

**Dynamical coupled-channels study of
 $S=-1$ hyperon resonances
and
possible new experiments at J-PARC**

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(RCNP, Osaka U.)**

**International Workshop on “J-PARC Hadron Physics 2016”
Tokai, March 2th-4th, 2016**

Current situation of $Y^*(= \Lambda^*, \Sigma^*)$ spectroscopy

$Y^* (= \Lambda^*, \Sigma^*)$ resonances are much less understood than N^* & Δ^* !!

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Λ^*			Σ^*		
Particle	J^P	Overall status	Particle	J^P	Overall status
$\Lambda(1116)$	1/2+	****	$\Sigma(1193)$	1/2+	****
$\Lambda(1405)$	1/2-	****	$\Sigma(1385)$	3/2+	****
$\Lambda(1520)$	3/2-	****	$\Sigma(1480)$		*
$\Lambda(1600)$	1/2+	***	$\Sigma(1560)$		**
$\Lambda(1670)$	1/2-	****	$\Sigma(1580)$	3/2-	*
$\Lambda(1690)$	3/2-	****	$\Sigma(1620)$	1/2-	**
$\Lambda(1800)$	1/2-	***	$\Sigma(1660)$	1/2+	***
$\Lambda(1810)$	1/2+	***	$\Sigma(1670)$	3/2-	****
$\Lambda(1820)$	5/2+	****	$\Sigma(1690)$		**
$\Lambda(1830)$	5/2-	****	$\Sigma(1750)$	1/2-	***
$\Lambda(1890)$	3/2+	****	$\Sigma(1770)$	1/2+	*
$\Lambda(2000)$		*	$\Sigma(1775)$	5/2-	****
$\Lambda(2020)$	7/2+	*	$\Sigma(1840)$	3/2+	*
$\Lambda(2100)$	7/2-	****	$\Sigma(1880)$	1/2+	**
$\Lambda(2110)$	5/2+	***	$\Sigma(1915)$	5/2+	****
$\Lambda(2325)$	3/2-	*	$\Sigma(1940)$	3/2-	***
$\Lambda(2350)$		***	$\Sigma(2000)$	1/2-	*
$\Lambda(2585)$		**	$\Sigma(2030)$	7/2+	****
			$\Sigma(2070)$	5/2+	*
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			$\Sigma(2100)$	7/2-	*
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			$\Sigma(2455)$		**
			$\Sigma(2620)$		**
			$\Sigma(3000)$		*
			$\Sigma(3170)$		*

PDG listing

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$\Lambda(2350)$	***	***	$\Sigma(2000)$	$1/2^-$	*
$\Lambda(2585)$	**	**	$\Sigma(2030)$	$7/2^+$	****
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- ✓ Before 2012, PDG listed only **Breit-Wigner (BW) mass and width** [except $\Sigma(1385)3/2^+$, $\Lambda(1520)3/2^-$]
 - N^* & Δ^* case:
Resonances **defined by poles of scattering amplitudes** are extensively studied;
PDG lists **BOTH pole and BW parameters**.

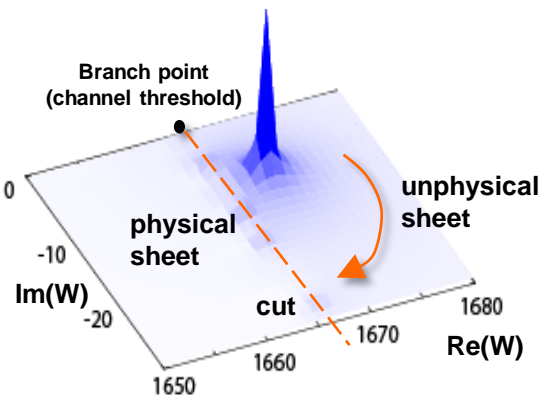
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Current situation of $Y^*(= \Lambda^*, \Sigma^*)$ spectroscopy

Resonance theory based on Gamow vectors:
[G. Gamow (1928), R. E. Peierls (1959), ...]

“Quantum resonance state is an (complex-)energy eigenstate of the **FULL** Hamiltonian of the **underlying theory** imposed by Purely Outgoing Boundary Condition (POBC).”



Energy eigenvalue = Pole energy
Transition matrix elements ~ Residues^{1/2} at the pole
btwn res. & scatt. states

Extracting poles of amplitudes from reaction data is nothing but obtaining “exact” energy eigenvalues of QCD !!!

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$\Lambda(2350)$		***
$\Lambda(2585)$		**
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- ✓ Comprehensive partial-wave analyses of $K^- p$ reactions to extract Y^* **defined by poles** have been accomplished **just recently** :

- Kent State University (KSU) group
(→ 2013, “KSU on-shell parametrization” of S-matrix)
Zhang et al., PRC88(2013)035204, 035205.
→ Reanalysis of KSU single-energy solution using an on-shell K-matrix model (Fernandez-Ramirez et al., arXiv:1510.07065)
- Our group
(→ 2014-2015, dynamical coupled-channels approach)
HK, Nakamura, Lee, Sato, PRC90(2014)065204; 92(2015)025205

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Dynamical Coupled-Channels (DCC) approach to Λ^* & Σ^* productions

Dynamical Coupled-Channels (DCC) model:

[Matsuyama, Sato, Lee, PR439(2007)193; HK, Nakamura, Lee, Sato, PRC88(2013)035209;90(2014)065204]

$$T_{a,b}^{(LSJ)}(p_a, p_b; E) = V_{a,b}^{(LSJ)}(p_a, p_b; E) + \underbrace{\sum_c}_{\text{CC effect}} \underbrace{\int_0^\infty q^2 dq}_{\text{off-shell effect}} V_{a,c}^{(LSJ)}(p_a, q; E) G_c(q; E) T_{c,b}^{(LSJ)}(q, p_b; E)$$

$$a, b, c = (\bar{K}N, \pi\Sigma, \pi\Lambda, \eta\Lambda, K\Xi, \boxed{\pi\Sigma^*, \bar{K}^*N}, \dots)$$

quasi two-body channels of
three-body $\pi\pi\Lambda$ & $\pi\bar{K}N$

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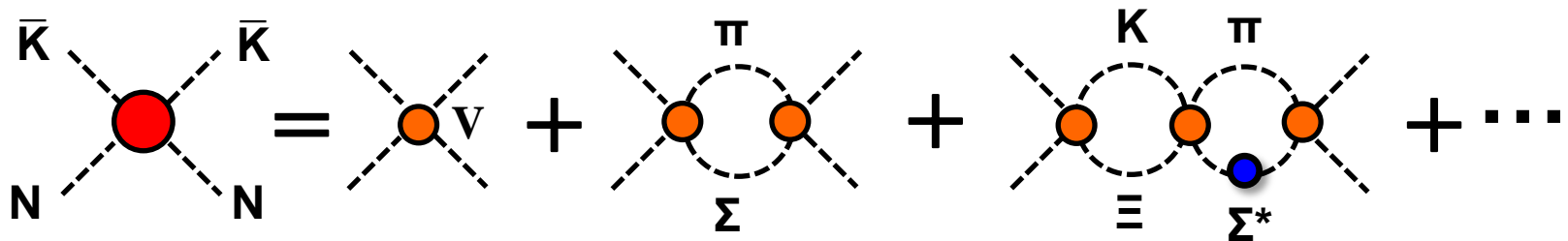
$$T_{a,b}^{(LSJ)}(p_a, p_b; E) = \underbrace{V_{a,b}^{(LSJ)}(p_a, p_b; E)}_{\text{CC effect}} + \underbrace{\sum_c \int_0^\infty q^2 dq V_{a,c}^{(LSJ)}(p_a, q; E) G_c(q; E) T_{c,b}^{(LSJ)}(q, p_b; E)}_{\text{off-shell effect}}$$

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quasi two-body channels of
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- ✓ Summing up all possible transitions between reaction channels !!
(\Rightarrow satisfies **multichannel two-** and **three-body unitarity**)

e.g.) $\bar{K}N$ scattering



- ✓ **Momentum integral** takes into account **off-shell rescattering effects** in the intermediate processes.

What we have done so far

With the DCC approach developed for the $S = -1$ sector, we made:

- ✓ Comprehensive analysis of **ALL** available data (**more than 17,000** data points) of $K^- p \rightarrow \bar{K}N, \pi\Sigma, \pi\Lambda, \eta\Lambda, K\Xi$ up to $W = 2.1$ GeV.

[HK, Nakamura, Lee, Sato, PRC90(2014)065204]

- ✓ Determination of threshold parameters (scattering lengths, effective ranges,...); the **partial-wave amplitudes** of $\bar{K}N \rightarrow \bar{K}N, \pi\Sigma, \pi\Lambda, \eta\Lambda, K\Xi$ for $S, P, D,$ and F waves.

[HK, Nakamura, Lee, Sato, PRC90(2014)065204]

- ✓ Extraction of $Y^* = (\Lambda^*, \Sigma^*)$ **resonance parameters** (mass, width, **couplings, ...**) defined by **poles of scattering amplitudes**.

[HK, Nakamura, Lee, Sato, PRC92(2015)025205]

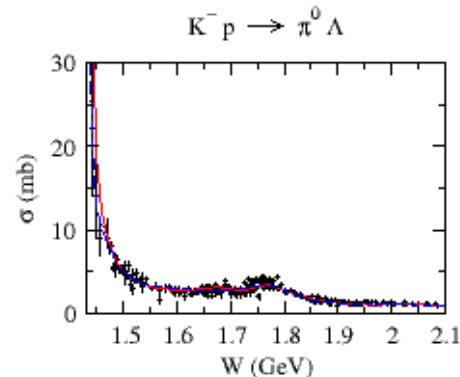
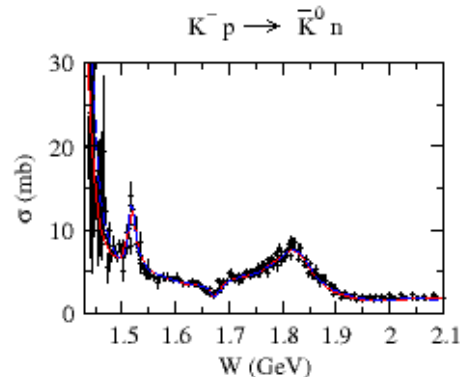
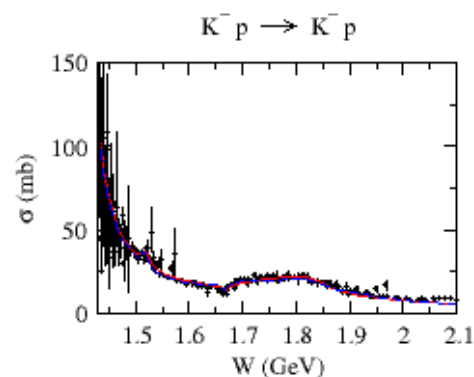
Supercomputers are necessary for the analysis !!



Results of the fits

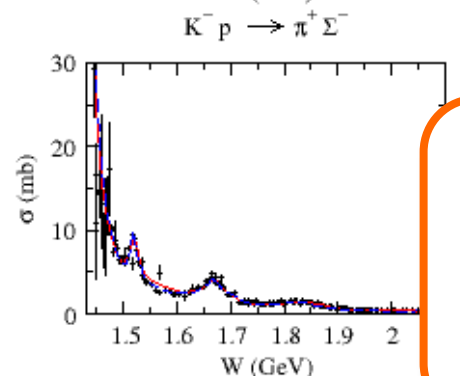
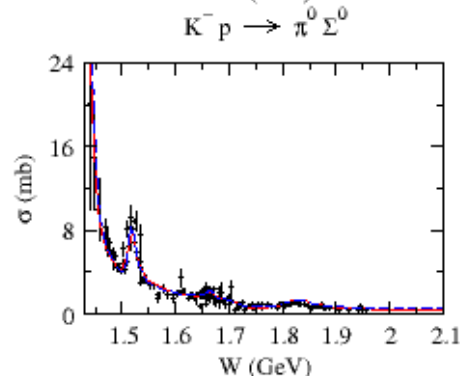
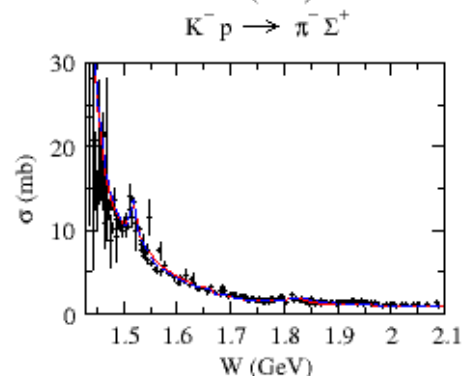
$K^- p \rightarrow \text{MB total cross sections}$

HK, Nakamura, Lee, Sato, PRC90(2014)065204

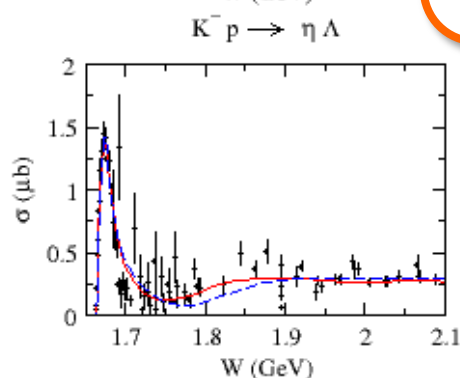
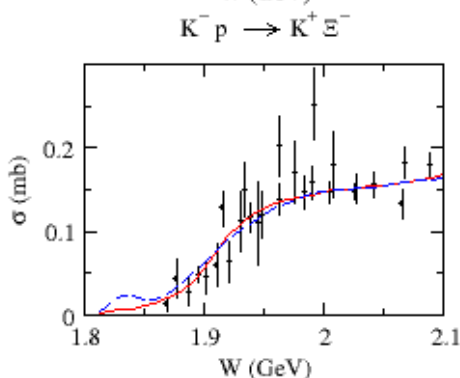
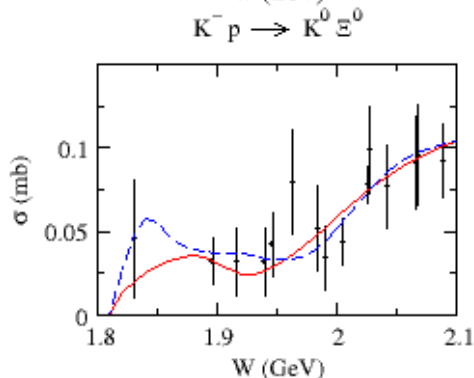


Red: Model A

Blue: Model B



“Incompleteness” of the current database allows us to have two parameter sets that give similar quality of the fit.

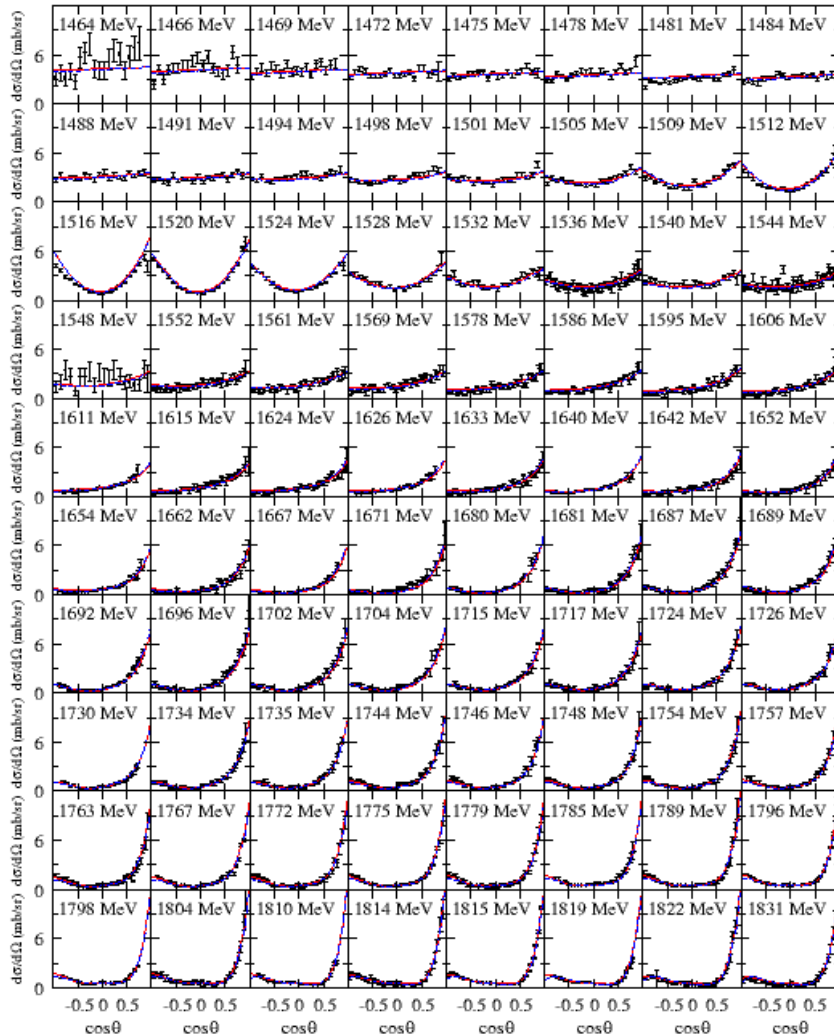


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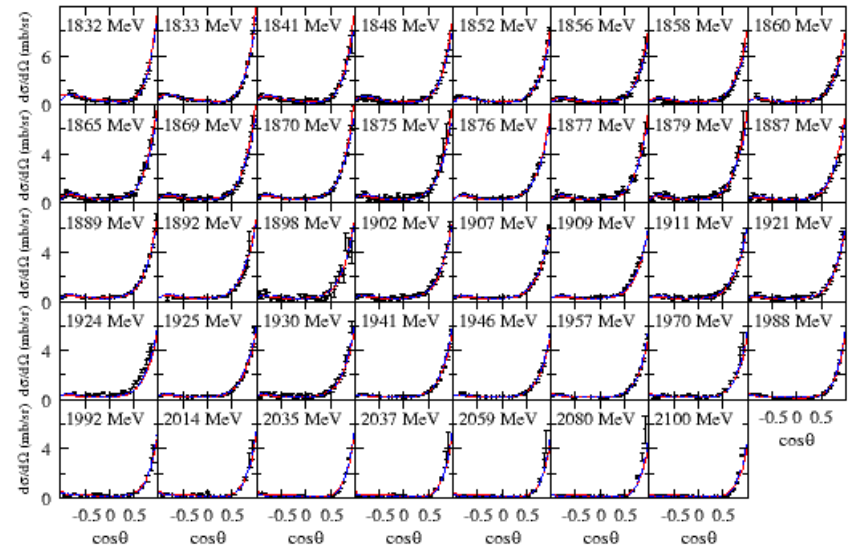
$K^- p \rightarrow K^- p$ scattering

