

Dibaryons with and without Heavy Flavor

Makoto Oka

*Tokyo Institute of Technology
and
ASRC, JAEA*

**International Workshop on Progress on
J-PARC Hadron Physics
March 03, 2016**

Dibaryons with and without **Strangeness**

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Why dibaryons?

- # **Dibaryon = (Baryon+ Baryon) = (6 quarks)**
- # **A “loosely bound” dibaryon (ex. deuteron) is a hadron molecule, which provides important information on the (long-range) baryon-baryon interaction.**
- # **A compact dibaryon could be a 6-quark state, which is related to the short-distance baryon-baryon interaction.**
- # **The dibaryon is a “good” resonance because there is no annihilation channel. The only outgoing channels are two baryons (neglecting meson emissions).**
- # **The quark model symmetries give simple and robust guidelines on the existence of compact dibaryons.**

Pauli principles and Spin dependence

Symmetries

- Recent Lattice QCD calculations (HALQCD) have confirmed that the short-range baryon-baryon interactions follow the quark model symmetry and dynamics.
- Two important effects are given by
 - Fermi-Dirac statistics among quarks (Pauli effect)
 - Spin dependent force: Color-magnetic interaction (CMI)
- Symmetries of internal degrees of freedom
 - spin \times flavor \times color \times orbital motion
 - $\underline{SU(2)_s \times SU(N_f)_f} \times SU(3)_c \times O(3)$
 - $SU(2N_f)_{\underline{sf}} \times SU(3)_c \times O(3)$

Pauli effect

- # $SU_{sf}(6) \supset SU(2)_s \times SU(3)_f$ symmetry of two-baryon states:
two $56 [3] = (8, 1/2) + (10, 3/2)$ baryons.

SU(6)

$$\begin{array}{ccccccc}
 & & \text{Sym} & & \text{Antisym} & & \\
 & \text{---} & \text{---} & \text{---} & \text{---} & \text{---} & \\
 [3] & \times & [3] & = & [6] & + & [42] & + & [51] & + & [33] \\
 56 & & 56 & & & & \text{odd L} & & \text{even L} & &
 \end{array}$$

Strong repulsion due to the Pauli Exclusion Principle

$$\begin{array}{ccccccc}
 L=0 & & & & & & \\
 [6] & \times & [51] & \times & [222] & \neq & [111111] \\
 \text{orbital} & & \text{flavor} & & \text{color} & & \text{Forbidden} \\
 & & \text{spin} & & \text{singlet} & &
 \end{array}$$

The totally symmetric orbital states are forbidden in the $[51]$ flavor-spin states.

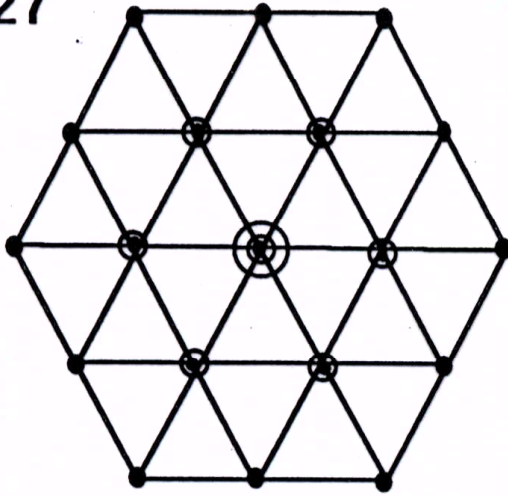
Strong short-range repulsion appears when the [6] symmetric orbital state is forbidden by the Pauli principle.

taken from D. Sc Thesis by M.O. (1980)

L	SU(4)	BB' (S,I)
even	{33}	$\Delta\Delta (3,0), \Delta\Delta (0,3)$
	{51} forbidden	$\Delta\Delta (3,2), \Delta\Delta (2,3), N\Delta (2,2), N\Delta (1,1)$
	{33} + {51}	$\Delta\Delta + N\Delta (2,1), \Delta\Delta + N\Delta (1,2)$ $NN + \Delta\Delta (1,0), NN + \Delta\Delta (0,1)$
odd	{6} forbidden	$\Delta\Delta (3,3)$
	{42}	$\Delta\Delta (3,1), \Delta\Delta (1,3), \Delta\Delta (2,0), \Delta\Delta (0,2)$ $N\Delta (2,1), N\Delta (1,2)$
	{6} + {42}	$N\Delta + \Delta\Delta (2,2), NN + \Delta\Delta (0,0)$
	{6} + {42} ²	$NN + N\Delta + \Delta\Delta (1,1)$

$B_8 B_8$ Flavor Symmetric \rightarrow singlet even/triplet odd

27



$NN(I=1)$

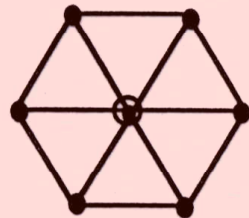
$\Sigma N(I=3/2), \Sigma N-\Lambda N(I=1/2)$

$\Sigma\Sigma(I=2), \Xi N-\Sigma\Sigma-\Sigma\Lambda(I=1), \Xi N-\Sigma\Sigma-\Lambda\Lambda(I=0)$

$\Xi\Sigma(I=3/2), \Xi\Sigma-\Xi\Lambda(I=1/2)$

$\Xi\Xi(I=1)$

8_s



$\Sigma N-\Lambda N(I=1/2)$

$\Xi N-\Sigma\Lambda(I=1), \Xi N-\Sigma\Sigma-\Lambda\Lambda(I=0)$

$\Xi\Sigma-\Xi\Lambda(I=1/2)$

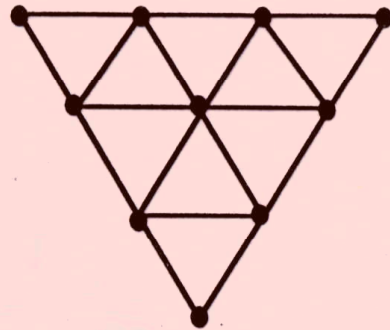
1



$\Xi N-\Sigma\Sigma-\Lambda\Lambda(I=0)$

$B_8 B_8$ Flavor Antisymmetric \rightarrow triplet even/singlet odd

10



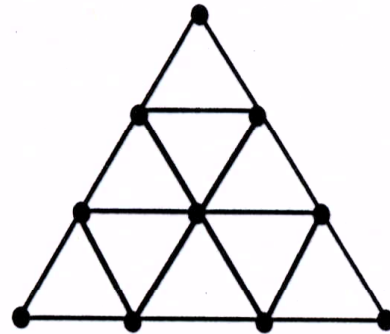
$\Sigma N(I=3/2)$

$\Xi N - \Sigma \Sigma - \Sigma \Lambda(I=1)$

$\Xi \Sigma - \Xi \Lambda(I=1/2)$

$\Xi \Xi(I=0)$

$\overline{10}$



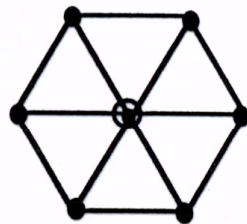
$NN(I=0)$

$\Sigma N - \Lambda N(I=1/2)$

$\Xi N - \Sigma \Lambda(I=1)$

$\Xi \Sigma(I=3/2)$

8_A



$\Sigma N - \Lambda N(I=1/2)$

$\Xi N - \Sigma \Sigma - \Sigma \Lambda(I=1), \Xi N(I=0)$

$\Xi \Sigma - \Xi \Lambda(I=1/2)$

Pauli effect

HAL QCD data are consistent with the quark Pauli effects.

