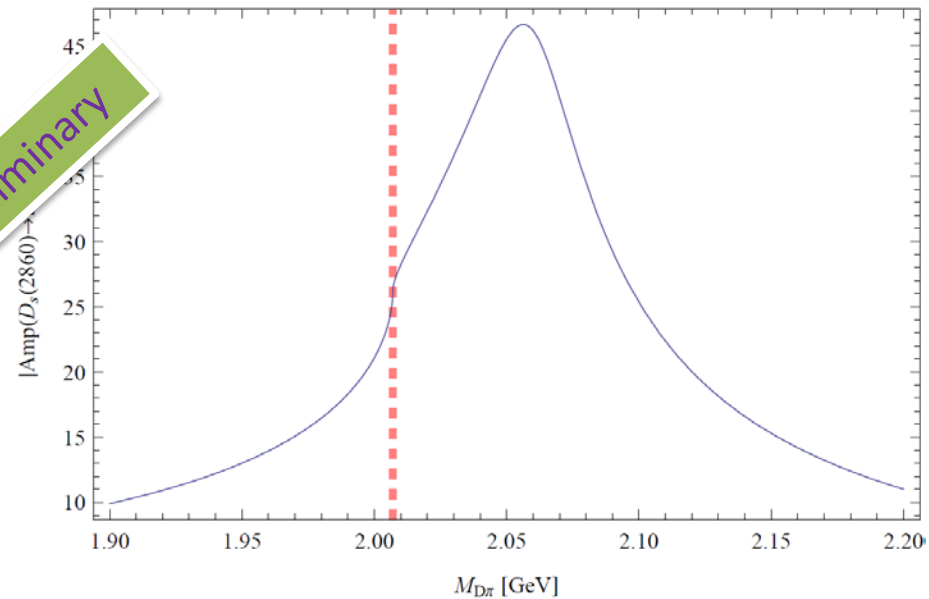
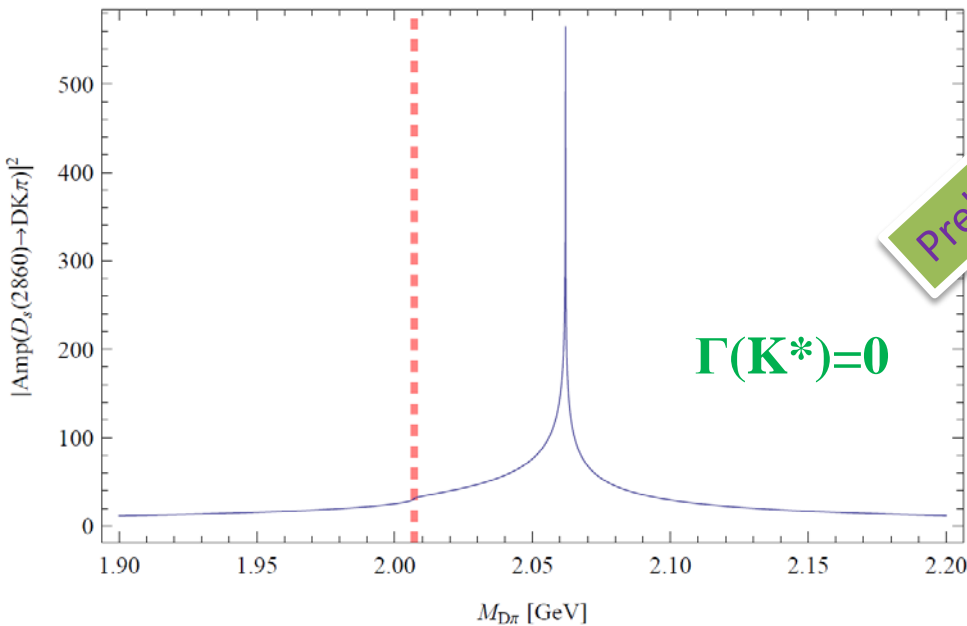
Anomalous threshold $\bar{s}_2 = s_n + \frac{m_1}{m_3}[(m_2 - m_3)^2 - m^2],$