HK, Nakamura, Lee, Sato, PRC90(2014)065204

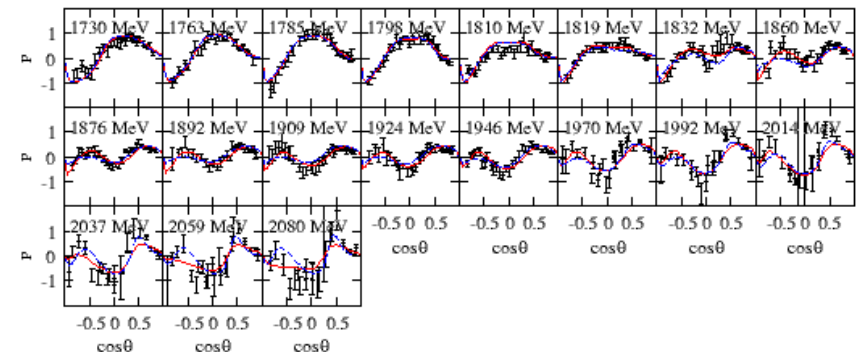
$d\sigma/d\Omega$ (1464 < W < 1831 MeV)



$d\sigma/d\Omega$ (1832 < W < 2100 MeV)

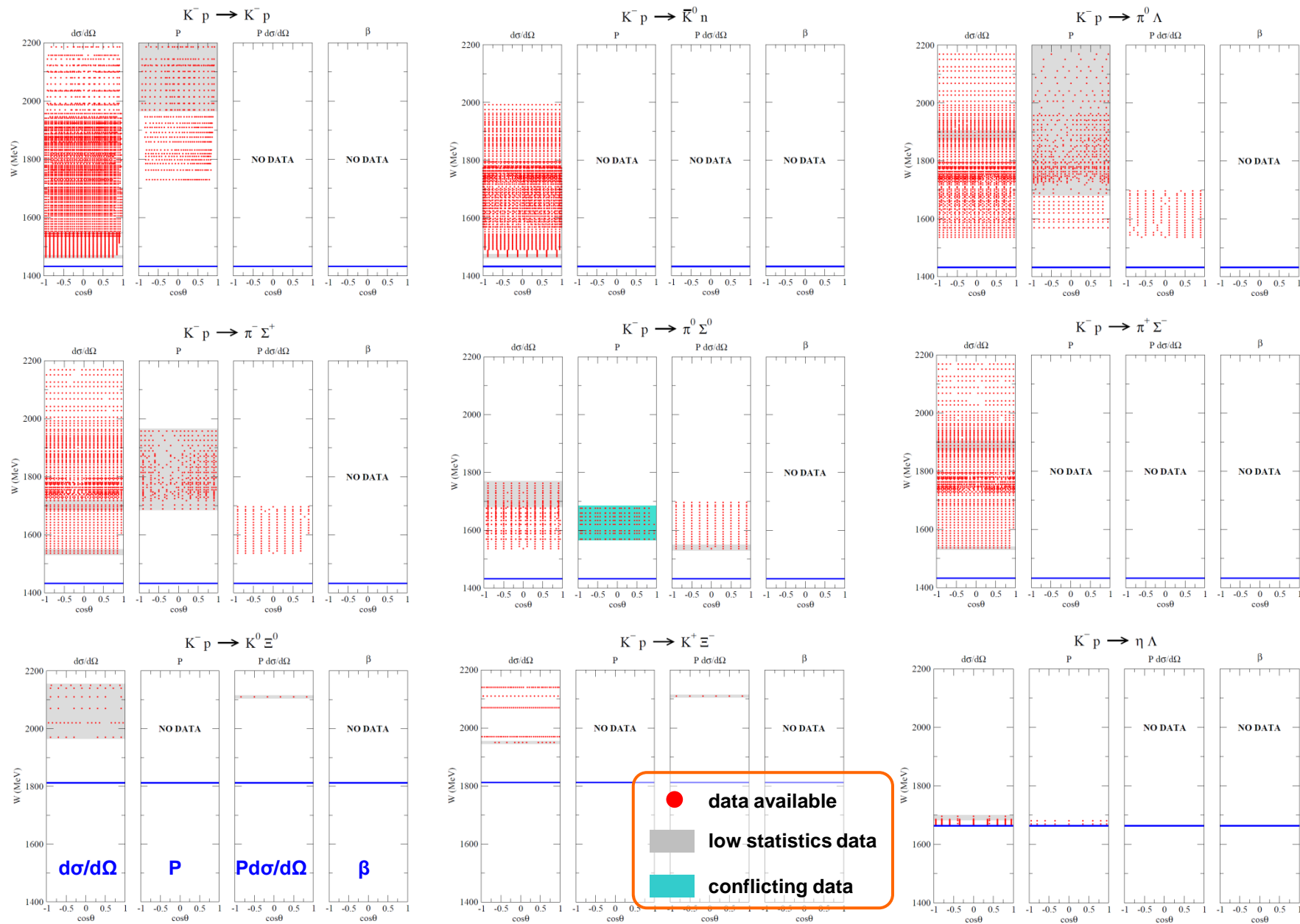


P (1730 < W < 2080 MeV)



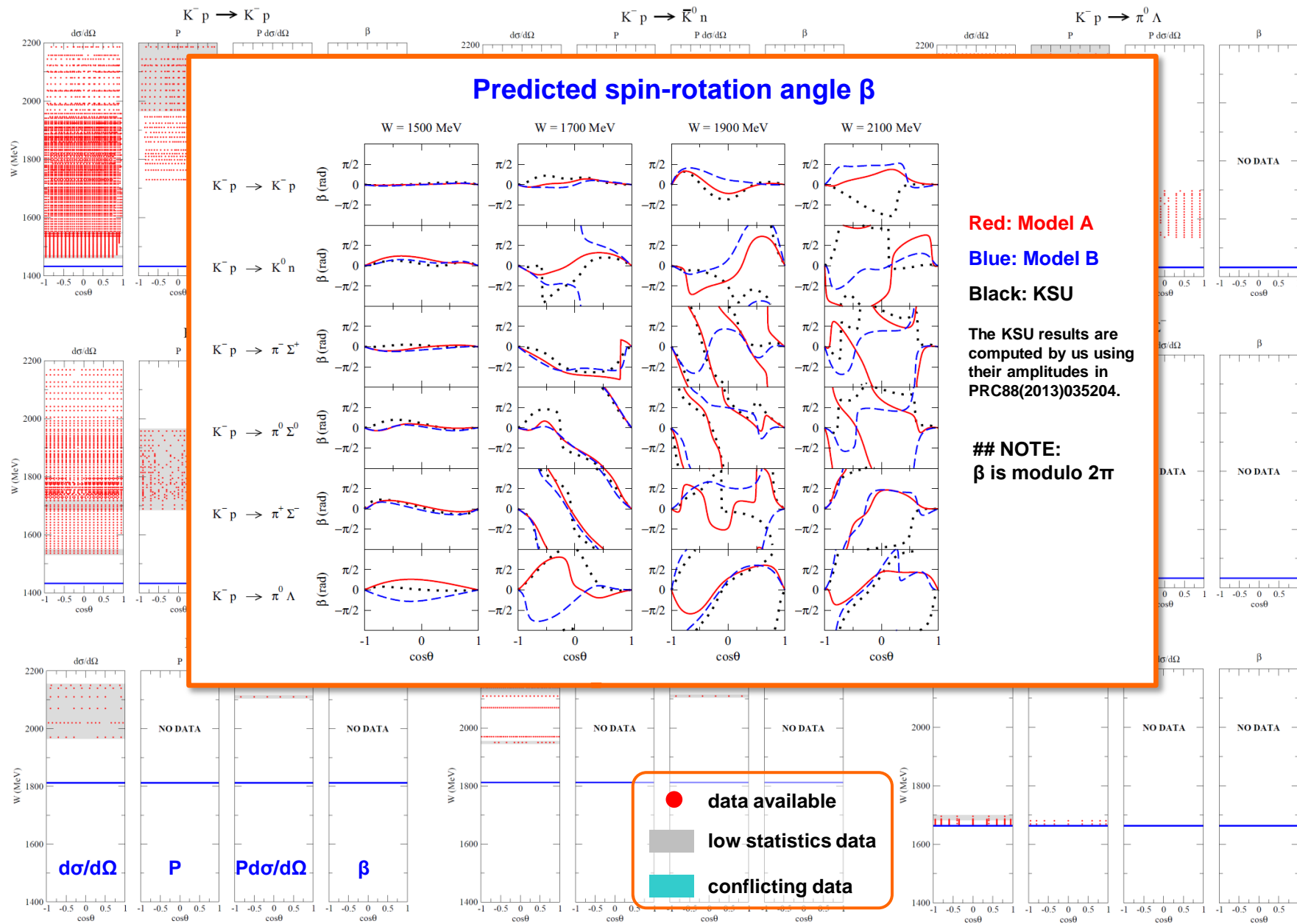
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Kinematical (W, cosθ) coverage of available $K^- p \rightarrow \bar{K} N, \pi \Sigma, \pi \Lambda, \eta \Lambda, K \Xi$ data



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Predicted spin-rotation angle β



Extracted Λ^* and Σ^* mass spectrum

Spectrum for Y^* resonances found above the $\bar{K}N$ threshold

HK, Nakamura, Lee, Sato, PRC92(2015)025205
(+ updates)

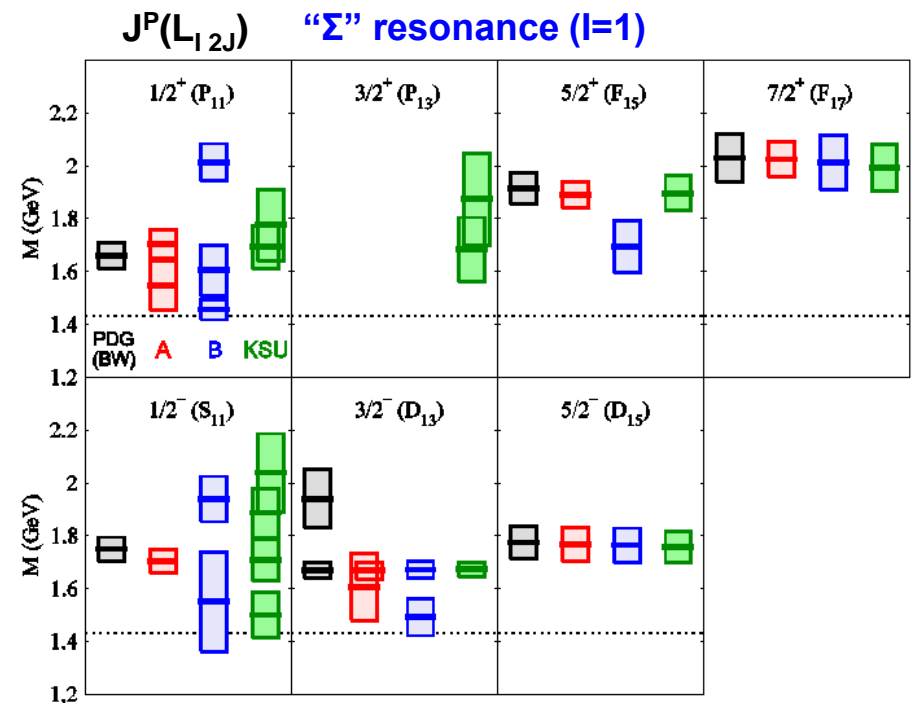
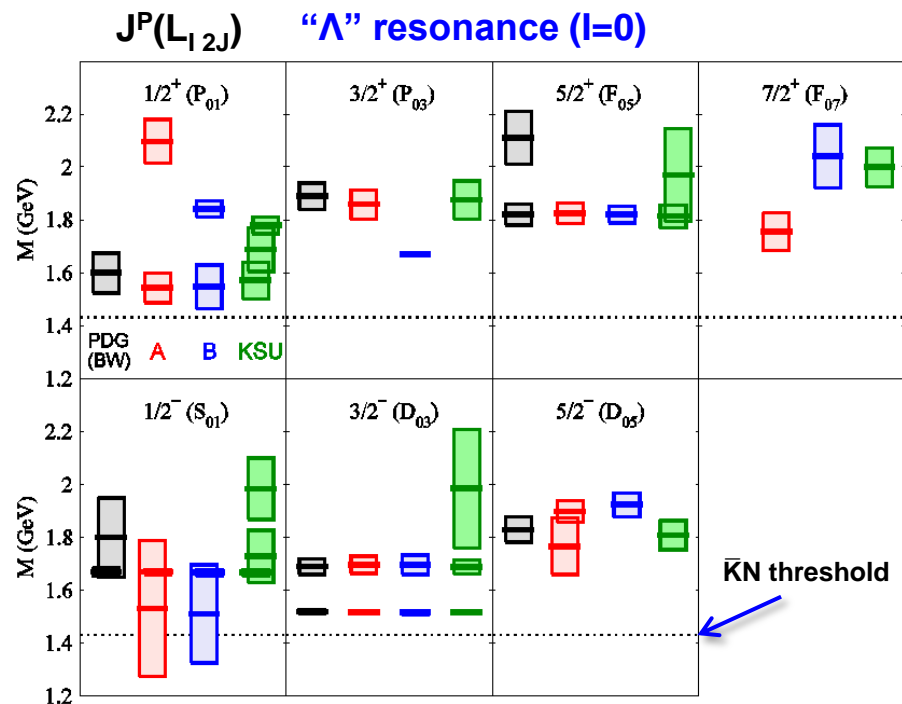
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Green: KSU[PRC88(2013)035205]

Black: PDG (only 4- & 3-star Y^* ;
Breit-Wigner)

$$\left. \begin{array}{l} -2\text{Im}(M_R) \\ \text{("width")} \end{array} \right\} \text{Re}(M_R) \quad M_R : \text{Resonance pole mass (complex)}$$