S=0

1	[33]	Allowed, $\Lambda\Lambda + N\Xi + \Sigma\Sigma \rightarrow H$
8 _s	[51]	Pauli forbidden , ΣN (I=1/2, S=0)
27	[33], [51]	55% Allowed, NN 1S_0

S=1

8 _a	[33], [51]	
10	[33], [51]	Almost forbidden , ΣN (I=3/2, S=1)
10*	[33], [51]	NN 3S_1

Pauli effect

T. Inoue et al., (HAL QCD) PTP 124, 591 (2010)

HAL QCD data are compared

$S=0$

1

[33]

All

8_s

[51]

Pauli

27

[33], [51]

55%

$S=1$

8_a

[33], [51]

10

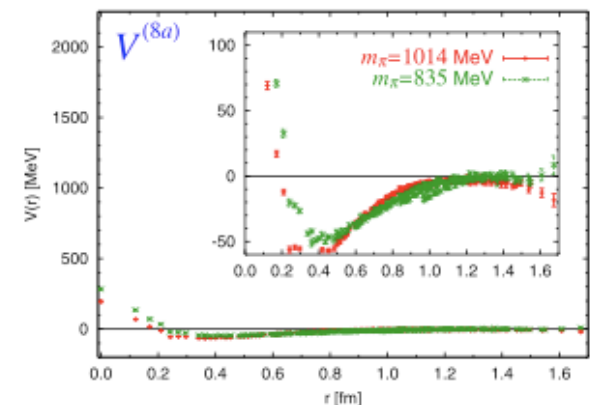
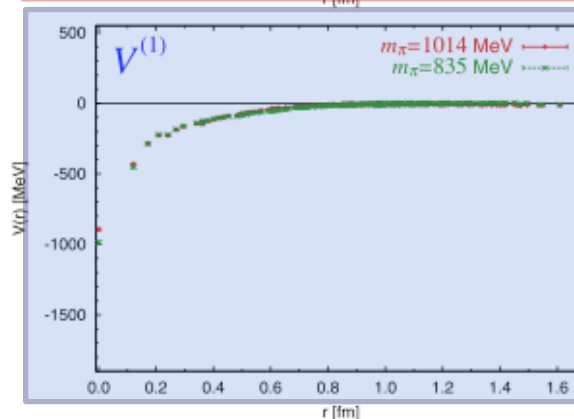
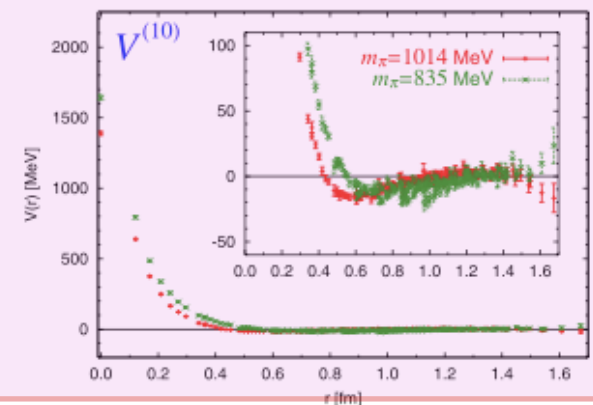
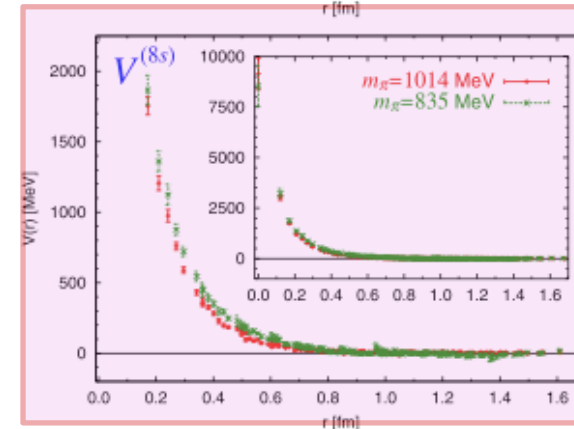
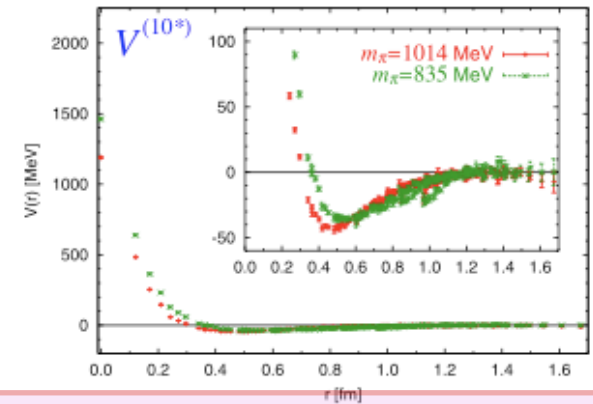
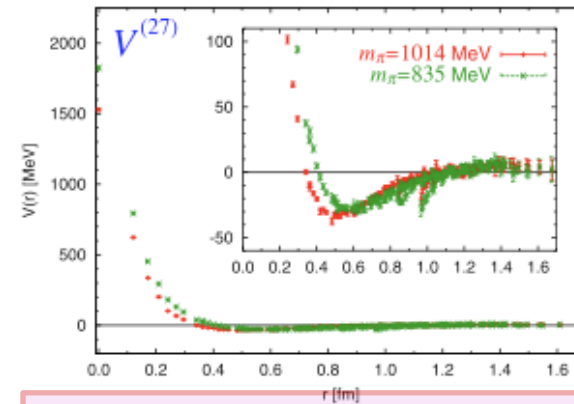
[33], [51]

All

10^*

[33], [51]

NN



Spin dependence

Spin-spin interaction aka Color-Magnetic Interaction (CMI)

$$V_{\text{CMI}} = -\alpha \sum_{i < j} (\vec{\lambda}_i \cdot \vec{\lambda}_j) (\vec{\sigma}_i \cdot \vec{\sigma}_j) f(r_{ij}) \quad f(r_{ij}) \sim \delta(r_{ij})$$

prefers symmetric color-spin states

$$\langle V_{\text{CMI}} \rangle_{(0s)^N} = \alpha \langle f(r) \rangle_{0s} \Delta_{\text{CM}} = V_0 \Delta_{\text{CM}}$$

$$\Delta_{\text{CM}} \equiv \langle - \sum_{i < j} (\vec{\lambda}_i \cdot \vec{\lambda}_j) (\vec{\sigma}_i \cdot \vec{\sigma}_j) \rangle$$

$$\Delta_{\text{CM}} = 8N - 2C_2[SU(6)_{cs}] + \frac{4}{3}S(S+1) + C_2[SU(3)_c]$$

$$C_2[SU(g)]([f_1, f_2, \dots, f_g]) = \sum_i f_i(f_i - 2i + g + 1) - \frac{N^2}{g}$$

$$C_2[\text{singlet}] = 0$$

Spin dependence

- ‡ CMI prefers color-spin symmetric states, i.e. flavor antisymmetric states.