when $W = m_1 + m_3,$

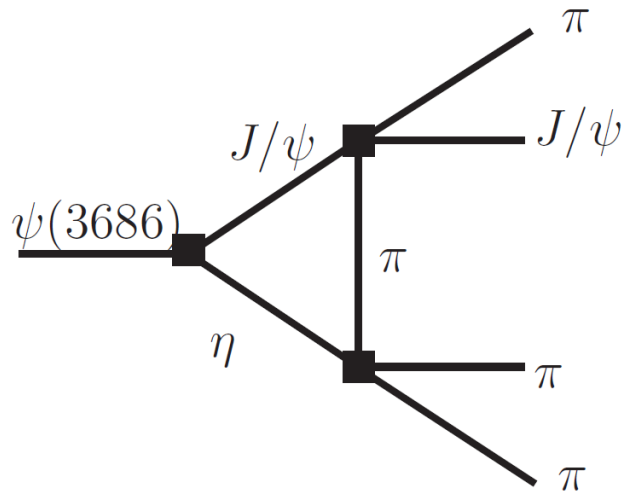
Anomalous Threshold Singularity



$$D_s(2860) \rightarrow DK^*[\pi] \rightarrow DK\pi$$

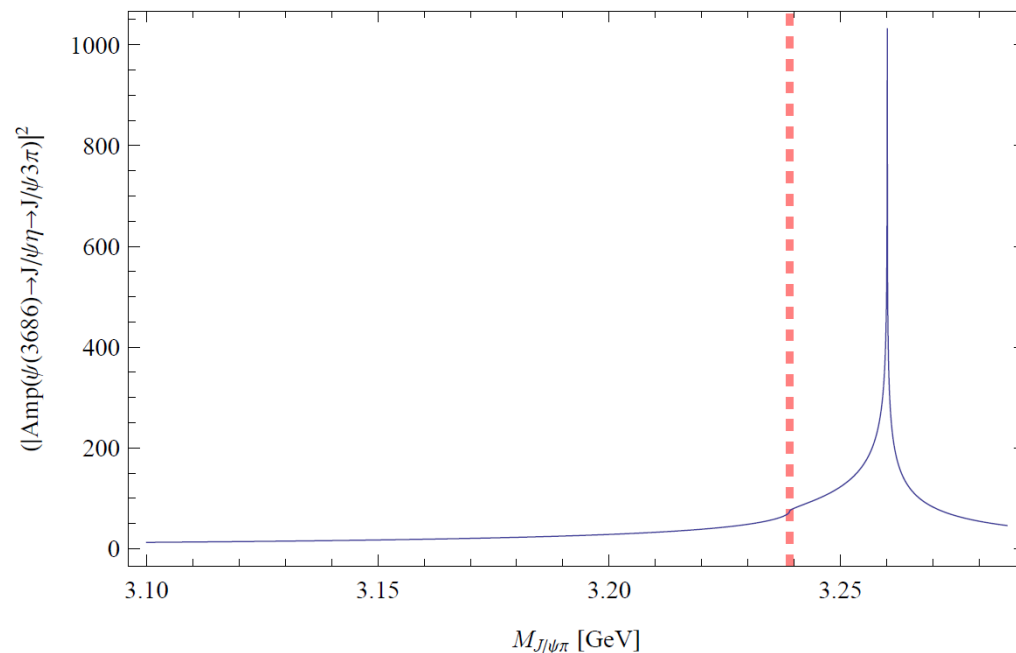


Anomalous Threshold Singularity



For investigating $J/\psi\pi$ interaction

Preliminary



Summary

- The lineshape behavior of the cross sections and distributions for the dipion transitions are studied. Couple-channel effects may largely affect the threshold behavior, especially that induced by the couplings between D-wave charmonia and TH charmed mesons, taking into account these leading order S-wave couplings will respect HQSS.
- Some interesting cusps are obtained, which may have some underlying connections with the XYZ states observed around the TH and HH thresholds.
- Kinematics plays a crucial role in generating the cusps.
- Anomalous threshold singularity would be used to discriminate coupled-channel effects and genuine resonances.