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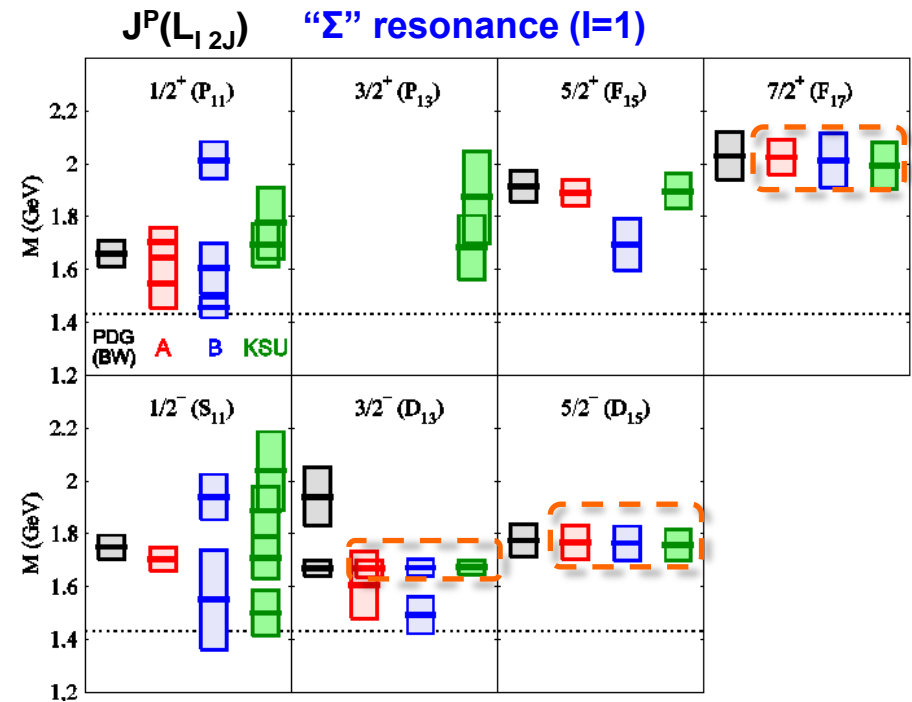
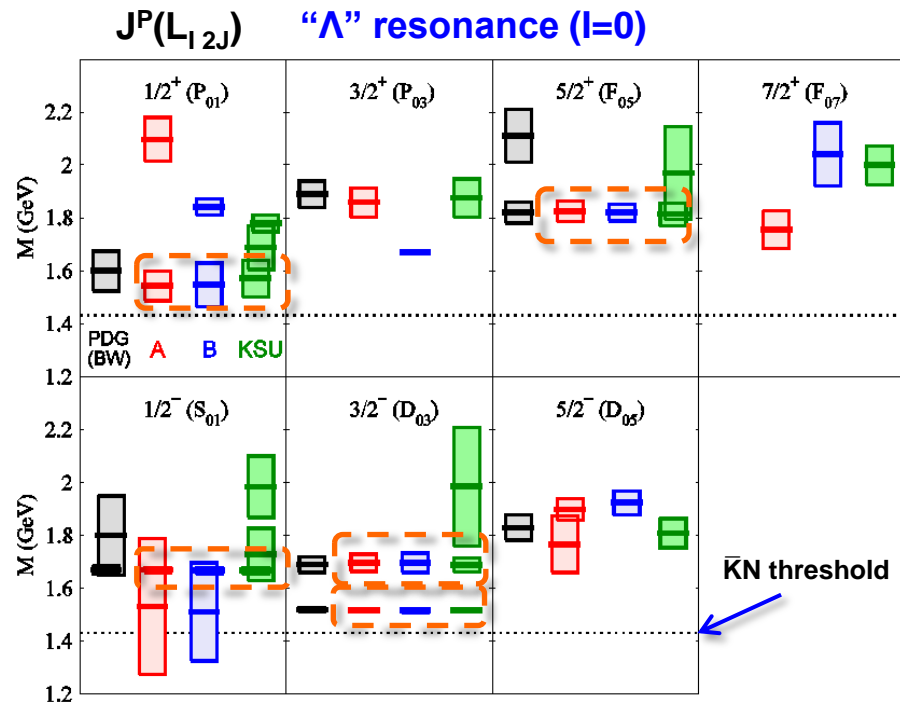
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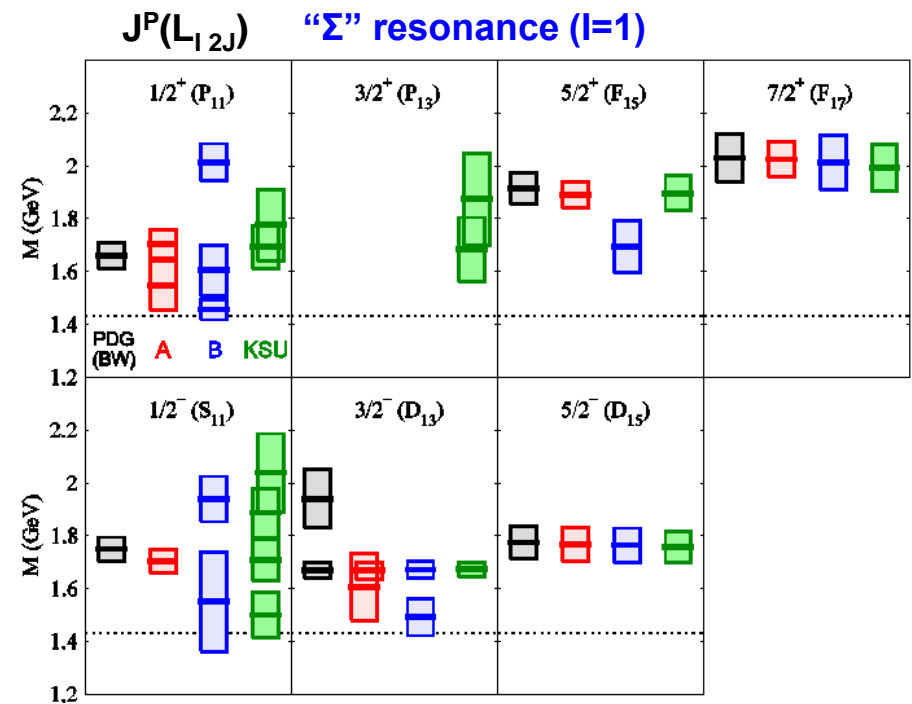
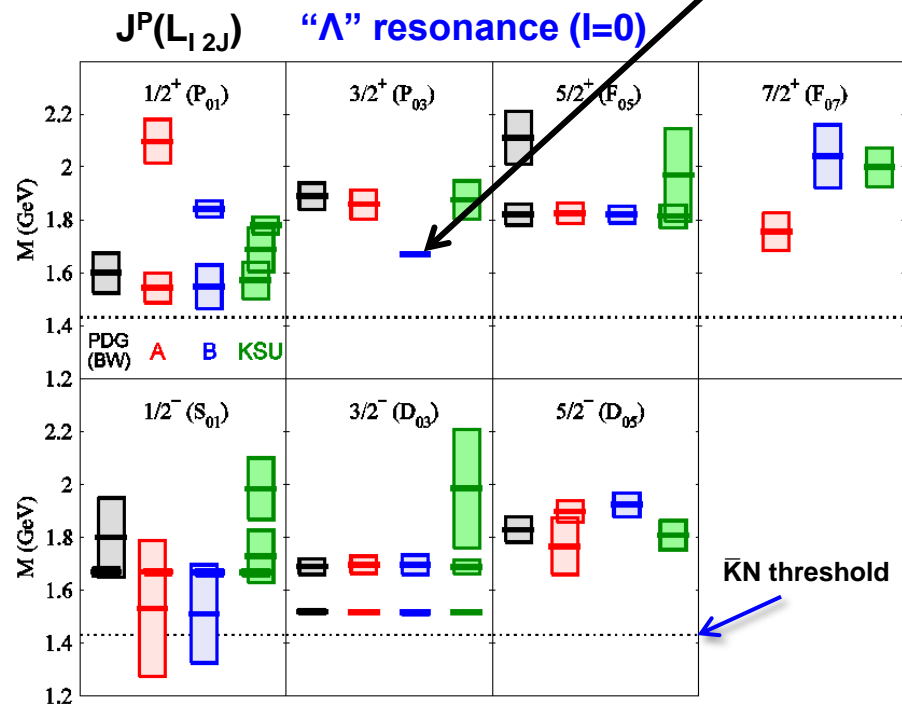
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New narrow $3/2^+$ resonance
 $M = 1671 - 5i$ MeV
near the $\eta\Lambda$ threshold !!

$$\left. \begin{array}{l} -2\text{Im}(M_R) \\ \text{("width")} \end{array} \right\} \text{Re}(M_R) \quad M_R : \text{Resonance pole mass (complex)}$$



Extracted Λ^* and Σ^* mass spectrum

Spectrum for Y^* resonances found above the $\bar{K}N$ threshold

HK, Nakamura, Lee, Sato, PRC92(2015)025205
(+ updates)

Red: Model A

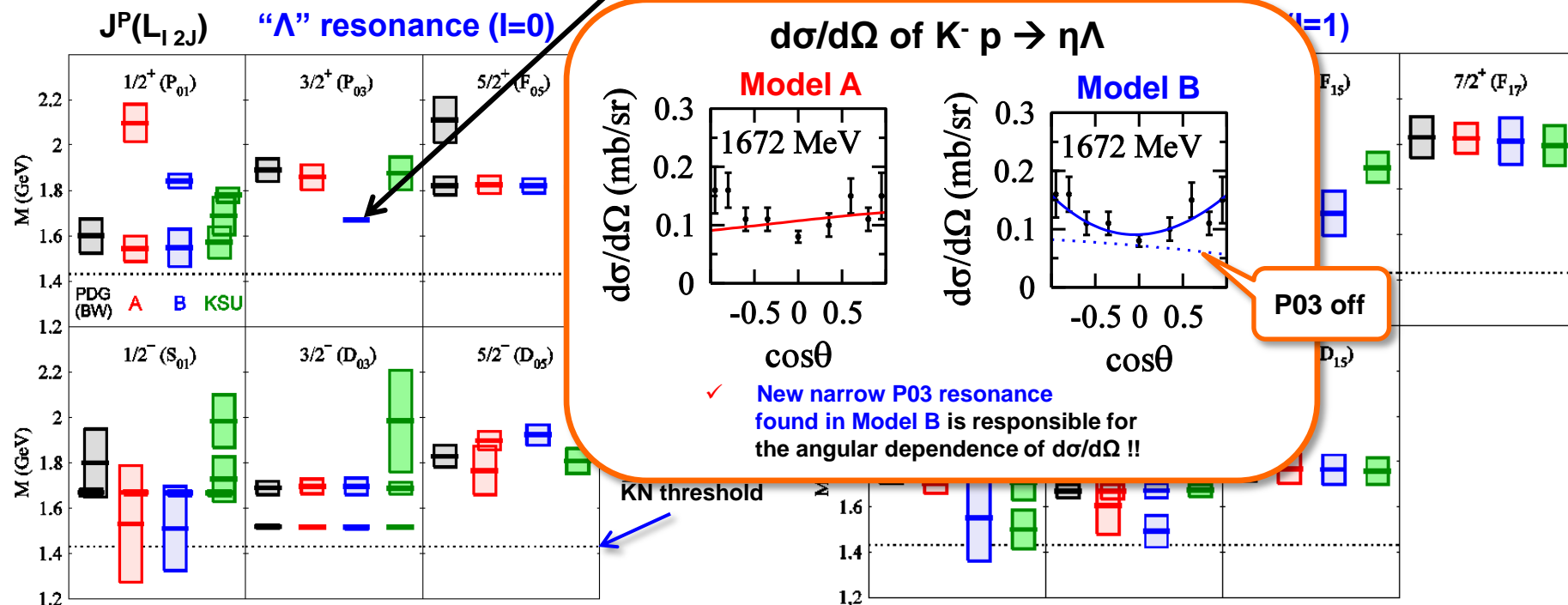
Blue: Model B

Green: KSU[PRC88(2013)035205]

Black: PDG (only 4- & 3-star Y^* ;
Breit-Wigner)

New narrow $3/2^+$ resonance
 $M = 1671 - 5i$ MeV
near the $\eta\Lambda$ threshold !!

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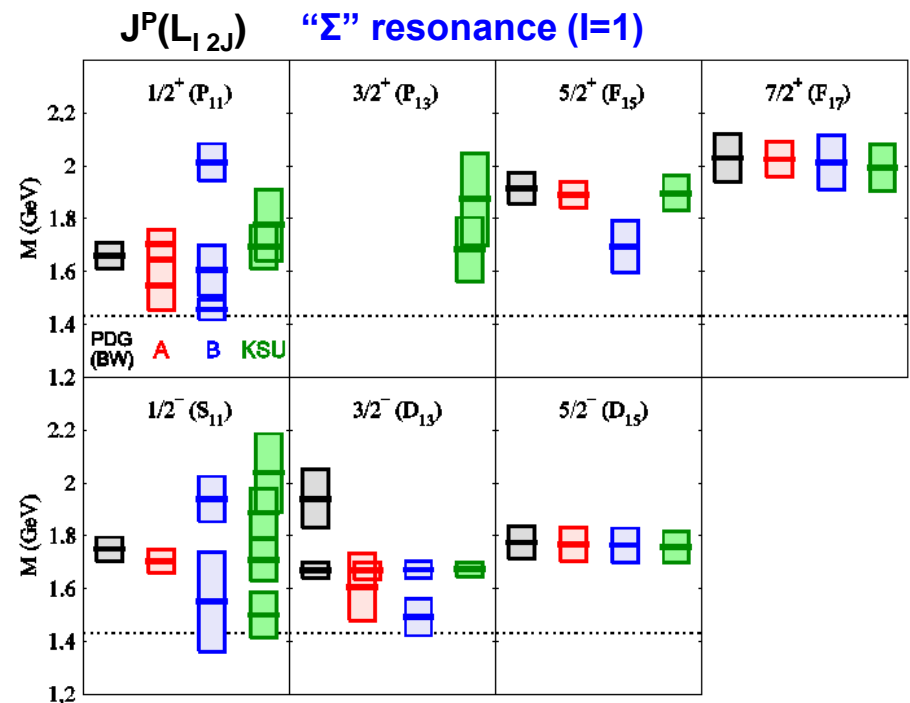
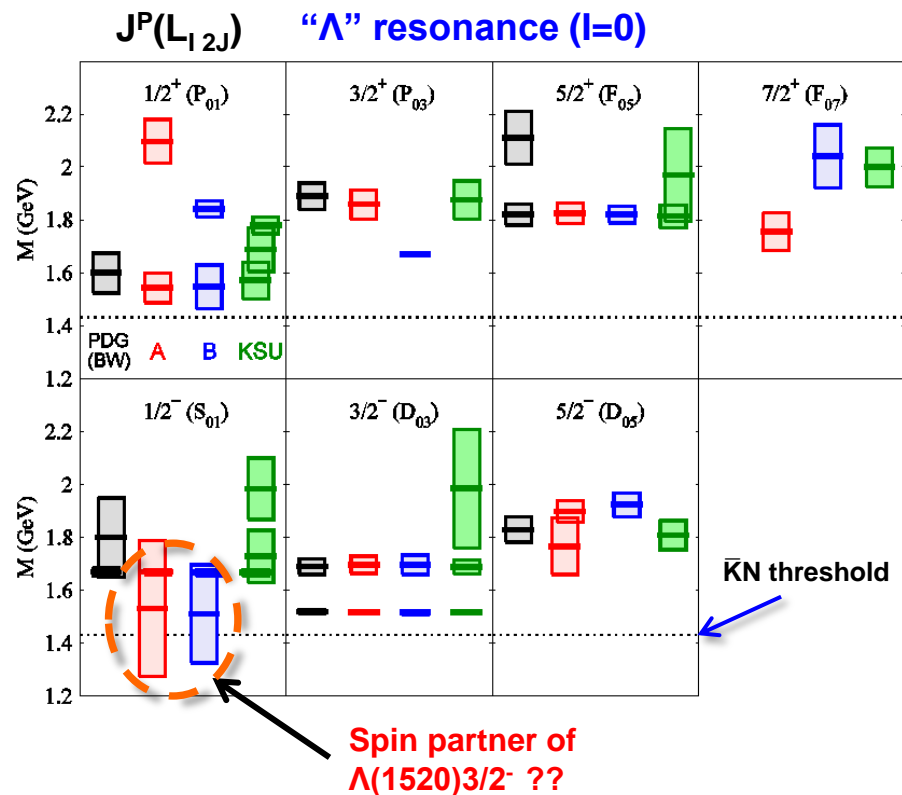
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(+ updates)

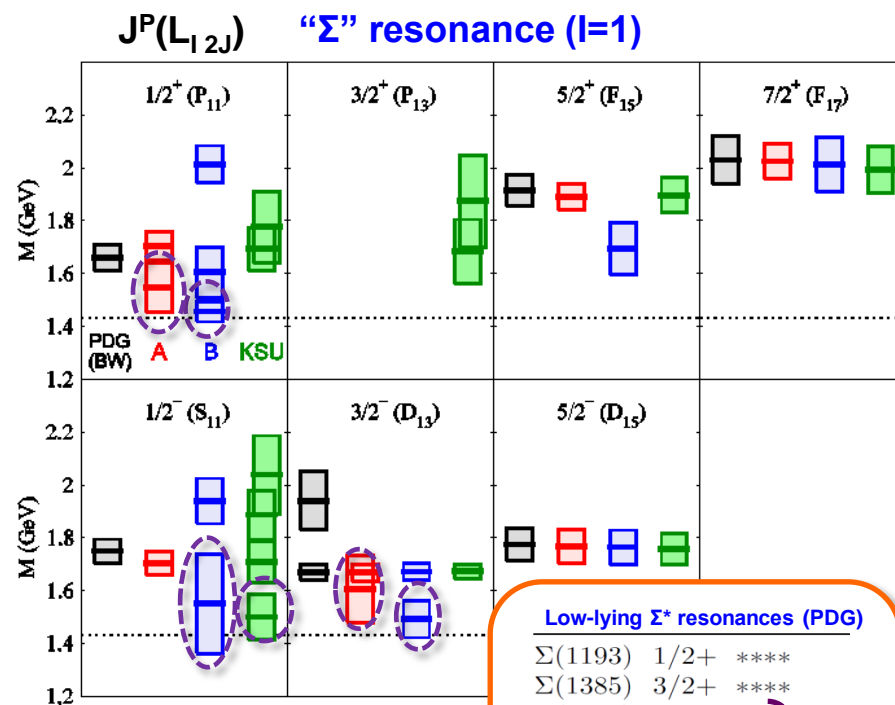
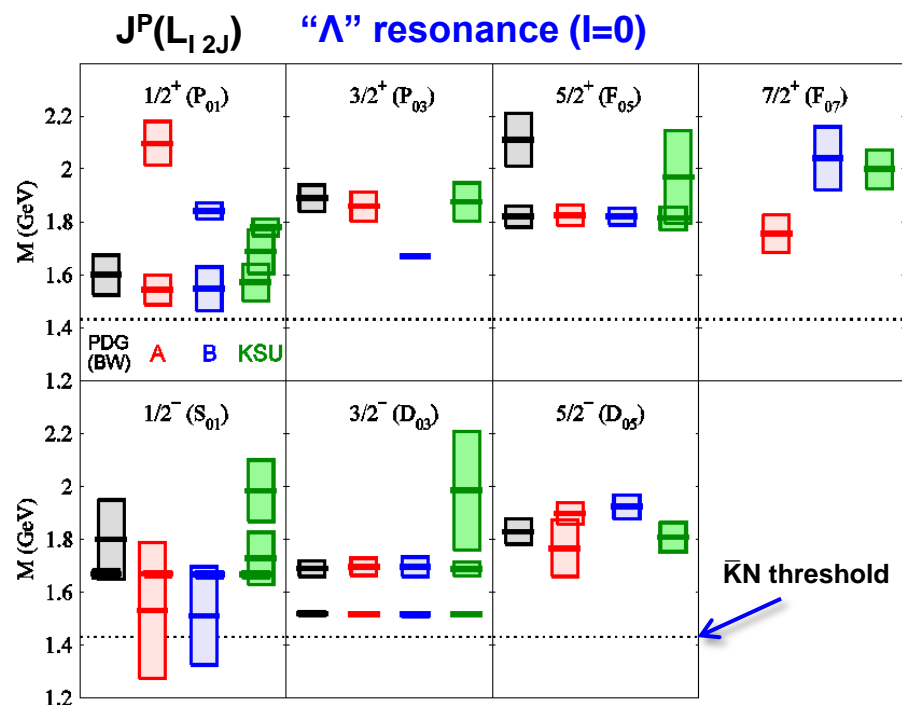
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Low-lying Σ^* resonances (PDG)

$\Sigma(1193)$	$1/2^+$	****
$\Sigma(1385)$	$3/2^+$	****
$\Sigma(1480)$		*
$\Sigma(1560)$		**
$\Sigma(1580)$	$3/2^-$	*
$\Sigma(1620)$	$1/2^-$	**
$\Sigma(1660)$	$1/2^+$	****
$\Sigma(1670)$	$3/2^-$	****

?