$$\Delta_{\text{CM}} = 8N - 2C_2[SU(6)_{cs}] + \frac{4}{3}S(S+1) + C_2[SU(3)_c]$$

$$\Delta_{\text{CM}}(\mathbf{10}) - \Delta_{\text{CM}}(\mathbf{8}) = 8 - (-8) = 16$$

$$M(\Delta) - M(N) = 16V_0 \sim 300 \text{ MeV}$$

$$V_0 \sim 300/16 \sim 19 \text{ MeV}$$

$$\Delta_{\text{CM}}(H) - 2\Delta_{\text{CM}}(\Lambda) = -24 - 2(-8) = -8 \quad \mathbf{H} (\Lambda\Lambda + \mathbf{N}\Xi + \Sigma\Sigma, \mathbf{S}=0)$$

$$\Delta_{\text{CM}}(D_\Delta) - 2\Delta_{\text{CM}}(\Delta) = 16 - 2 \times 8 = 0 \quad \mathbf{D}_\Delta (\Delta\Delta, \mathbf{I}=0, \mathbf{S}=3)$$

H dibaryon

H dibaryon

$H = u^2d^2s^2$ ($S = -2$, $J=0^+$ $I=0$) predicted by Jaffe (1977)

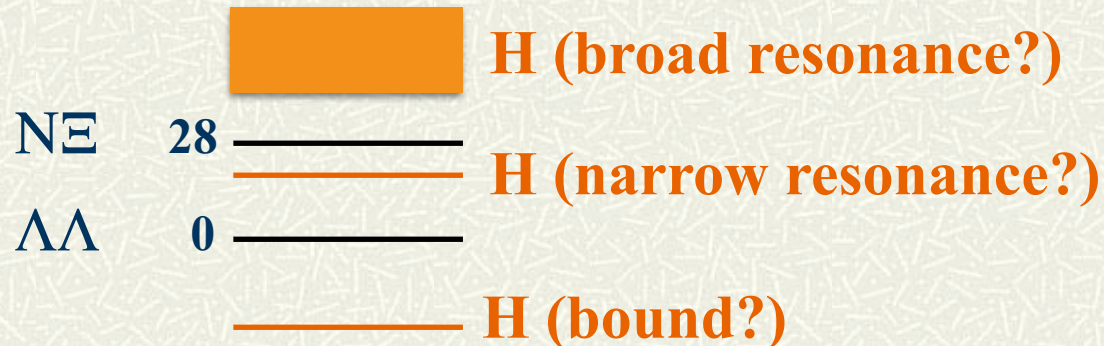
CMI prefers

symmetric color-spin state \Leftrightarrow antisymmetric flavor state

Most favored state is the flavor singlet state.

(MeV)
 $\Sigma\Sigma$ 150

$$|F = 1\rangle = -\sqrt{\frac{1}{8}}|\Lambda\Lambda\rangle + \sqrt{\frac{4}{8}}|N\Xi\rangle + \sqrt{\frac{3}{8}}|\Sigma\Sigma\rangle$$



H dibaryon

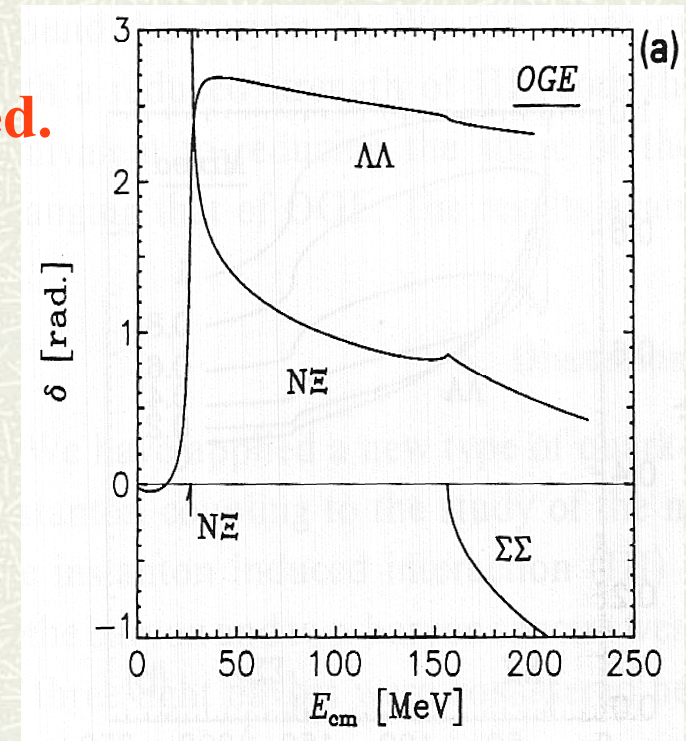
Quark cluster model approach to the coupled channel $\Lambda\Lambda$, $N\Xi$, $\Sigma\Sigma$ system, with the linear + OGE potential for quarks.

MO, K. Shimizu, K. Yazaki (1983)

- The B_8B_8 ($F=1$) channel is **Pauli super-allowed**.
- There appears a very sharp resonance just below the $N\Xi$ threshold.
- Additional long range attraction will form a bound state below the $\Lambda\Lambda$ threshold.

S. Takeuchi and MO (1991)

- The instanton induced interaction yields 3-body repulsive force to H, resulting no bound state.



$$|\text{Singlet}\rangle = \sqrt{\frac{1}{8}}|\Lambda\Lambda\rangle + \sqrt{\frac{4}{8}}|N\Xi\rangle - \sqrt{\frac{3}{8}}|\Sigma\Sigma\rangle$$