Importance of $2 \rightarrow 3$ reactions: Branching ratios of high-mass Y^* resonances

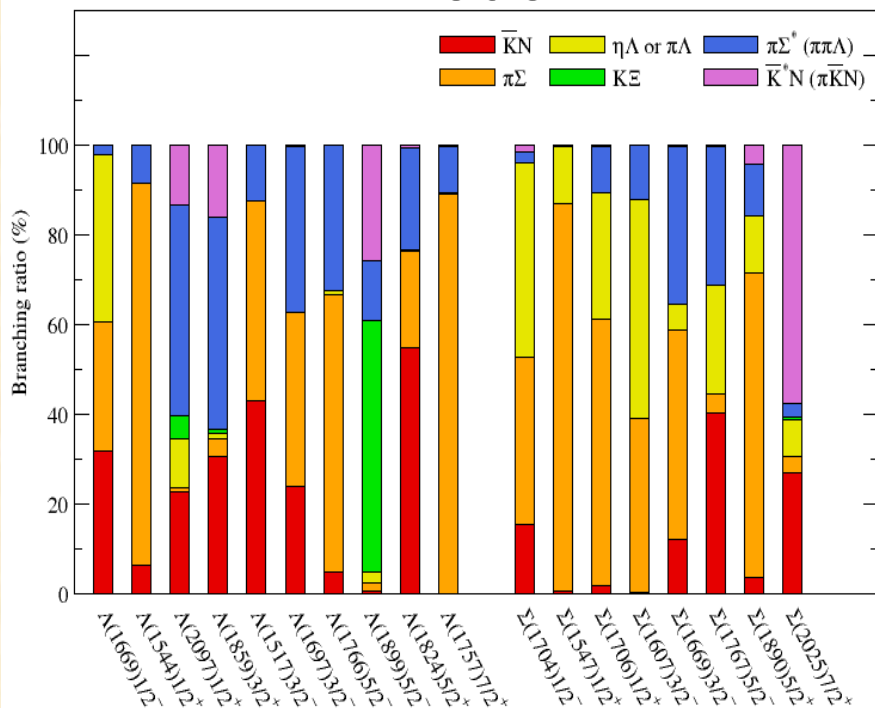
✓ High-mass Y^* have large branching ratio to $\pi\Sigma^*$ ($\pi\pi\Lambda$) & \bar{K}^*N ($\pi\bar{K}N$)

➤ $K^- p \rightarrow \pi\pi\Lambda, \pi\bar{K}N, \dots$ data would play a crucial role for establishing high-mass Y^* .

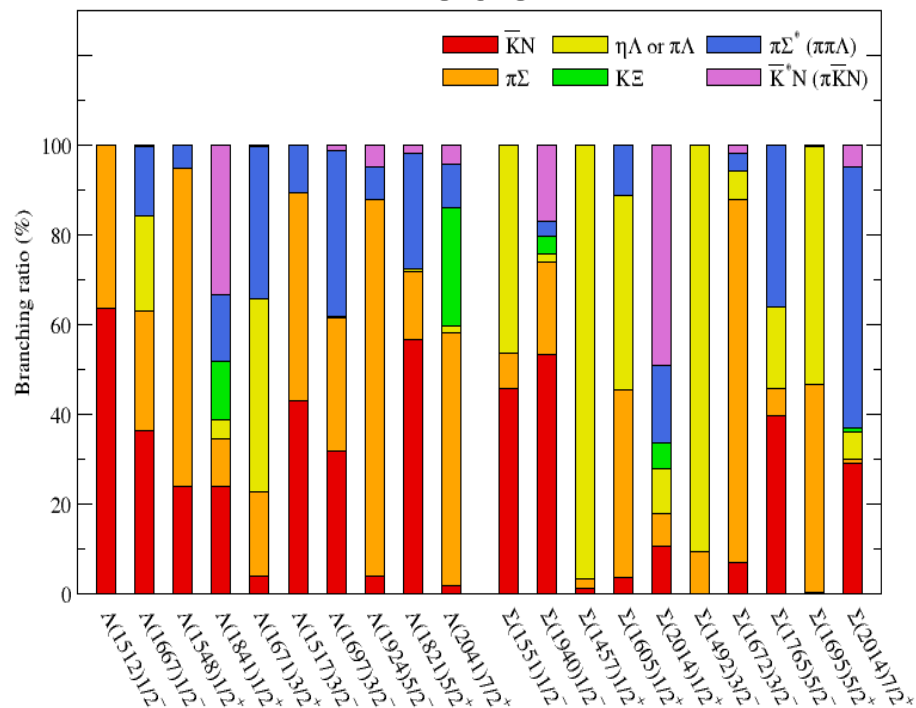
➔ Similar to high-mass N^* and Δ^* case, where $\pi\pi N$ channel plays a crucial role.

(e.g., measurement of $\pi N \rightarrow \pi\pi N$ reactions at **J-PARC E45**)

Model A



Model B



Importance of $2 \rightarrow 3$ reactions: Branching ratios of high-mass Y^* resonances

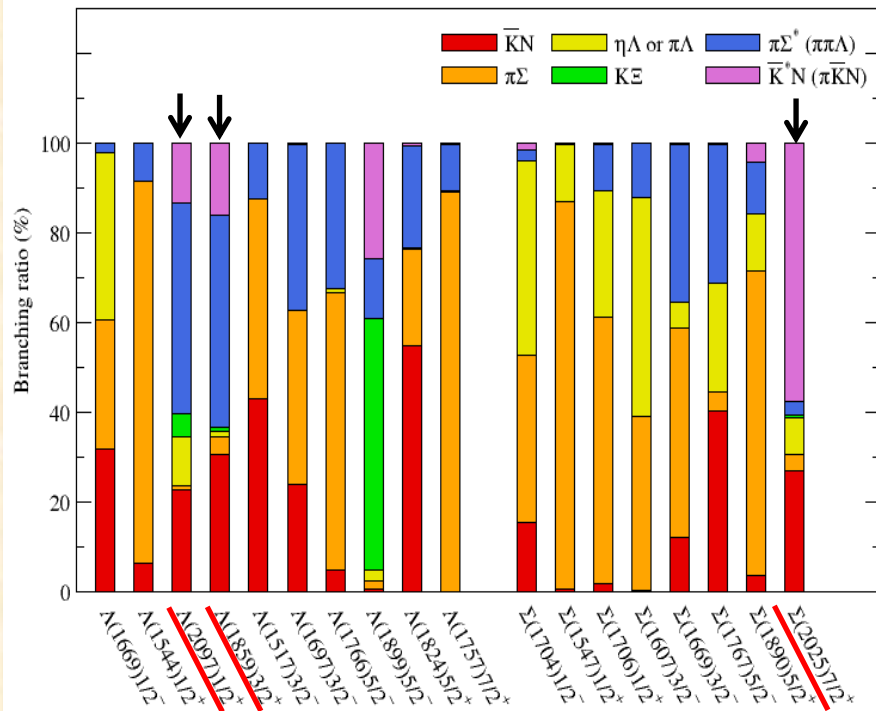
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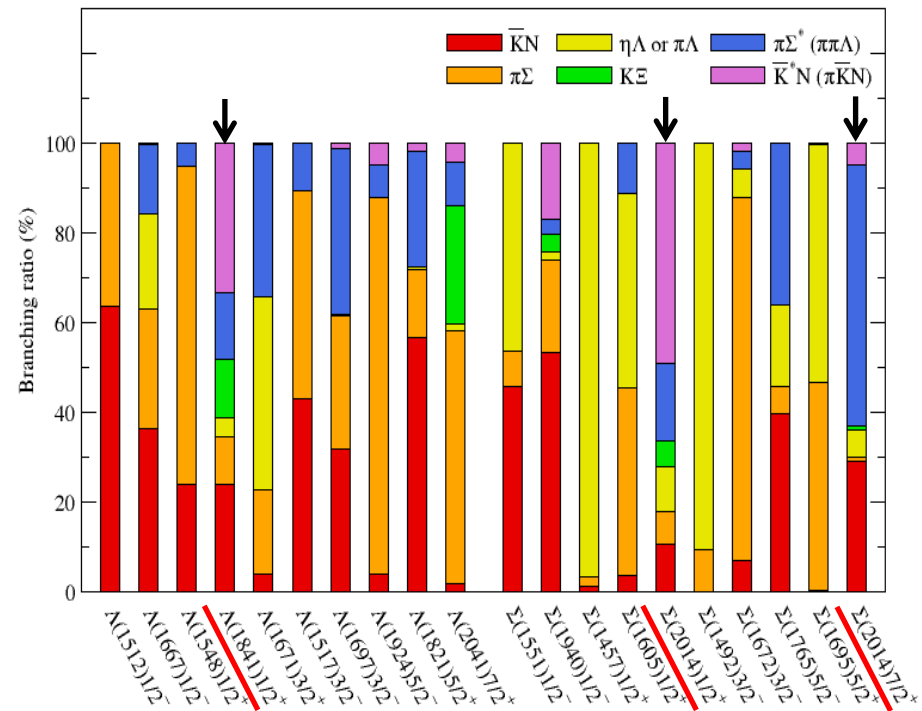
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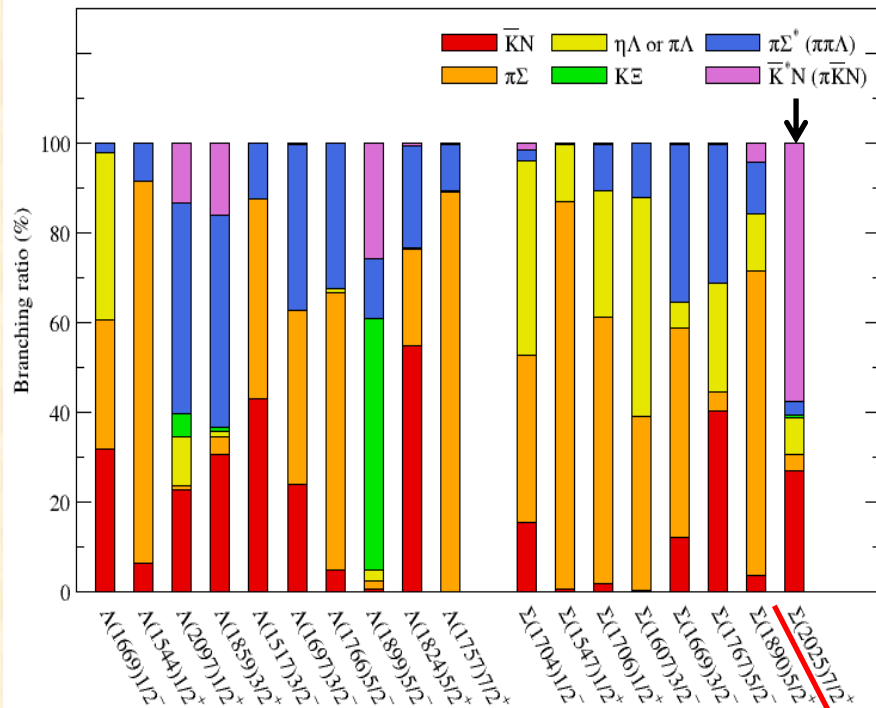
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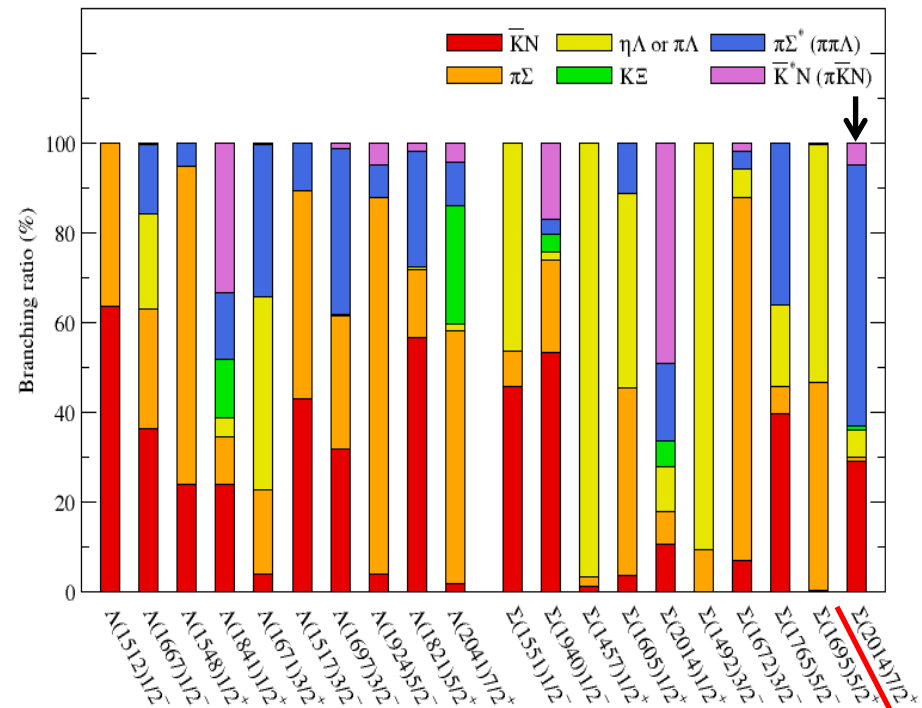
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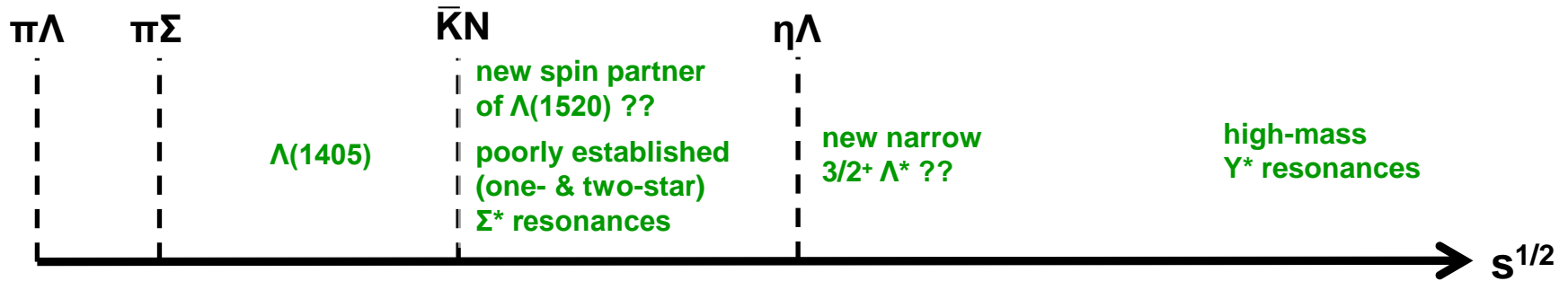
Model A



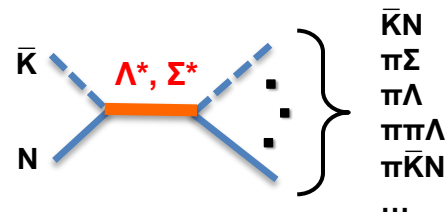
Model B



Strategy for establishing Y^* resonances using antikaon-induced reactions

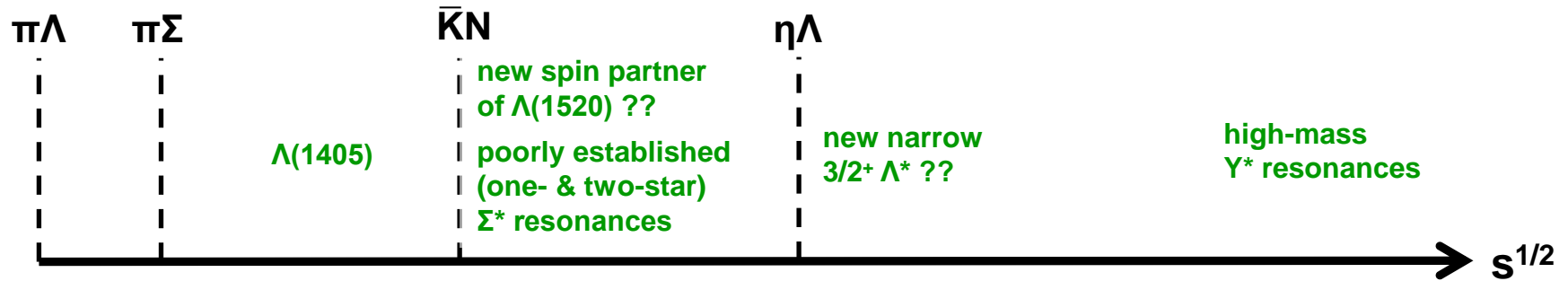


“Complete experiments” for $\bar{K}N \rightarrow \bar{K}N, \pi Y, \eta Y, K\Xi, \omega Y, \eta' Y, \dots$

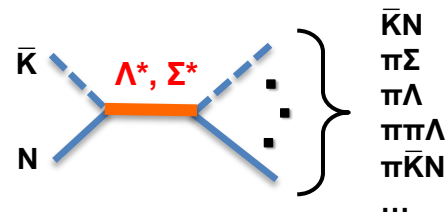


2 \rightarrow 3 reactions: $\bar{K}N \rightarrow \pi\pi\Lambda, \pi\bar{K}N, \dots$

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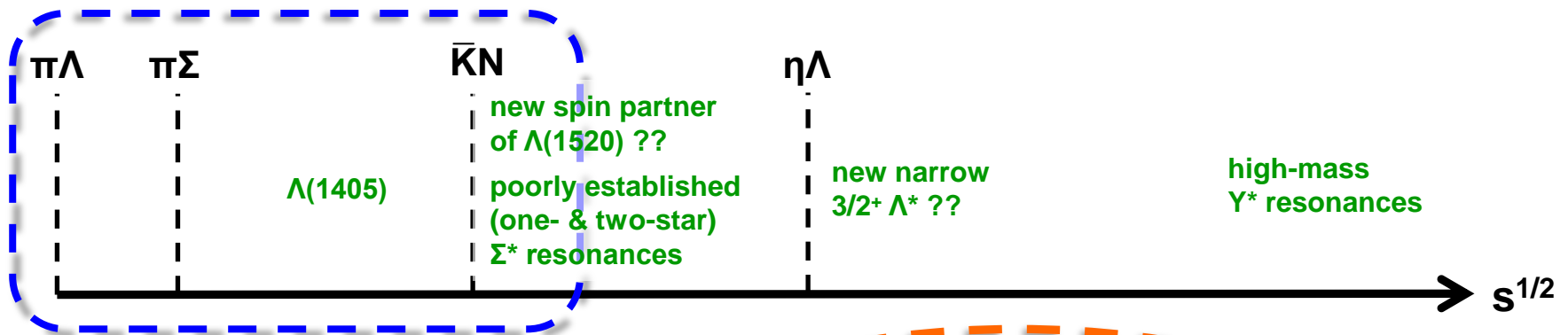


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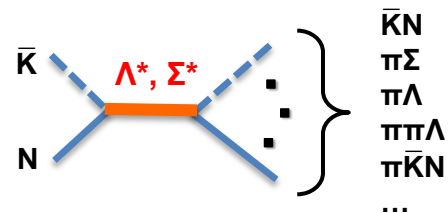
New measurements at J-PARC ??

[discussions with H. Sako, K. Hicks et al.]

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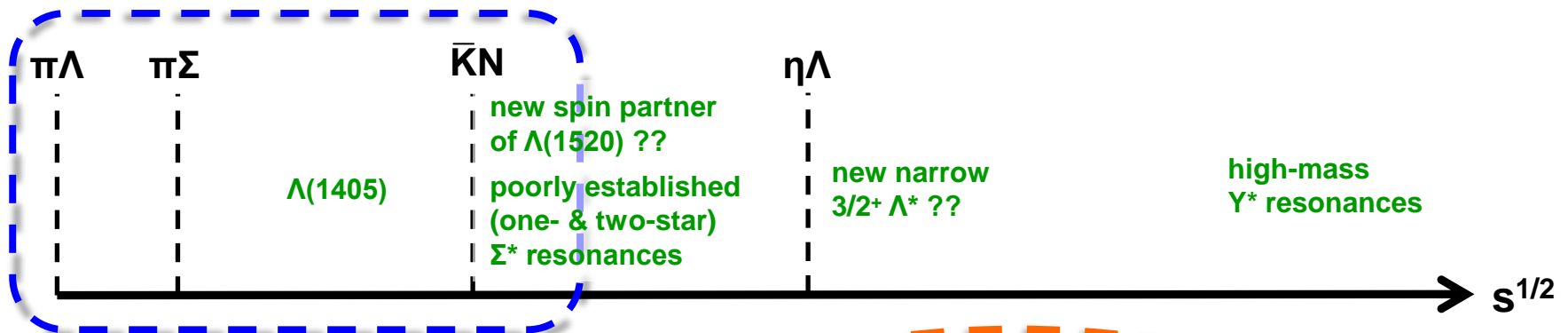


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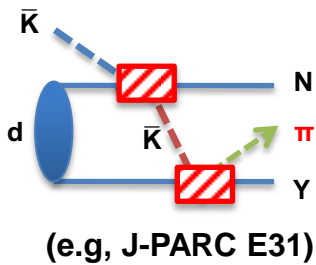
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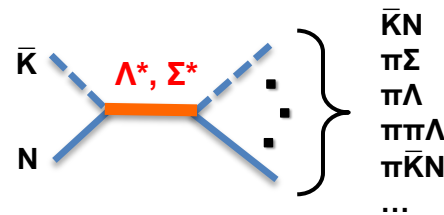


$\bar{K}d \rightarrow \pi YN$



Application of our DCC approach to $\bar{K}d$ reactions is underway (HK & Lee), aiming at **COMBINED analysis of $\bar{K}N$ and $\bar{K}d$ reactions** !!!

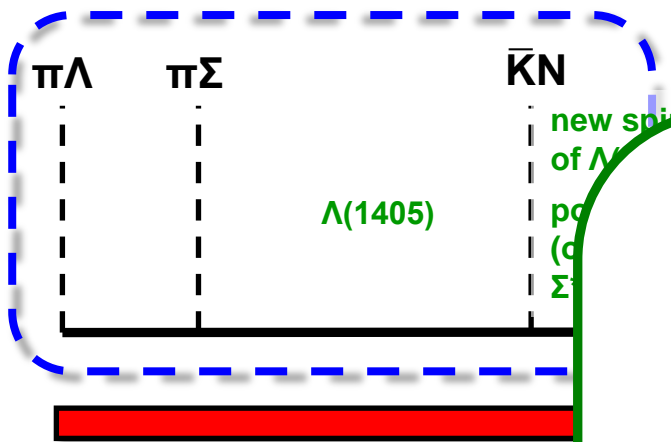
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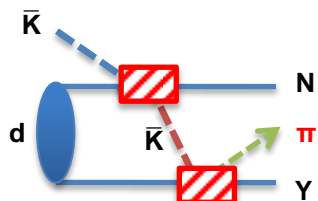
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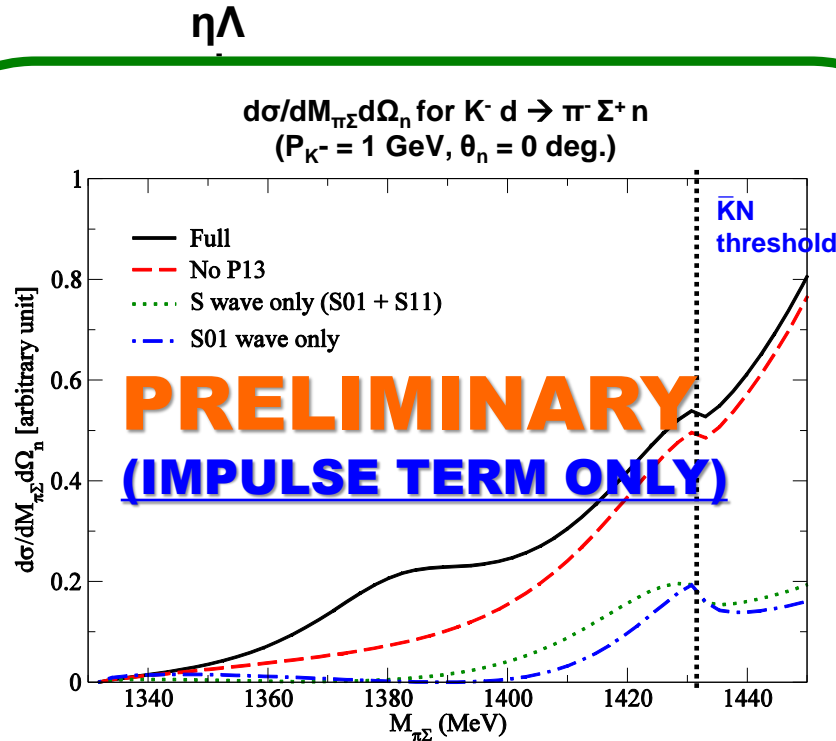


$\bar{K}d \rightarrow \pi YN$



(e.g., J-PARC E31)

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Construction of $K^- d$ reaction model is underway !!!

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New measurements at J-PARC ??
[discussions with H. Sako, K. Hicks et al.]

Summary

- ✓ Accomplished comprehensive analysis of $K^- p \rightarrow \bar{K}N, \pi\Sigma, \pi\Lambda, \eta\Lambda, K\Xi$ up to $W = 2.1$ GeV for the first time within a dynamical coupled-channels approach.
- ✓ Successfully extracted partial-wave amplitudes (up to F wave) and Y^* resonance parameters defined by poles of amplitudes.
 - New narrow $J^P = 3/2^+ \Lambda^*$ resonance ($M_R = 1672-i5$ MeV) located near the $\eta\Lambda$ threshold
 - New $J^P = 1/2^- \Lambda^*$ resonance ($\text{Re } M_R \sim 1520$ MeV) with mass close to $\Lambda(1520)3/2^-$
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- ✓ New accurate data for both $\bar{K}N$ and $\bar{K}d$ reactions are much appreciated !!!
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 - Deuteron-target reaction ($\bar{K}d \rightarrow \pi YN, \dots$) to determine low-lying Y^*

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The help of J-PARC is definitely needed for “complete” determination of Y^* resonance mass spectrum !!!

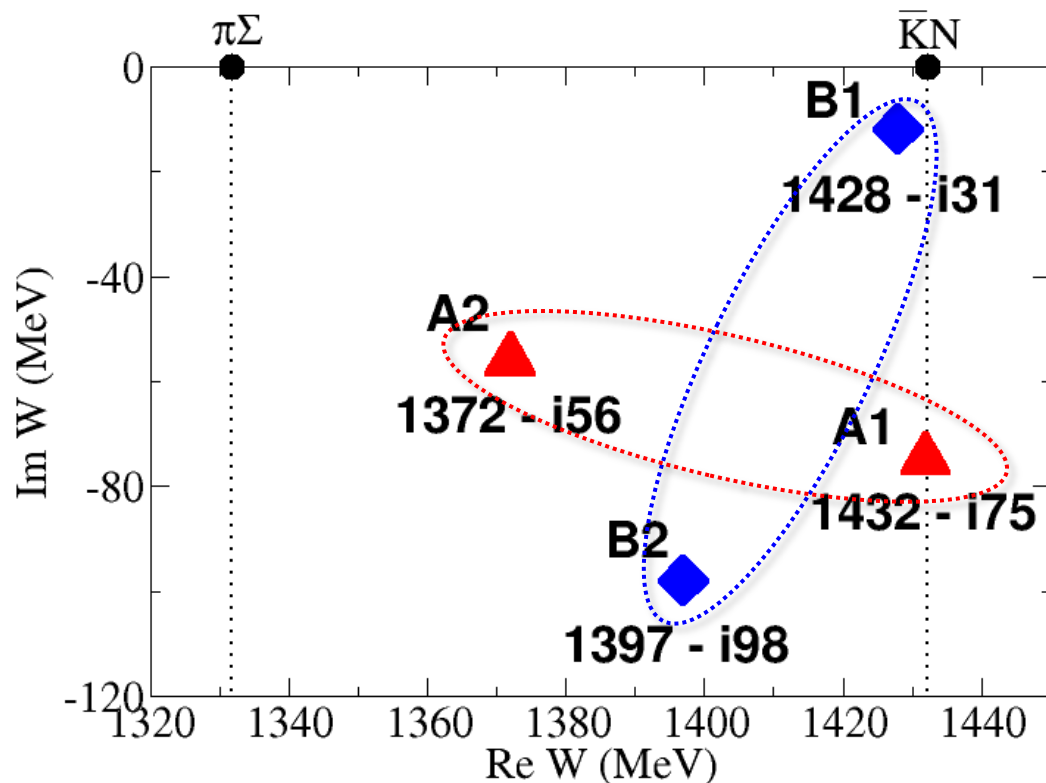
Back up

S-wave resonances below $\bar{K}N$ threshold from the current analysis

HK, Nakamura, Lee, Sato, PRC92(2015)025205

NOTE: Further extensive analysis including the data below $\bar{K}N$ threshold is necessary to have *conclusive* results for the $\bar{K}N$ subthreshold region.

“Predicted” Λ^* ($J^P = 1/2^-$) resonance poles below $\bar{K}N$ threshold




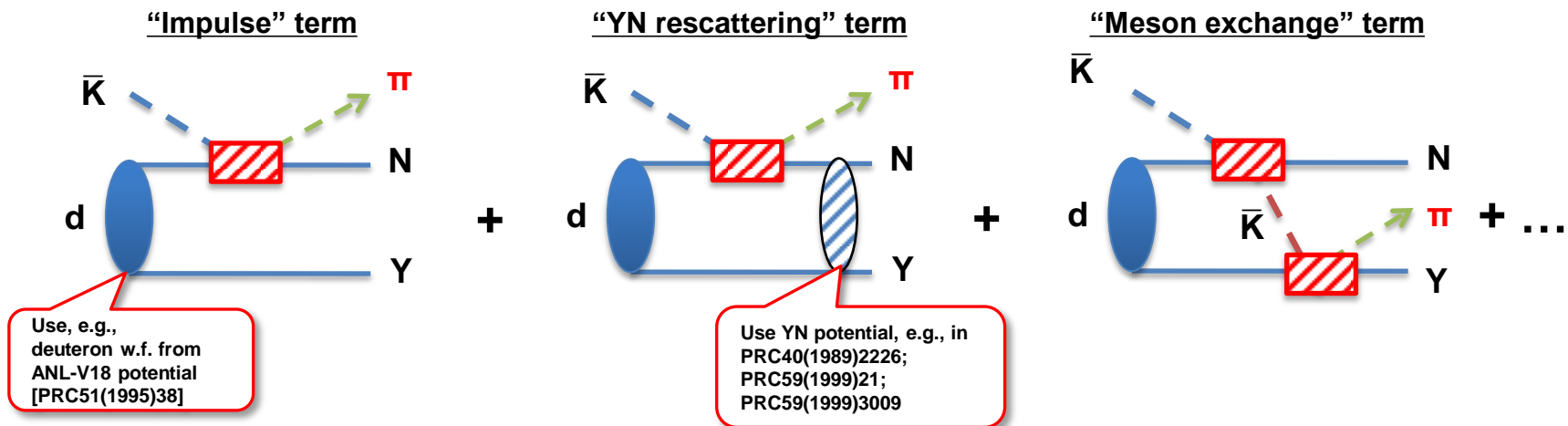
✓ Two resonance poles are found in both Models **A** and **B**.

- **A1** & **B1** seem correspond to $\Lambda(1405)$
- Another Λ resonance with mass 30-60 MeV lower than $\Lambda(1405)$ (**A2** & **B2**) is also found to exist.

Red triangles: **Model A**
Blue diamonds: **Model B**

Model for deuteron-target reactions

- ✓ Multistep processes are treated “perturbatively”.
- ✓ **Full-off-shell** amplitudes for meson-baryon sub-processes () are taken from **our dynamical coupled-channels model**. HK, Nakamura, Lee, Sato, PRC90(2014)065203



Unique feature of our work:

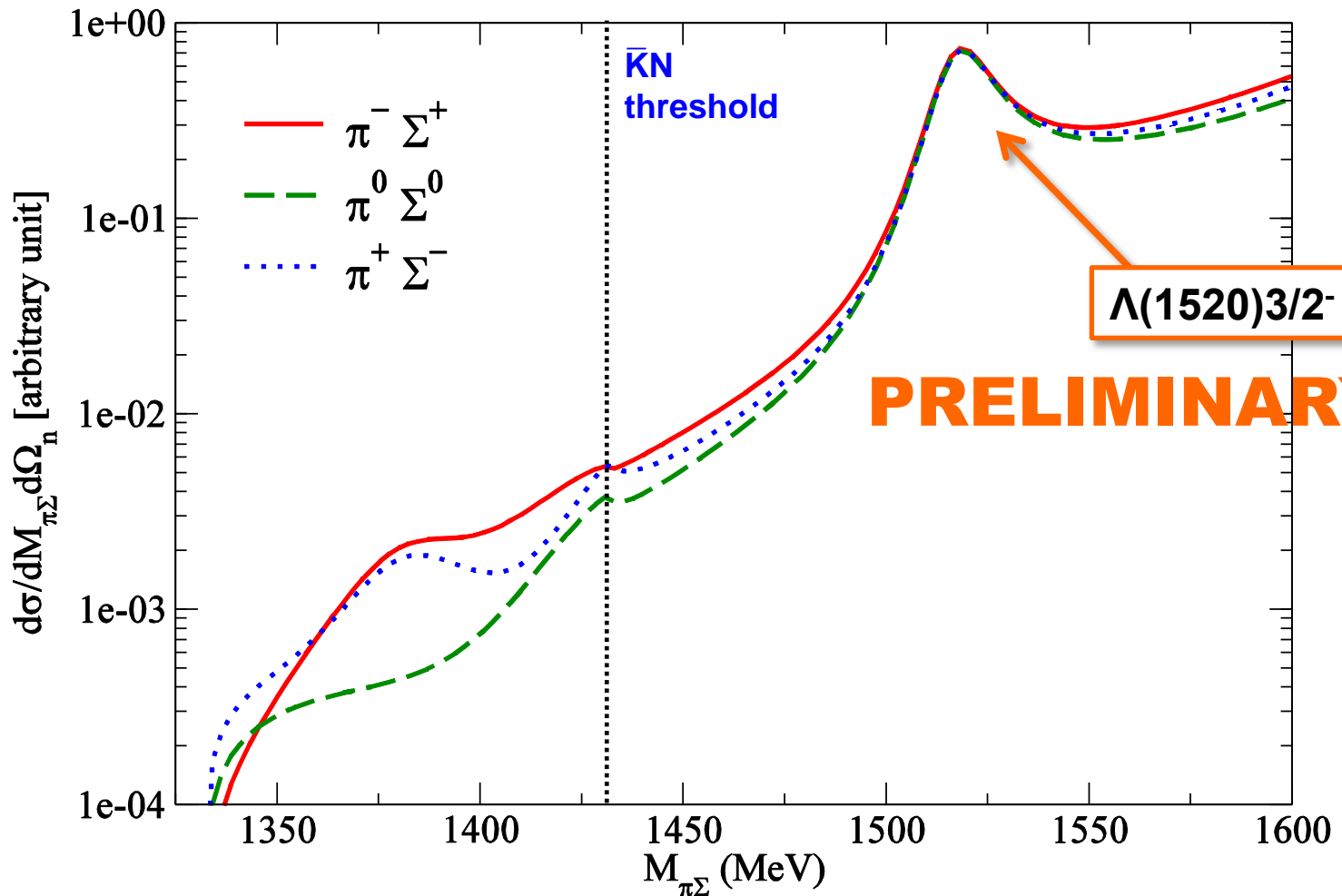
For meson-baryon sub-processes, we have **full-off-shell** amplitudes

- well-tested by $K^- p \rightarrow \bar{K}N, \pi\Sigma, \pi\Lambda, \eta\Lambda, K\Xi$ up to $W = 2.1$ GeV.
- not only for S wave, **but also P, D, F waves**.