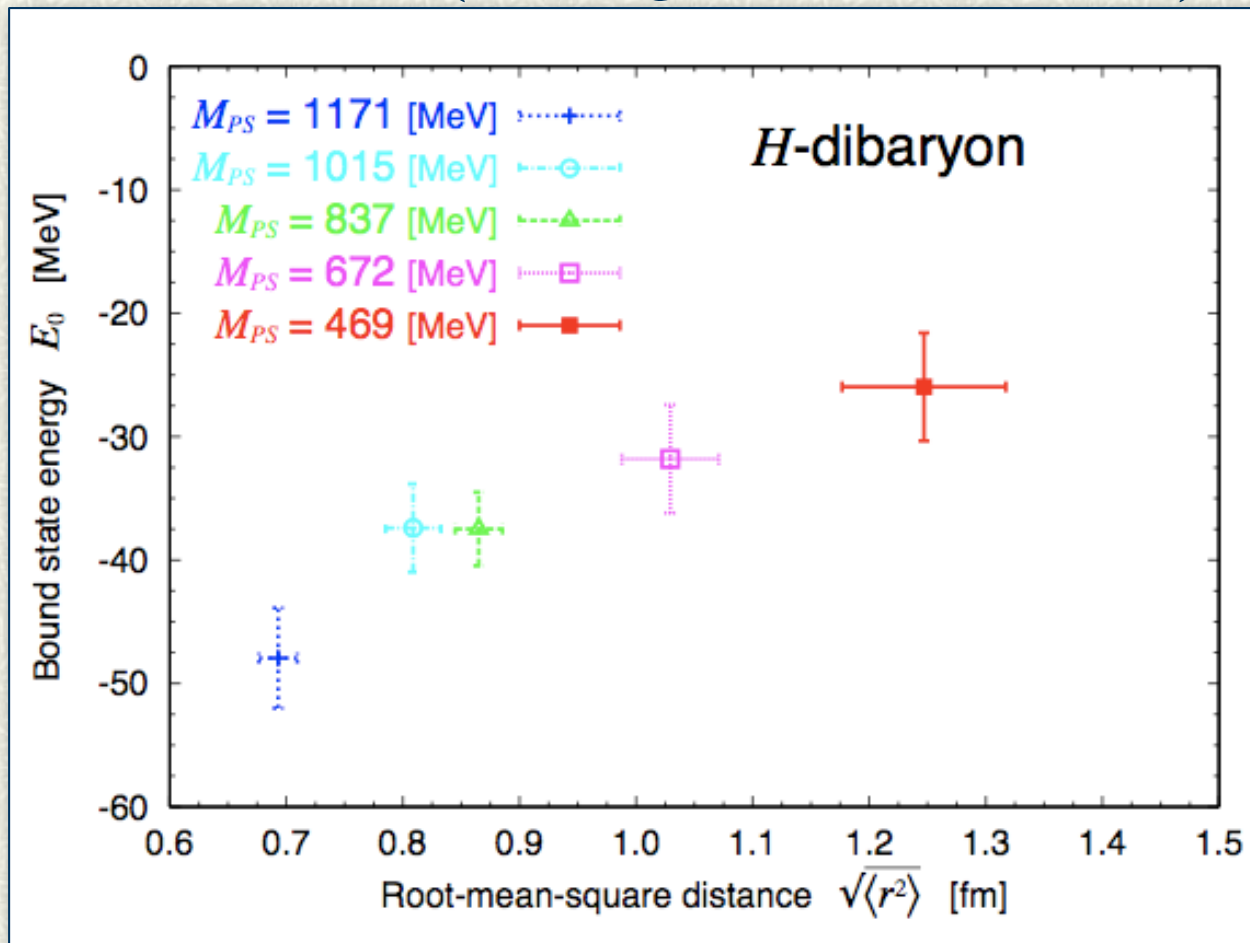
H dibaryon on Lattice

- # **New Lattice QCD calculations of H dibaryon**
- **Bound H di-baryon in Flavor SU(3) Limit of Lattice QCD**
Takashi Inoue (HAL QCD Collaboration)
PRL 106, 162002 (2011)
- **Evidence for a Bound H di-baryon from Lattice QCD**
S. R. Beane et al. (NPLQCD Collaboration)
PRL 106, 162001 (2011)