Results (IMPULSE TERM ONLY !!)

$d\sigma/dM_{\pi\Sigma}d\Omega_n$ for $K^- d \rightarrow (\pi \Sigma)_0 n$
($P_{K^-} = 1 \text{ GeV}$, $\theta_n = 0 \text{ deg.}$)

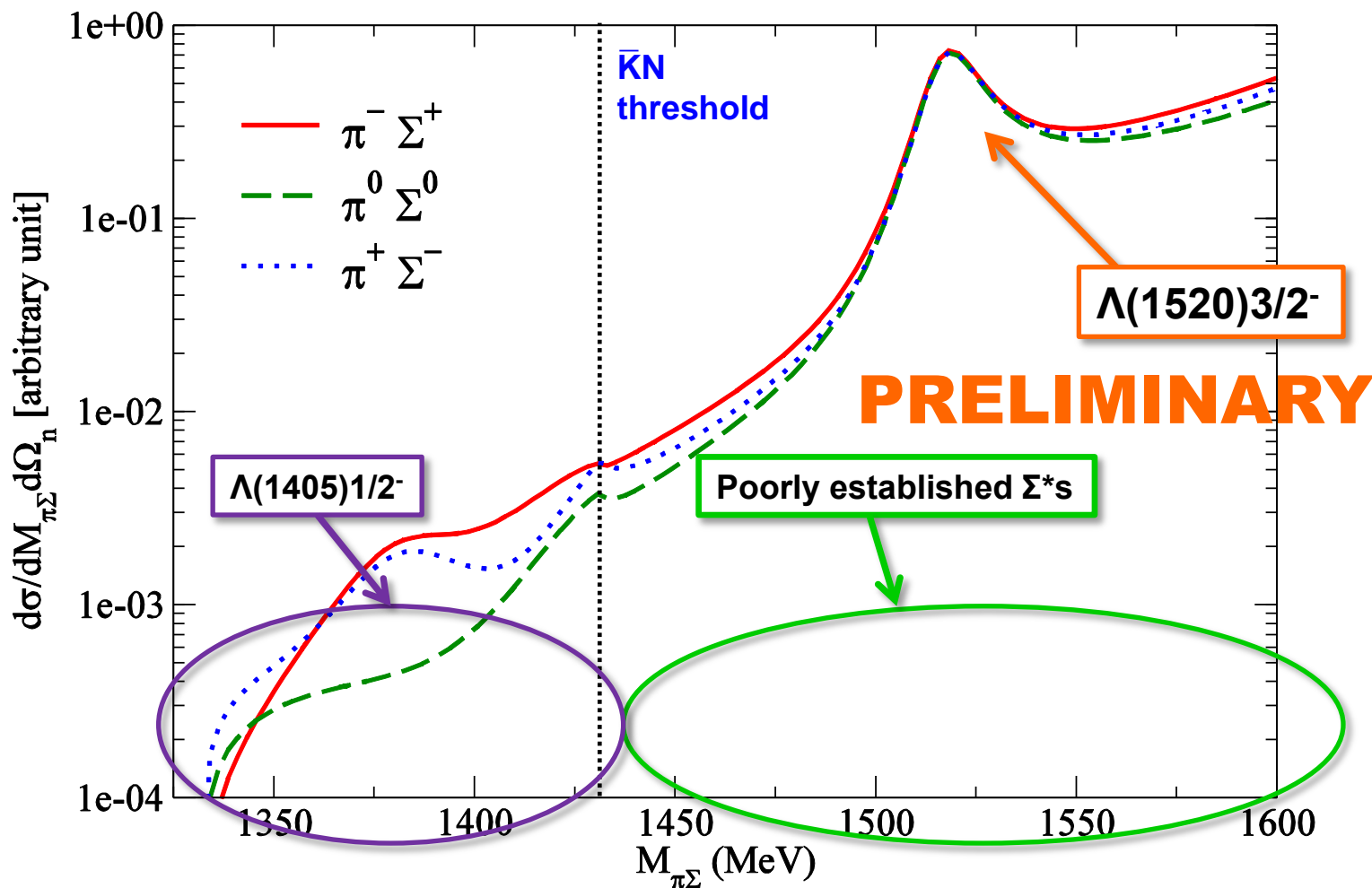
Same kinematics
as J-PARC E31



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($P_{K^-} = 1$ GeV, $\theta_n = 0$ deg.)

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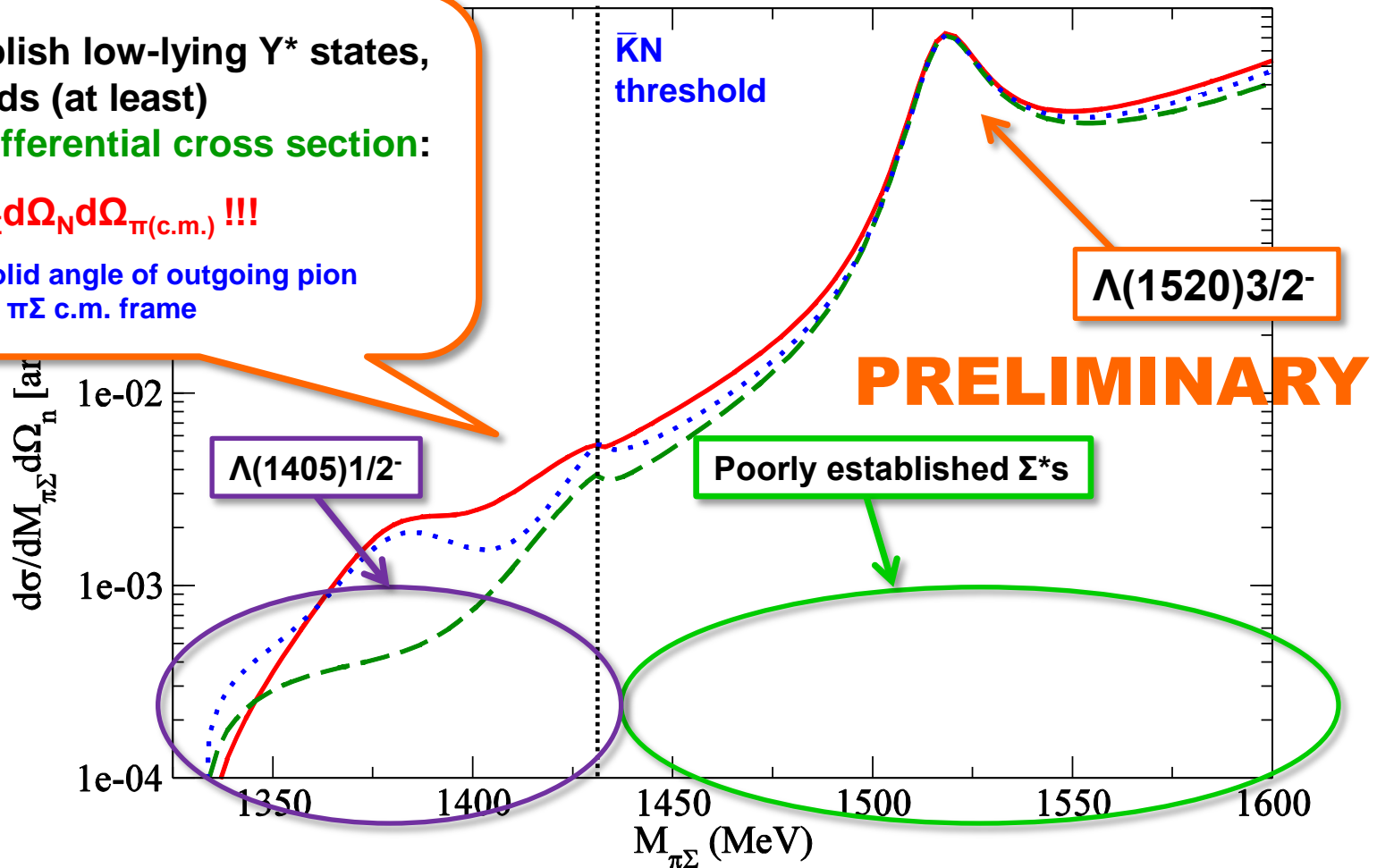
Same kinematics
as J-PARC E31

To establish low-lying Y^* states,
one needs (at least)

5-fold differential cross section:

$d\sigma/dM_{\pi\Sigma}d\Omega_Nd\Omega_{\pi(c.m.)} !!!$

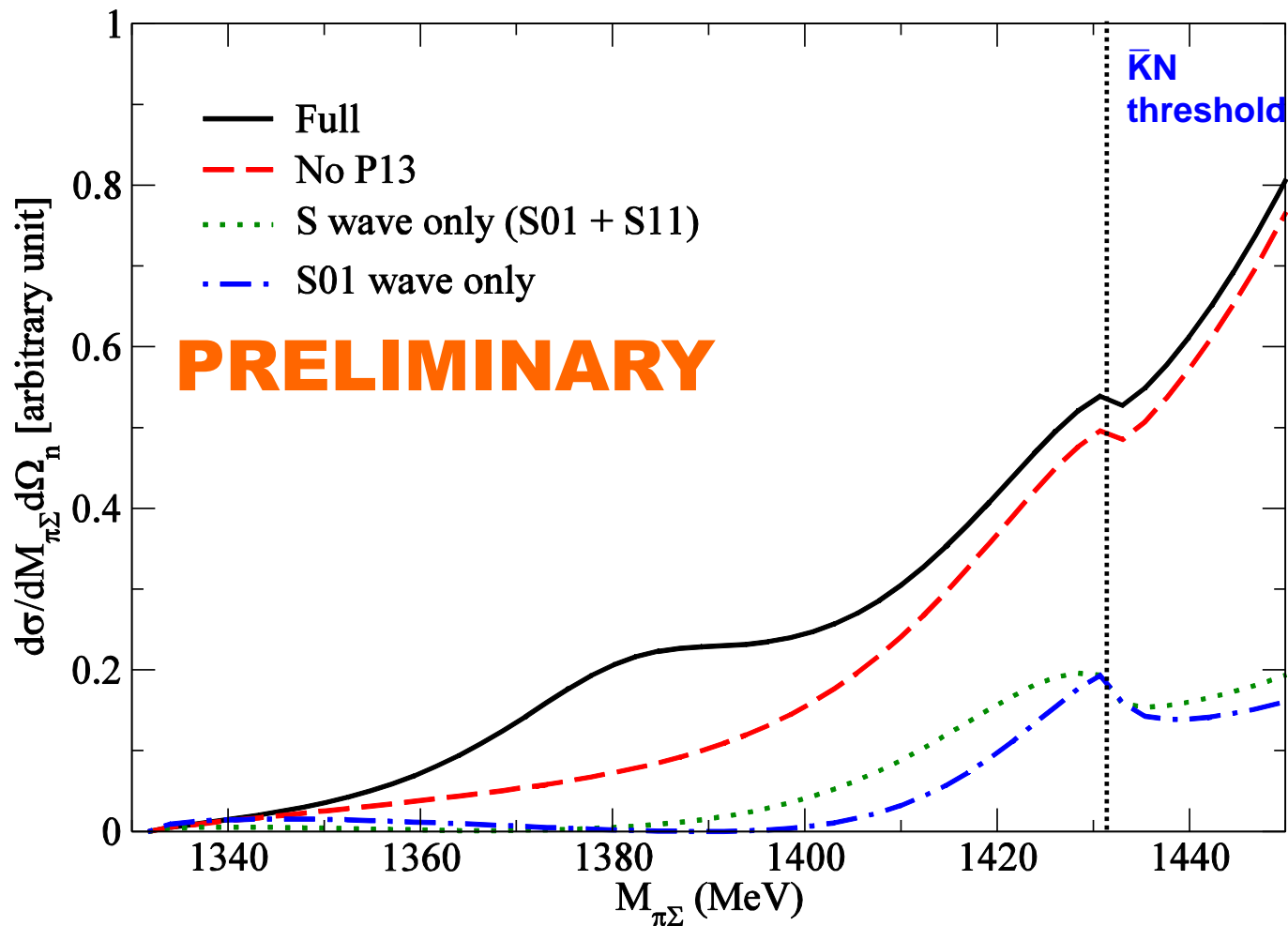
$\Omega_{\pi(c.m.)}$ = solid angle of outgoing pion
in $\pi\Sigma$ c.m. frame



Results (IMPULSE TERM ONLY !!)

$d\sigma/dM_{\pi\Sigma}d\Omega_n$ for $K^- d \rightarrow \pi^- \Sigma^+ n$
($P_{K^-} = 1$ GeV, $\theta_n = 0$ deg.)