H dibaryon on Lattice

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○

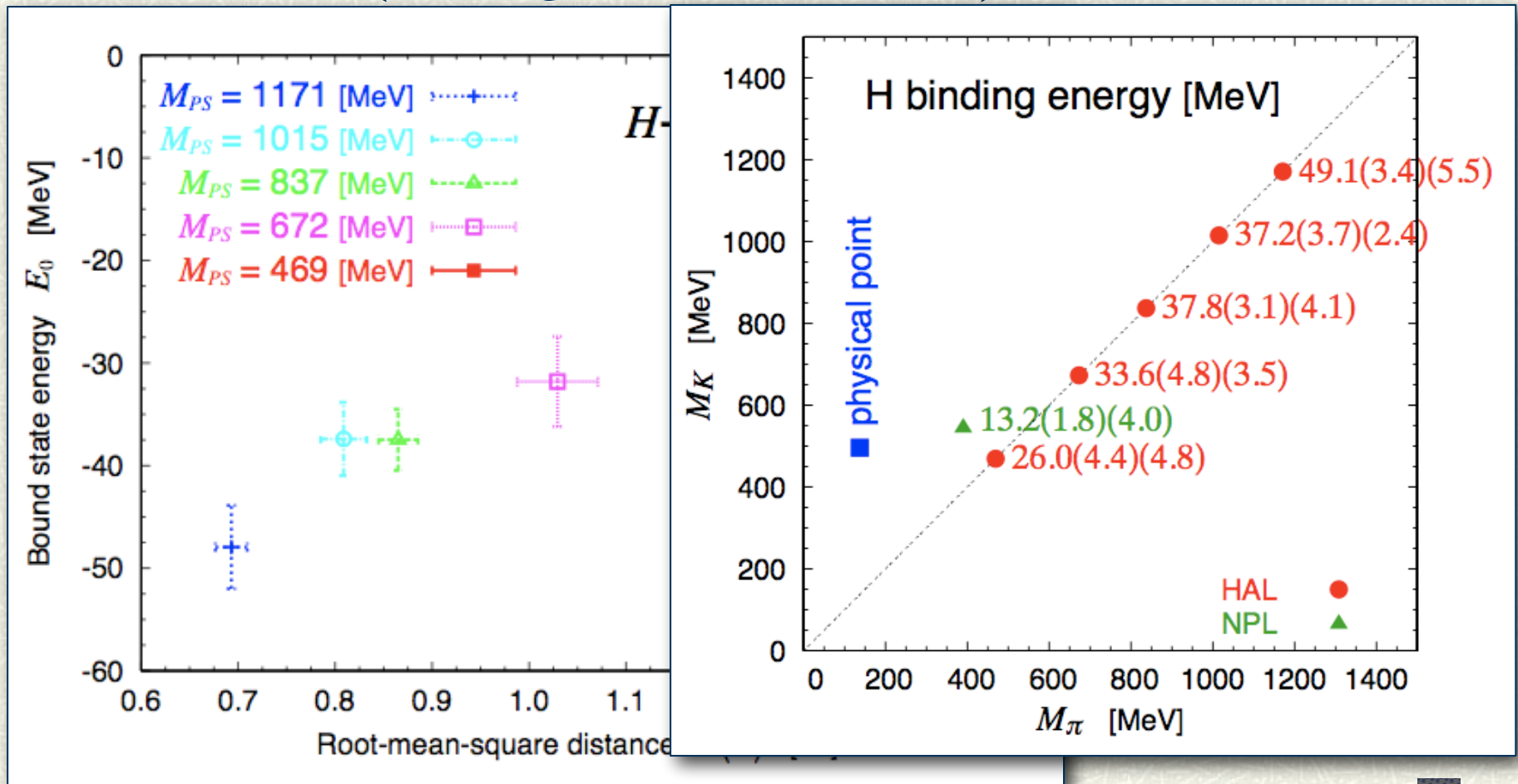


tice QCD

H dibaryon on Lattice

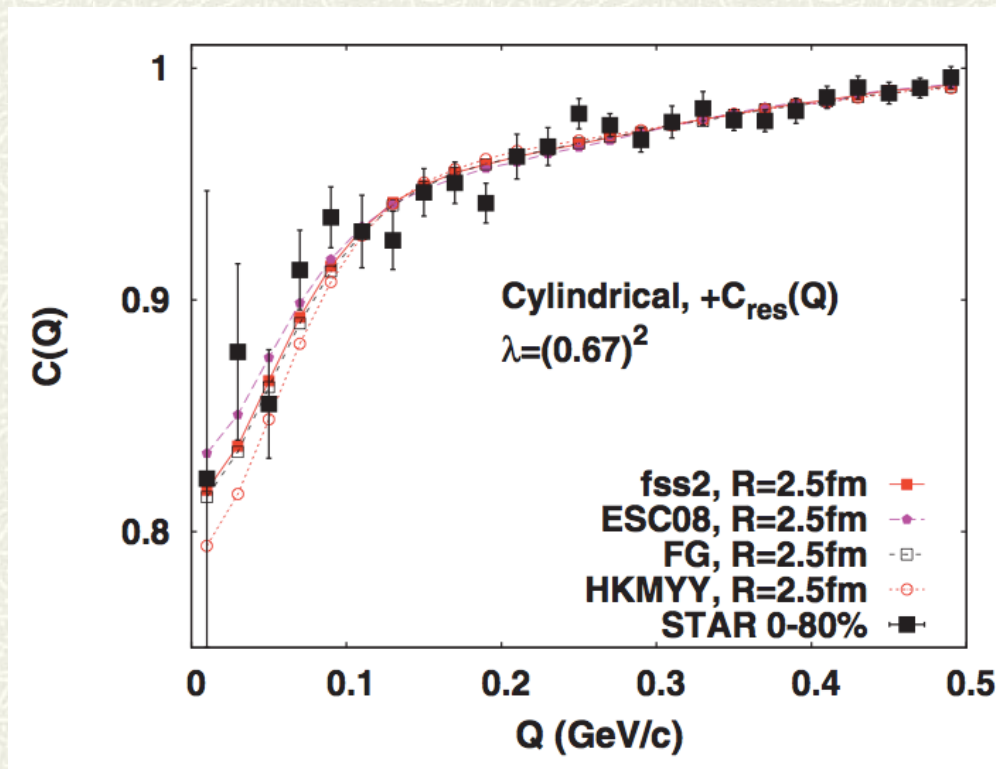
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Takashi Inoue (HAL QCD Collaboration)

○



$\Lambda\Lambda$ correlation in Heavy Ion Collisions

- STAR collaboration, PRL 114, 022301 (2015)
 $\Lambda\Lambda$ correlation function in Au+Au collisions at $\sqrt{s_{NN}}=200$ GeV
- K. Morita, T. Furumoto, A. Ohnishi, PRC 91, 024916 (2015)
 $\Lambda\Lambda$ interaction from relativistic heavy-ion collisions

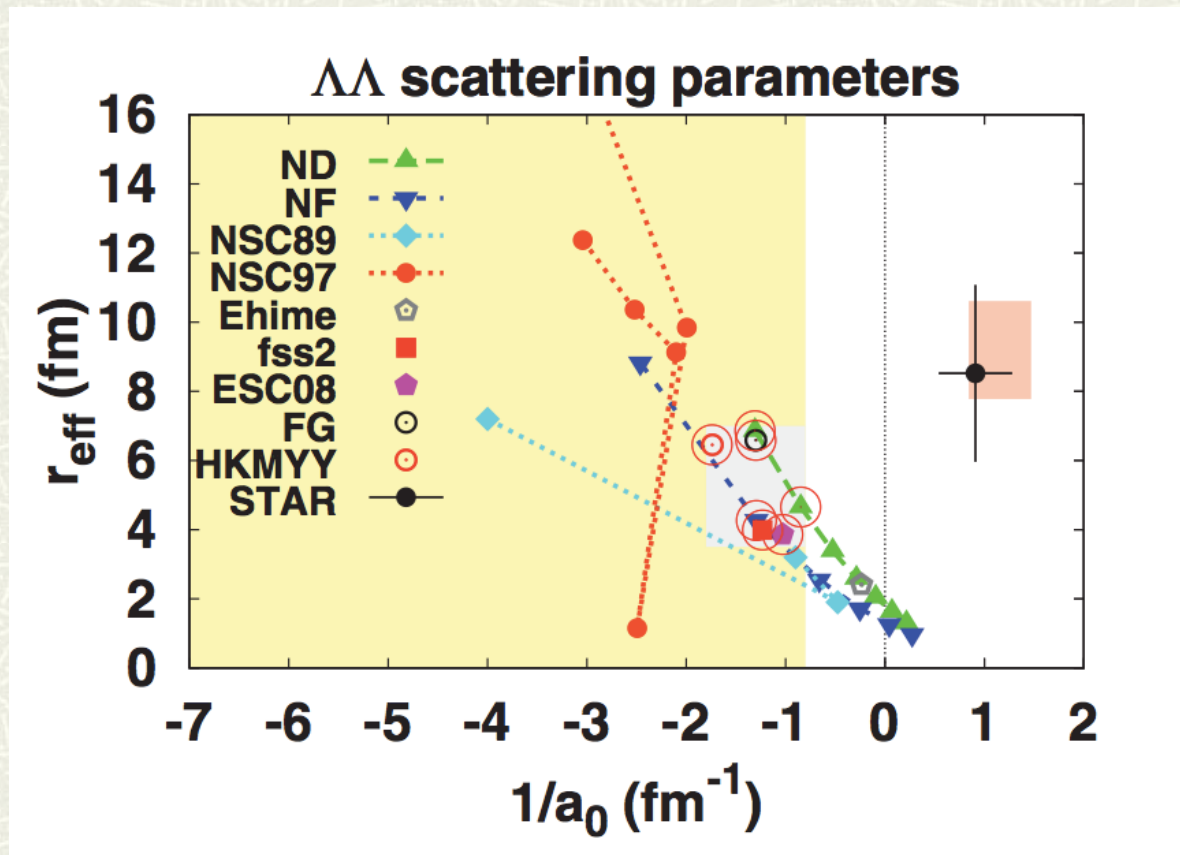


$\Lambda\Lambda$ correlation in Heavy Ion Collisions

- # K. Morita, T. Furumoto, A. Ohnishi, PRC 91, 024916 (2015)
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$D_\Delta (\Delta\Delta)_{I=0}$ dibaryon

$D_\Delta (\Delta\Delta)_{I=0}$ dibaryon

‡ $S=3, I=0$ (Δ^2) bound state

→ relatively narrow $NN\pi\pi$ ($I=0$) resonance

NUCLEAR FORCE IN A QUARK MODEL

M. OKA and K. YAZAKI

*Department of Physics, Faculty of Science, University of Tokyo,
Bunkyo-ku, Tokyo 113, Japan*

Received 24 July 1979

Phys. Lett. 90B (1980) 41

The problem of the nuclear force in a nonrelativistic quark model is studied by the resonating group method which has been extensively used in treating the interaction between composite particles. The calculated phase shifts for the 3S_1 and 1S_0 states of two nucleons indicate the presence of a strong repulsive force at short distance, while an attractive force is predicted for the 7S_3 ($(S, T) = (3, 0)$) state of two Δ 's. These features are due to an interplay between the Pauli principle and the spin-spin interaction between quarks.

Classification of two baryon systems without strangeness.

The spin-flavor $SU(6)$ is reduced to the spin-isospin $SU(4)$.