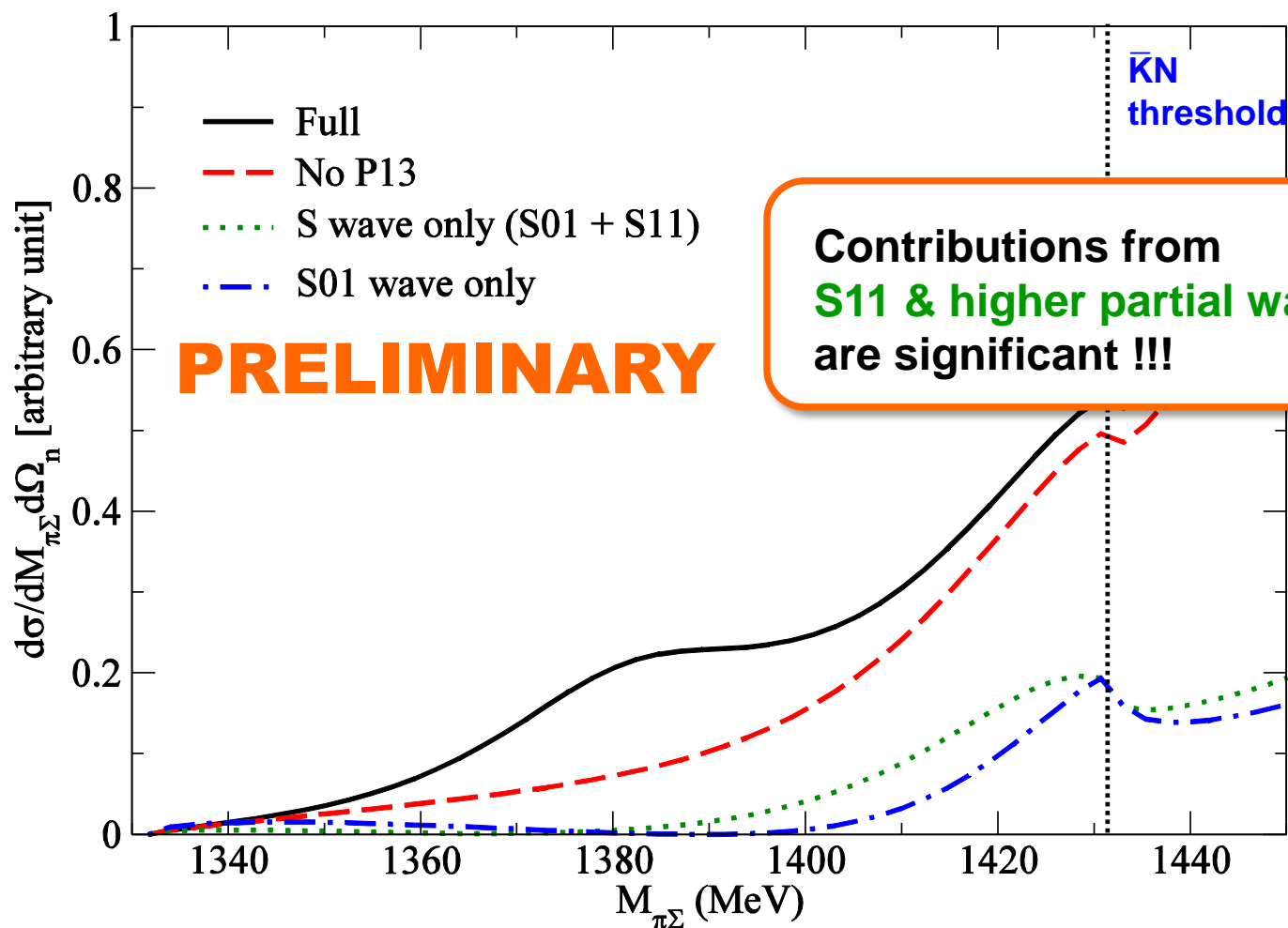
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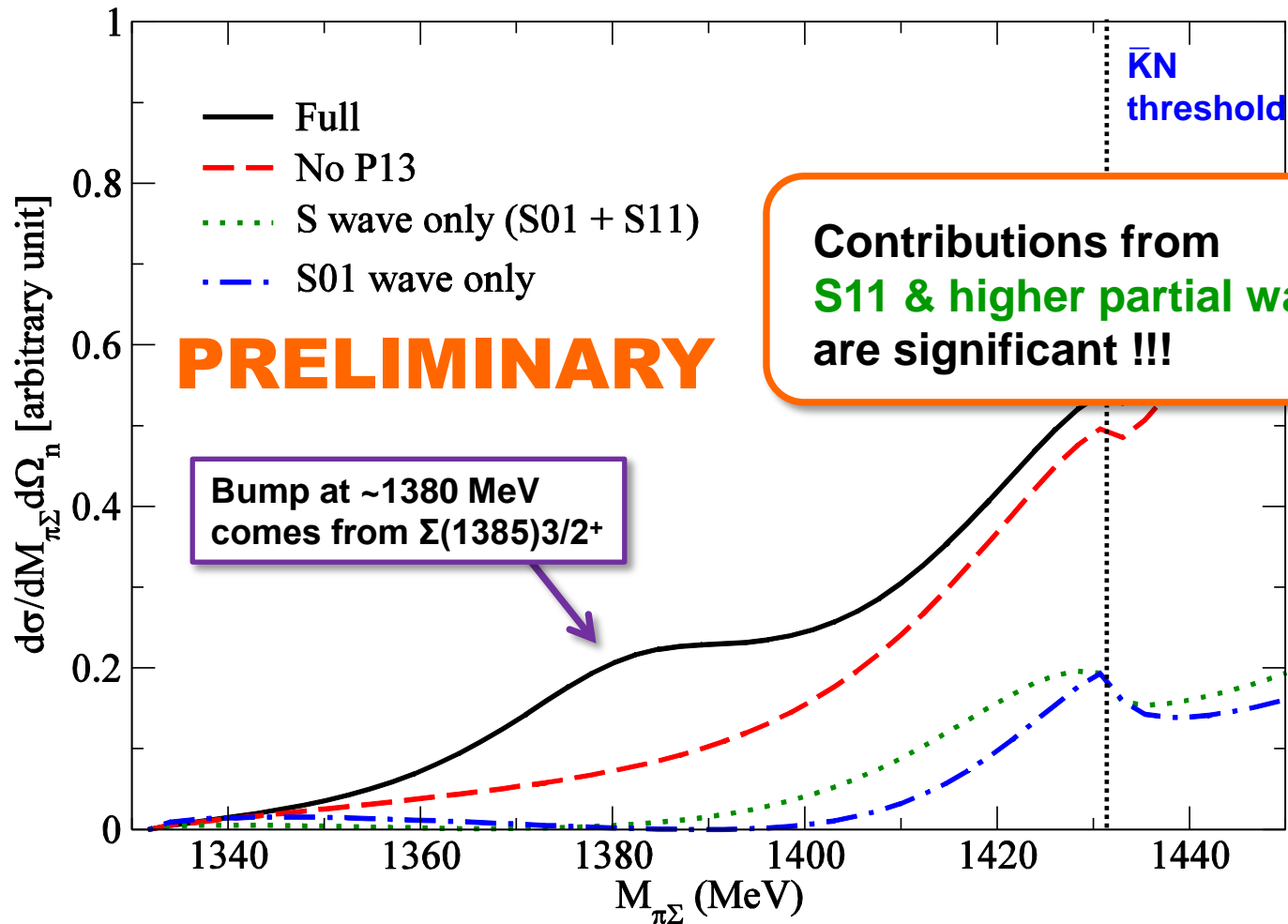


Contributions from
S11 & higher partial waves
are significant !!!

Results (IMPULSE TERM ONLY !!)

$d\sigma/dM_{\pi\Sigma}d\Omega_n$ for $K^- d \rightarrow \pi^- \Sigma^+ n$
($P_{K^-} = 1$ GeV, $\theta_n = 0$ deg.)

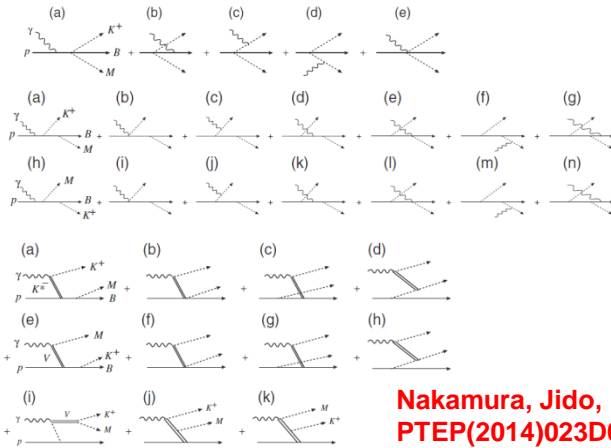
Same kinematics
as J-PARC E31



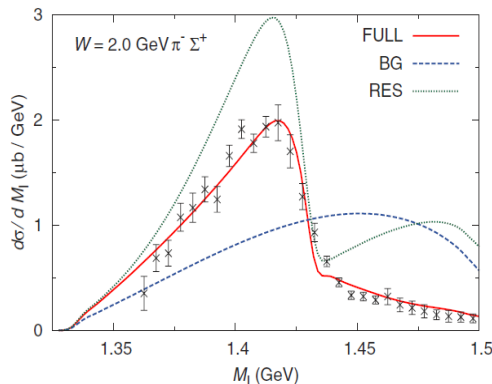
How we study the region below the $\bar{K}N$ threshold ?

$\gamma p \rightarrow K^+ \pi^- \Sigma$ @CLAS

At the CLAS energy, **many** production processes contribute and sizably affect mass distributions as **backgrounds**.



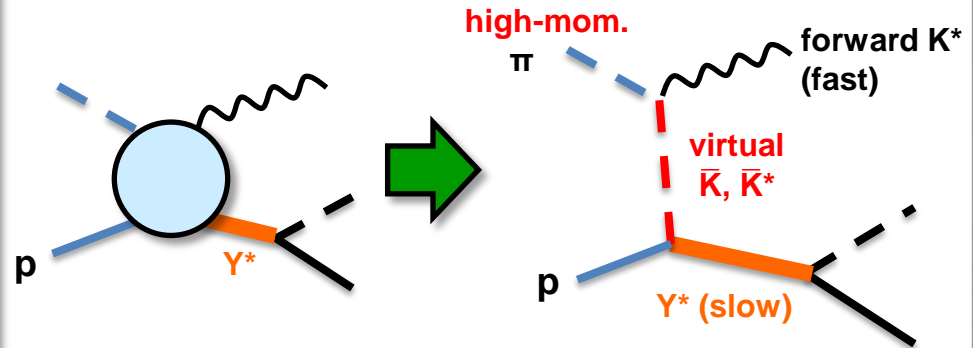
Nakamura, Jido,
PTEP(2014)023D01



Large model/analysis
dependence from
complicated production
processes.

→ Makes unambiguous
determination
of $\Lambda(1405)$ difficult.

Forward $p(\pi, K^*)X$ reactions with
high-momentum pion beam (→ J-PARC E50)



- For forward K^* (small t), the processes are dominated by diffractive t -channel exchange processes.
- We DO have fully unitarized $\bar{K}N \rightarrow MB$ and $\bar{K}^*N \rightarrow MB$ half off-shell amplitudes !!
- 12 GeV JLab can do a similar measurement by replacing incident π by high-energy photon.

➤ Useful also for
determining
low-lying Σ^* resonances

Low-lying Σ^* resonances(PDG)

$\Sigma(1193)$	$1/2+$	****
$\Sigma(1385)$	$3/2+$	****
$\Sigma(1480)$		*
$\Sigma(1560)$		**
$\Sigma(1580)$	$3/2-$	*
$\Sigma(1620)$	$1/2-$	**
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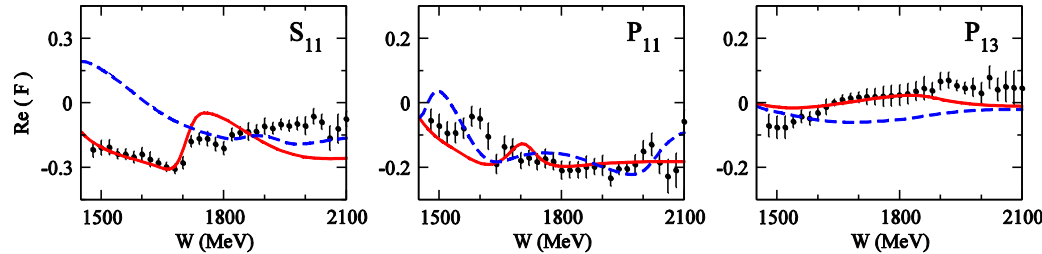
?

New data can eliminate analysis dependence ??

Extracted $\bar{K}N \rightarrow \pi\Lambda$ amplitudes

HK, Nakamura, Lee, Sato,
PRC90(2014)065204;92(2015)025205

Real parts



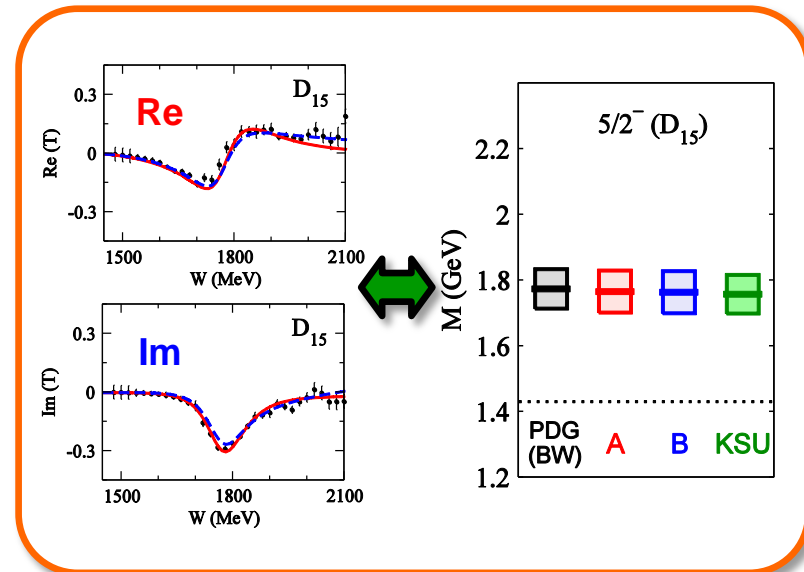
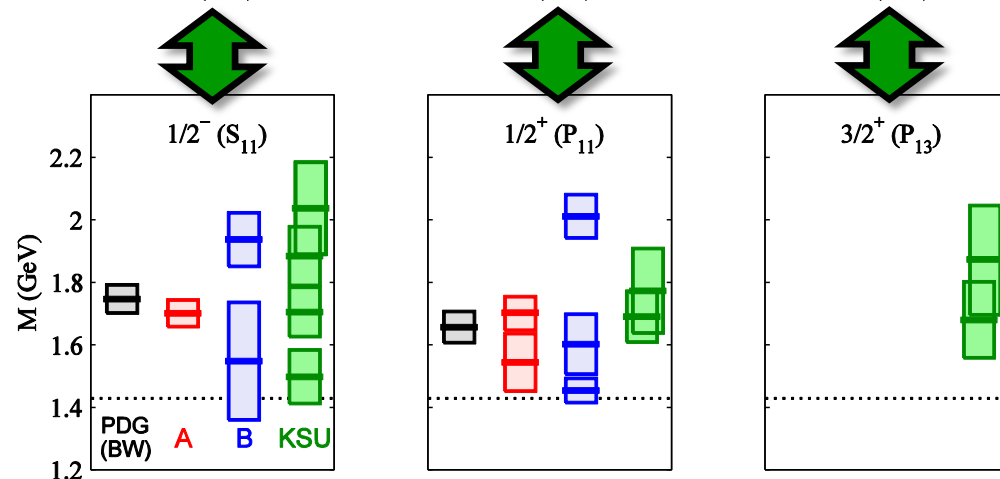
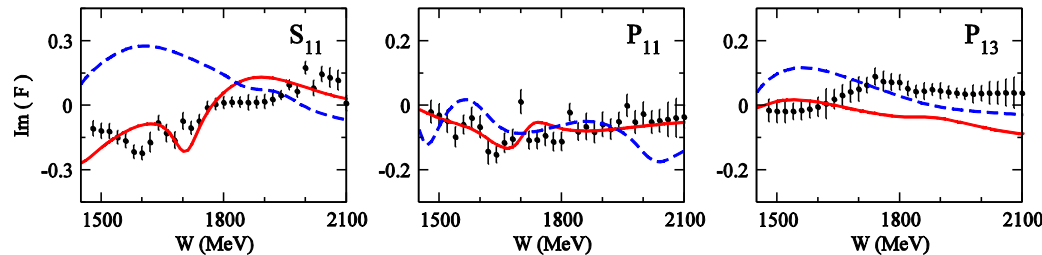
L_{12J} : $L = S, P, \dots$; $I =$ isospin;
 $J =$ Total angular mom.

Red : Model A

Blue: Model B

Circles: KSU single-energy solution
[PRC88(2013)035204]

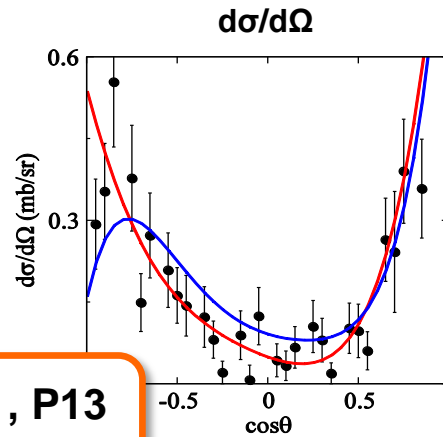
Imaginary parts



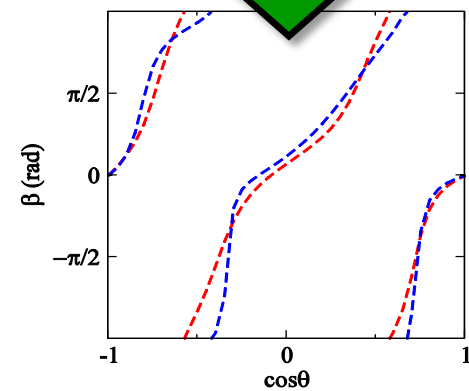
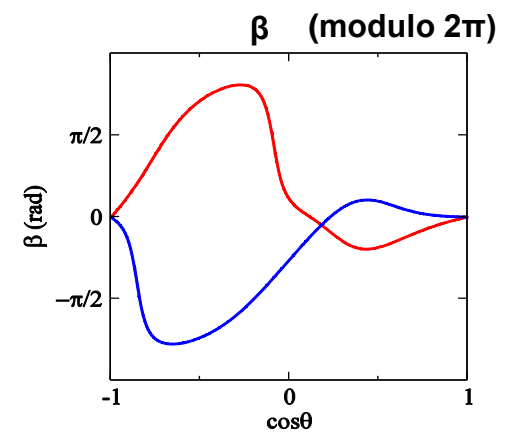
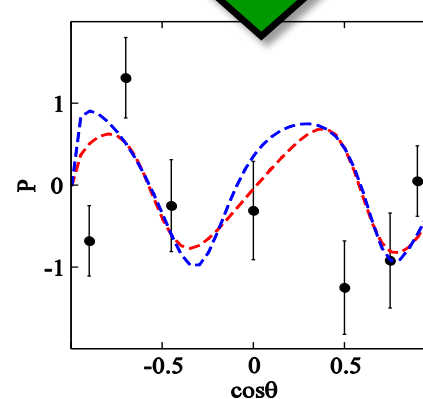
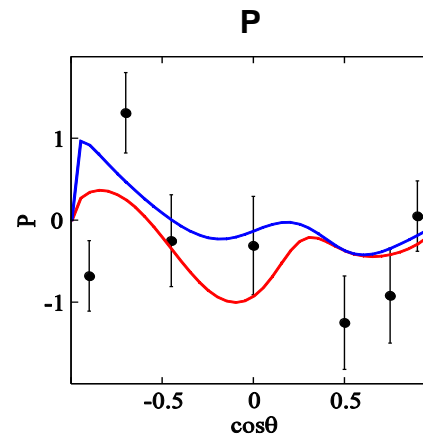
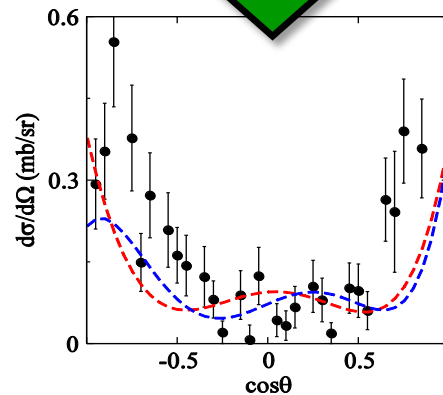
New data can eliminate analysis dependence ??

$K^-p \rightarrow \pi^0 \Lambda$ observables @ $W = 1700$ MeV

Red: Model A
Blue: Model B

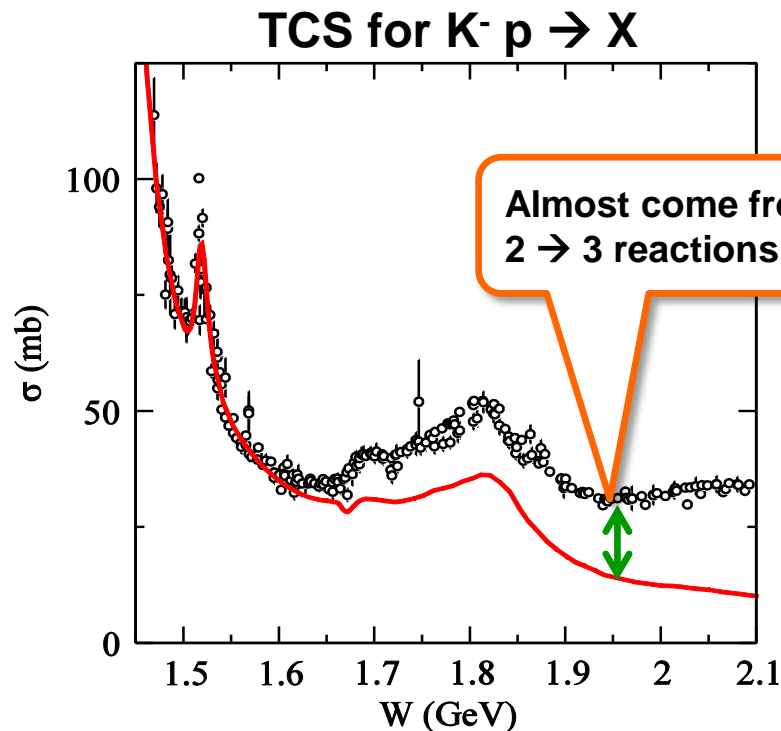


S11, P11, P13
waves off



- At this energy, the difference between Models A & B **mostly comes from S11, P11, P13 waves**.
- High statistics data (of **P and β in particular**) will reduce significantly the analysis dependence !!!

Importance of $2 \rightarrow 3$ reactions: Dominance of cross sections at high W



— Sum of $K^- p \rightarrow \bar{K}N, \pi\Sigma, \pi\Lambda, \eta\Lambda, K\Xi$
(Computed with Model A)

TCS for $2 \rightarrow 3$ reactions
($K^- p \rightarrow \pi\pi\Lambda, \pi\bar{K}N, \dots$):

- significant above $W \sim 1.7$ GeV.
- even larger than the $2 \rightarrow 2$ TCS above $W \sim 1.9$ GeV !!

Effects of **3-body channels** on Y^* resonance parameters are expected to be sizable.



However, at present **essentially no differential cross section data** are available for $2 \rightarrow 3$ reactions that can be used for **detailed partial wave analyses** !!

Extracted scattering lengths and effective ranges

HK, Nakamura, Lee, Sato, PRC90(2014)065204

Scattering length and effective range

	Model A		Model B	
	$I = 0$	$I = 1$	$I = 0$	$I = 1$
$a_{\bar{K}N}$ (fm)	$-1.37 + i0.67$	$0.07 + i0.81$	$-1.62 + i1.02$	$0.33 + i0.49$
$a_{\eta\Lambda}$ (fm)	$1.35 + i0.36$	-	$0.97 + i0.51$	-
$a_{K\Xi}$ (fm)	$-0.81 + i0.14$	$-0.68 + i0.09$	$-0.89 + i0.13$	$-0.83 + i0.03$
$r_{\bar{K}N}$ (fm)	$0.67 - i0.25$	$1.01 - i0.20$	$0.74 - i0.25$	$-1.03 + i0.19$
$r_{\eta\Lambda}$ (fm)	$-5.67 - i2.24$	-	$-5.82 - i3.32$	-
$r_{K\Xi}$ (fm)	$-0.01 - i0.33$	$-0.42 - i0.49$	$0.13 - i0.20$	$-0.22 - i0.11$

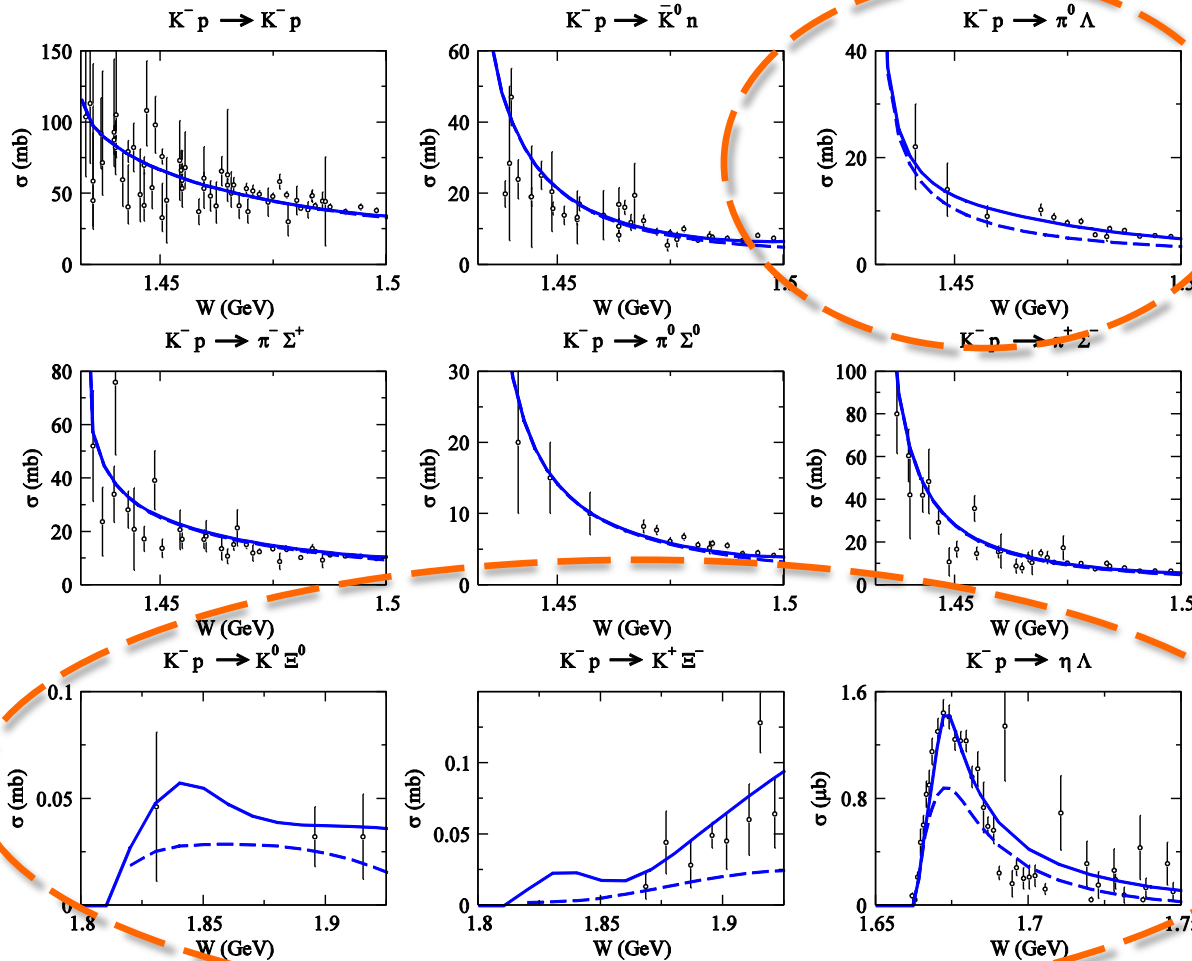
$$a_{K-p} = -0.65 + i0.74 \text{ fm (Model A)}$$

$$a_{K-p} = -0.65 + i0.76 \text{ fm (Model B)}$$

S-wave dominance ??

$K^- p \rightarrow MB$ total cross sections near threshold

Model B



Solid: Full
Dashed: S wave only

For $K^- p \rightarrow \pi \Lambda, \eta \Lambda, K \Xi$,
higher partial waves
visibly contribute
to the cross sections
even in the threshold
region.

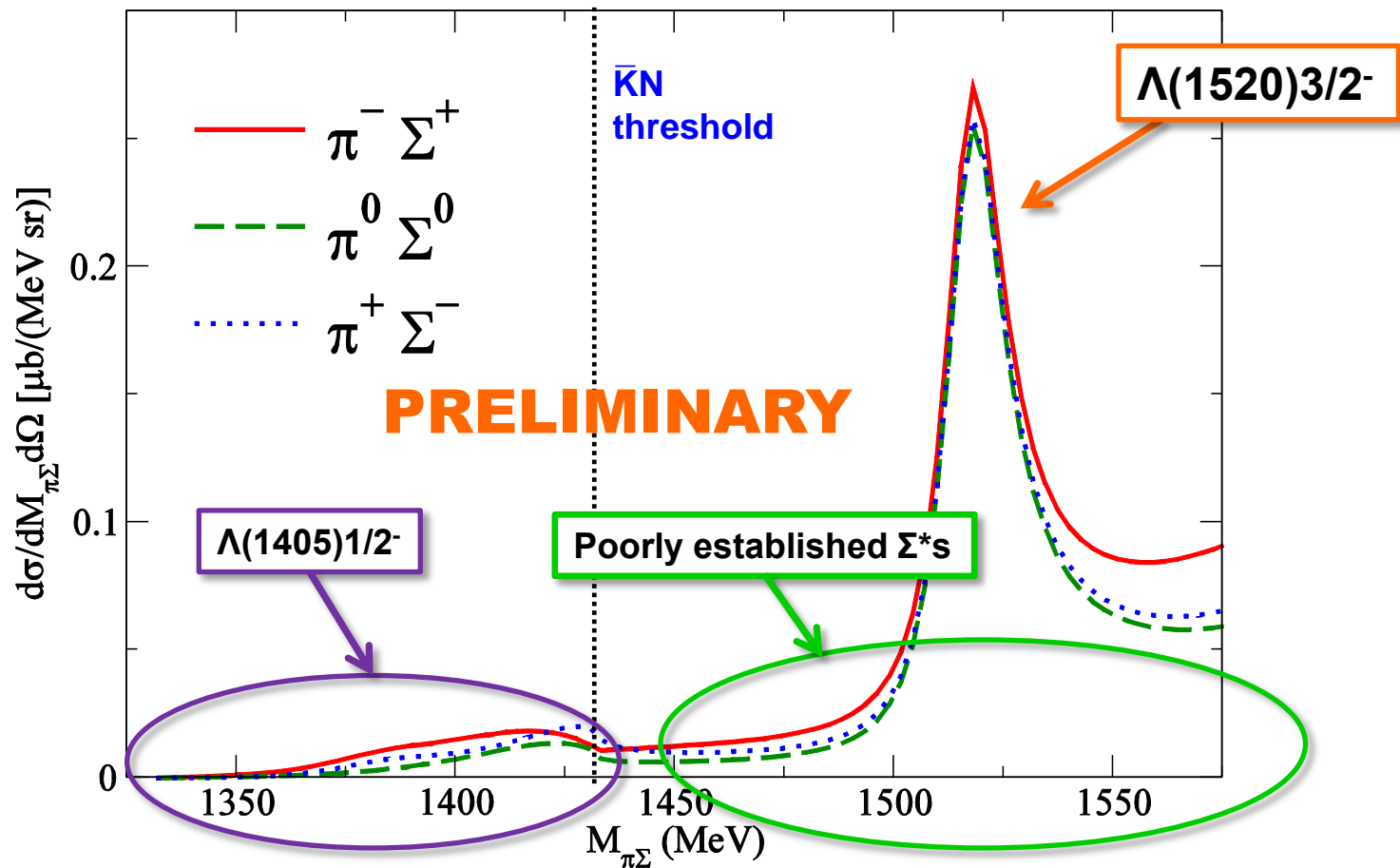
→ consistent with the observation in
Jackson et al., PRC91(2015)065208



Naïve expectation for
S-wave dominance
near the threshold
sometimes does not hold !!

Results (IMPULSE TERM ONLY !!)

$d\sigma/dM_{\pi\Sigma}d\Omega_n$ for $K^- d \rightarrow (\pi \Sigma)_0 n$
($P_{K^-} = 600$ MeV, $\theta_n = 0$ deg.)



Results (IMPULSE TERM ONLY !!)

$$d\sigma/dM_{\pi\Sigma}d\Omega_n \text{ for } K^- d \rightarrow (\pi\Sigma)_0 n$$

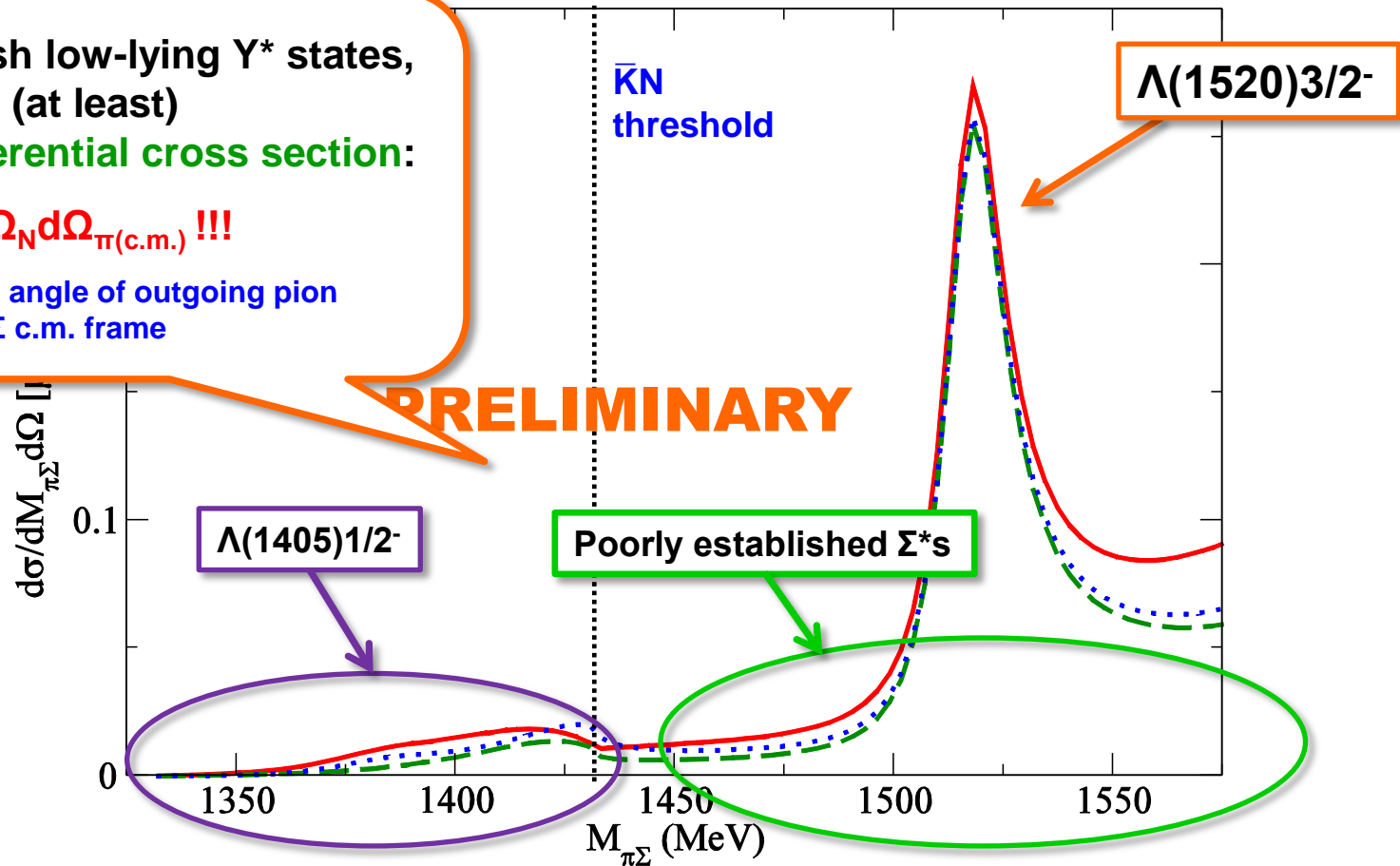
$$(P_{K^-} = 600 \text{ MeV}, \theta_n = 0 \text{ deg.})$$

To establish low-lying Y^* states,
one needs (at least)

5-fold differential cross section:

$$d\sigma/dM_{\pi\Sigma}d\Omega_Nd\Omega_{\pi(c.m.)} !!!$$

$\Omega_{\pi(c.m.)}$ = solid angle of outgoing pion
in $\pi\Sigma$ c.m. frame



Results (IMPULSE TERM ONLY !!)

$d\sigma/dM_{\pi\Sigma}d\Omega_N$ for $K^- d \rightarrow \pi^- \Sigma^+ n$
($P_{K^-} = 600$ MeV, $\theta_n = 0$ deg.)