$S(I)$ denotes the total spin (isospin) of the system.

taken from D. Sc Thesis by M.O. (1980)

L	$SU(4)$	$BB' (S, I)$
even	$\{33\}$	$\Delta\Delta (3,0), \Delta\Delta (0,3)$ <i>avored by $SU(6)$</i>
	$\{51\}$	$\Delta\Delta (3,2), \Delta\Delta (2,3), N\Delta (2,2), N\Delta (1,1)$
	$\{33\} + \{51\}$	$\Delta\Delta + N\Delta (2,1), \Delta\Delta + N\Delta (1,2)$ $NN + \Delta\Delta (1,0), NN + \Delta\Delta (0,1)$
odd	$\{6\}$	$\Delta\Delta (3,3)$
	$\{42\}$	$\Delta\Delta (3,1), \Delta\Delta (1,3), \Delta\Delta (2,0), \Delta\Delta (0,2)$ $N\Delta (2,1), N\Delta (1,2)$
	$\{6\} + \{42\}$	$N\Delta + \Delta\Delta (2,2), NN + \Delta\Delta (0,0)$
	$\{6\} + \{42\}^2$	$NN + N\Delta + \Delta\Delta (1,1)$

$$\Gamma_{\text{CM}} \equiv - \sum_{i < j} (\lambda_i^a \lambda_j^a) (\sigma_i^k \sigma_j^k) = 8n - 2C_6 + \frac{4}{3}S(S+1)$$

$$C_6 \equiv C_2[SU(6)_{\text{cs}}] = \sum_i f_i(f_i - 2i + 7) - \frac{n^2}{6}$$

CMI strength: $V_0 = 300/16 \sim 18(\text{MeV})$ from $\Gamma_{\text{CM}}(\Delta) = +8$
 $\Gamma_{\text{CM}}(N) = -8$

$SU(6)_{\text{cs}}$ representation	$4C_6$	$SU(3)_{\text{f}}$ representation	
490	144	$\underline{1}$	$H = \Lambda\Lambda(I = S = 0) \quad V = V_0 \times (-8)$
896	120	$\underline{8}$	
280	96	$\underline{10}$	
175	96	$\underline{10}^*$	$\Delta\Delta(I = 0, S = 3) \quad V = V_0 \times 0$
189	80	$\underline{27}$	<i>avored</i>
35	48	$\underline{35}$	
1	0	$\underline{28}$	$\Delta\Delta(I = 3, S = 0) \quad V = V_0 \times 32$
			<i>less favored</i>

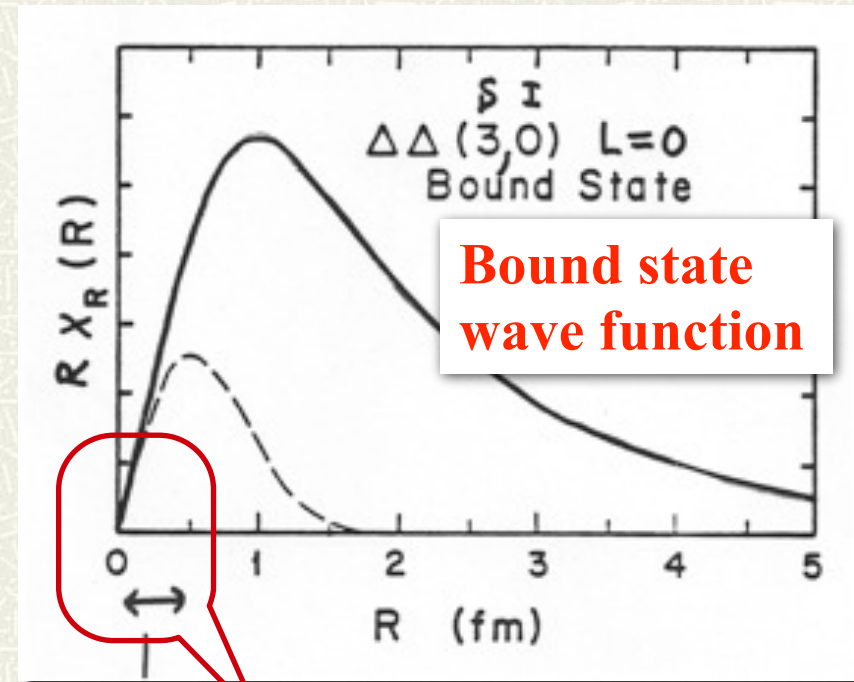
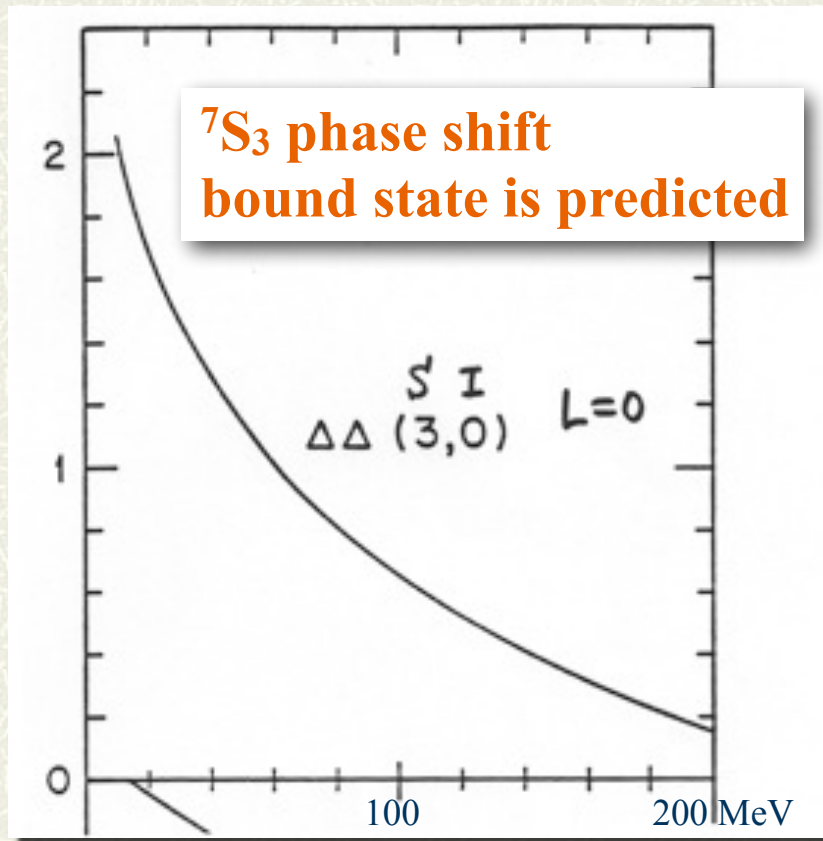
Perhaps a Stable Dihyperon*

R.L. Jaffe, PRL 38 (1977) 195

$D_\Delta (\Delta\Delta)_{I=0}$ dibaryon

Quark Cluster Model: $S=3, I=0$ (Δ^2) bound state

MO, K. Yazaki, Phys. Lett. 90B (1980) 41

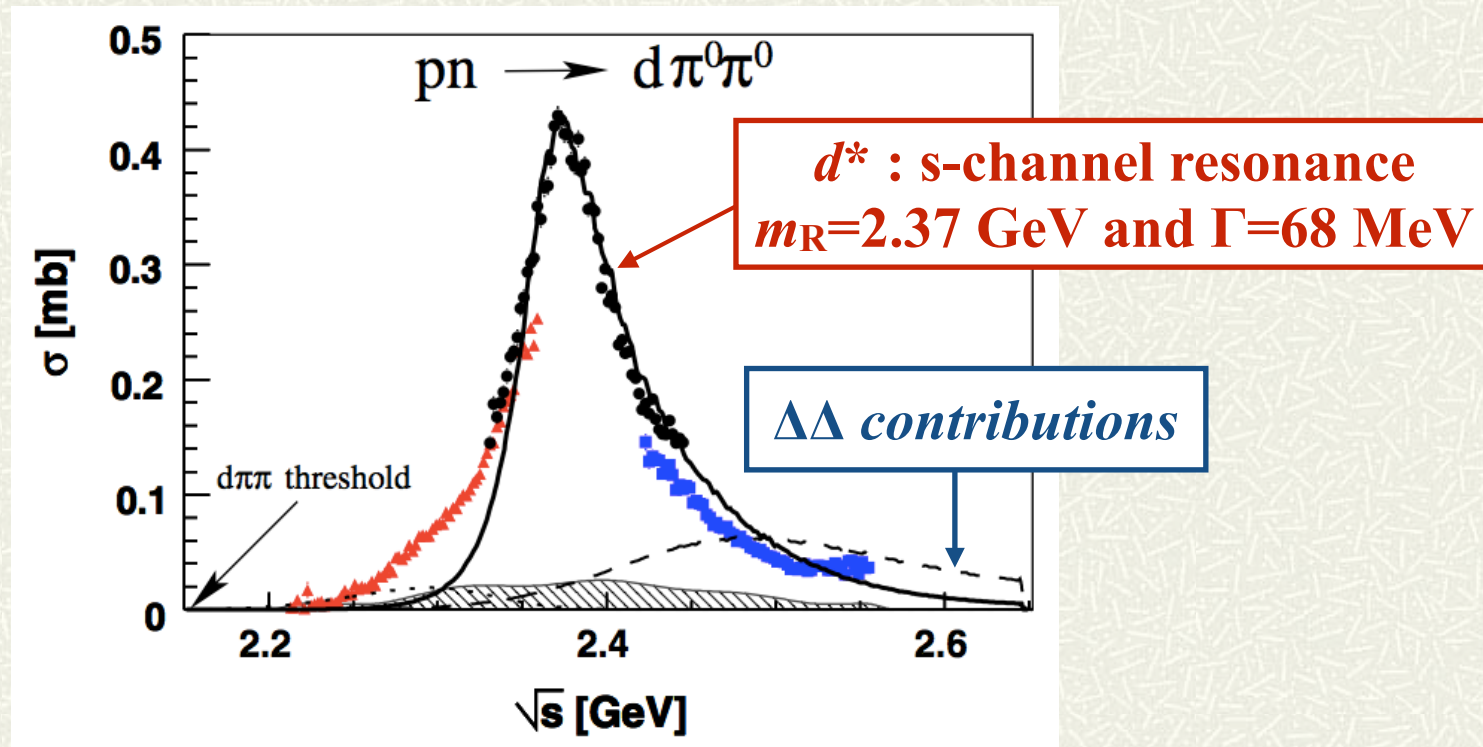


No repulsive core

d^* resonance

WASA@COSY, PRL 106, 242302 (2011)

$p + n(d) \rightarrow d + \pi^0 + \pi^0$ (+ $p_{\text{spectator}}$) at $T_p=1.0, 1.2, 1.4$ GeV

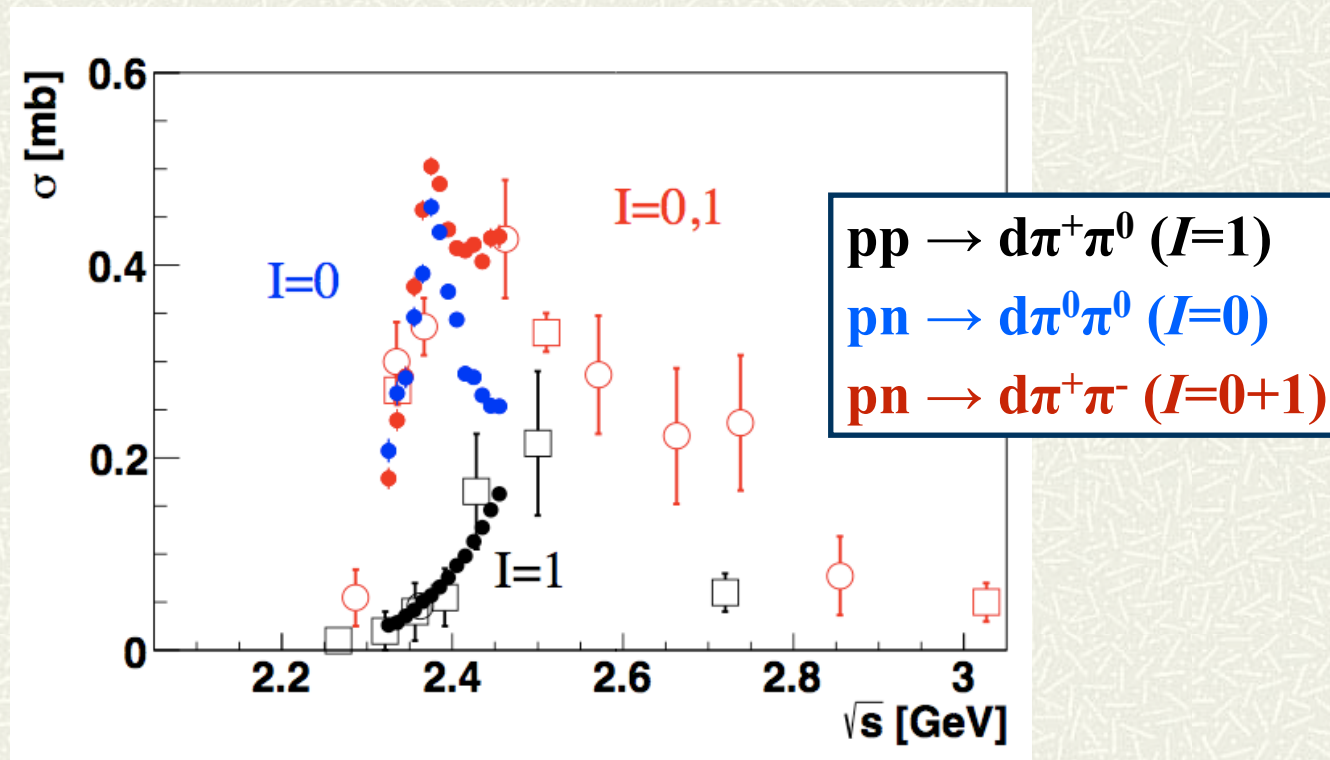


A di-baryon resonance, d^* ($I=0, J^\pi=3^+$) (in pn and $\Delta\Delta$) is confirmed.

d^* resonance

WASA@COSY, PLB 721 (2013) 229

Isospin decomposition of the basic double-pionic fusion in the region of the ABC effect



The ($I=1$) production is consistent with the $\Delta\Delta$ production.

d^* resonance

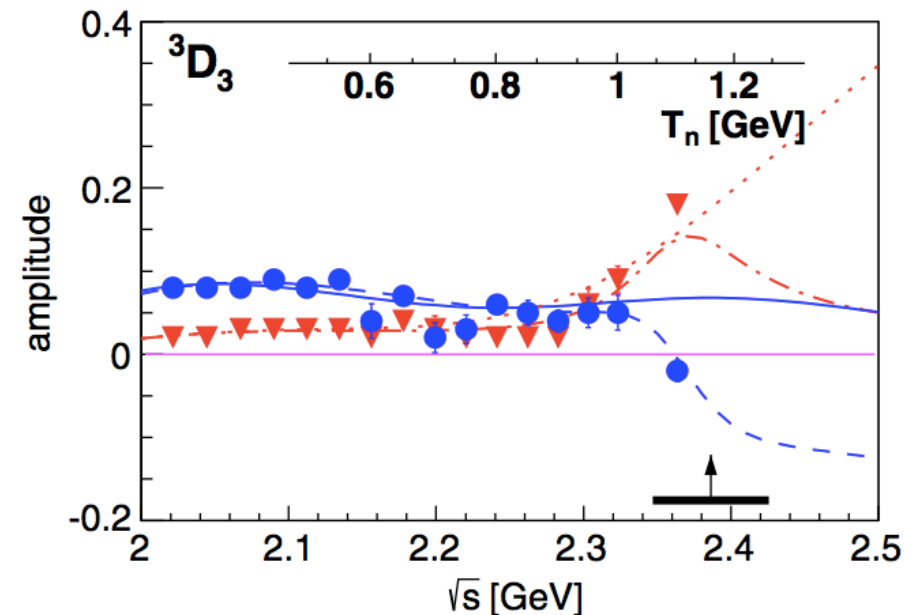
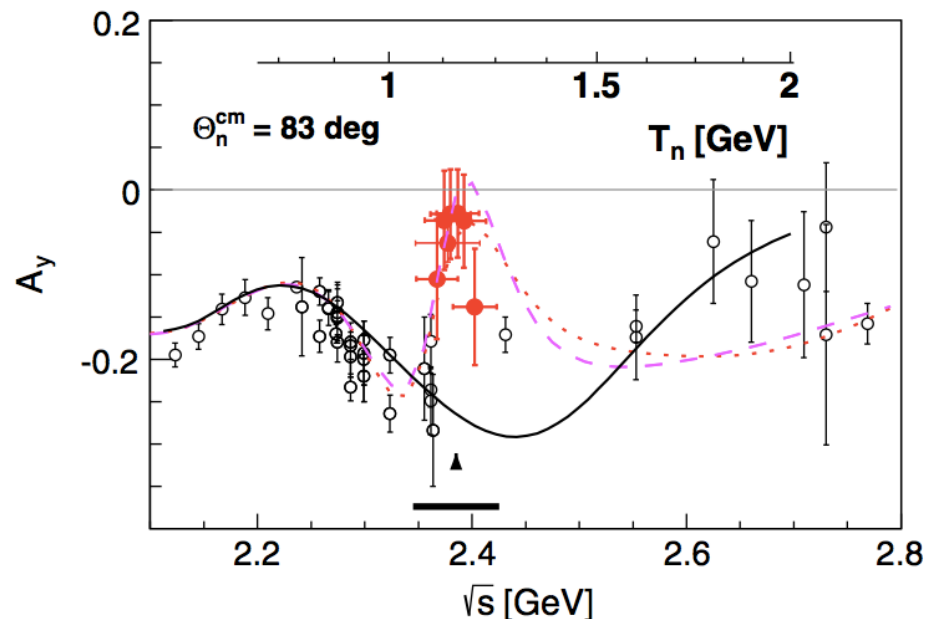
■ **WASA@COSY+SAID, PRL 112, 202301 (2014)**

Evidence for a new resonance from polarized n-p scattering

$d(\uparrow) + p \rightarrow np + p_{\text{spectator}}$

np analyzing power, $A_y(\theta)$, at $T_n=1.108\text{-}1.197$ GeV

A phase shift analysis of 3D_3 (3^+) amplitudes shows a narrow resonance at $M=2380$ MeV and $\Gamma\sim 70$ MeV.



d^* resonance

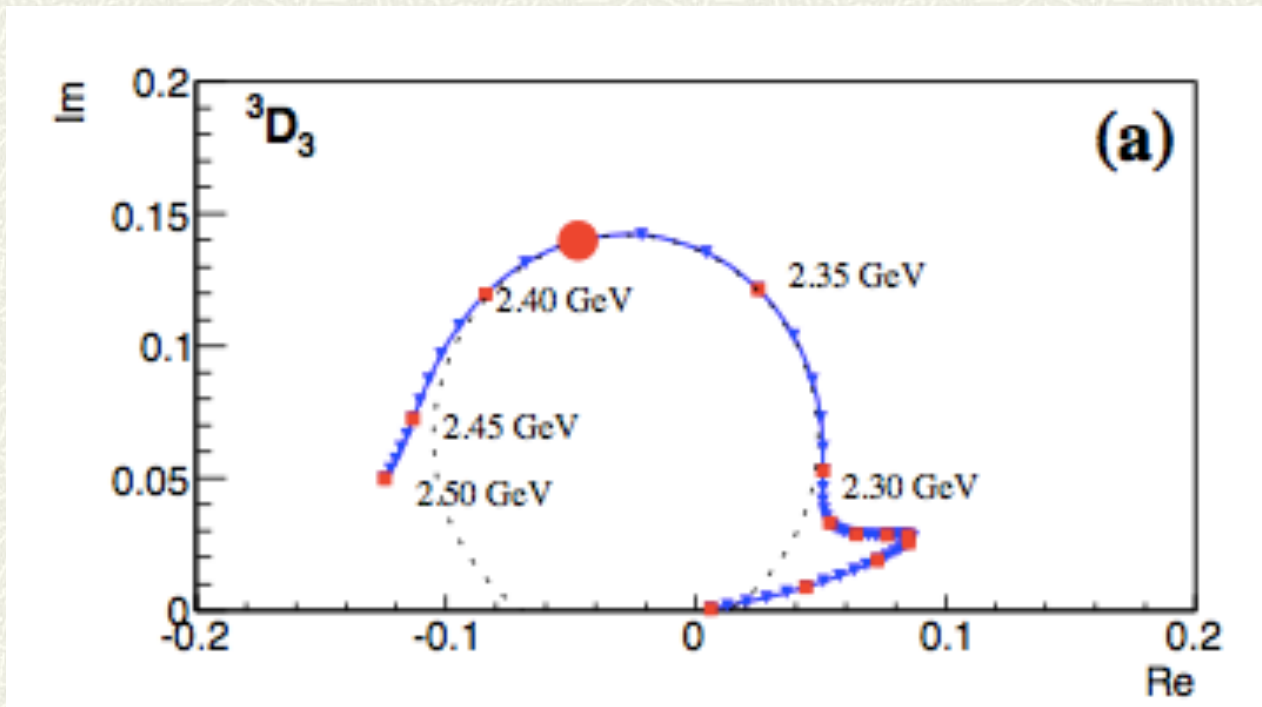
WASA@COSY+SAID, PRL 112, 202301 (2014)

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Conclusion

- # “Dibaryon” is a long-standing but still exciting subject. Its existence should be correlated to the short-range baryonic interactions.
- # LQCD has confirmed the Pauli effect and the CMI for the short-range baryon-baryon interactions.
- # The quark model symmetries, $SU(6)_{sf}$ for the Pauli effect and $SU(3)_f$ for the CMI, give guideline for possible compact dibaryons.
- # H ($F=1$) is the most-likely dibaryon.
- # $D_\Delta = (\Delta\Delta)$ ($I=0, S=3$) is another favorable state.
The d^* resonance at WASA-COSY is a strong candidate of a “compact” dibaryon.