

# $\Lambda$ - $\Lambda$ and p- $\Xi^-$ interaction measurements with femtoscopy in ALICE

VALENTINA MANTOVANI SARTI FOR THE ALICE COLLABORATION

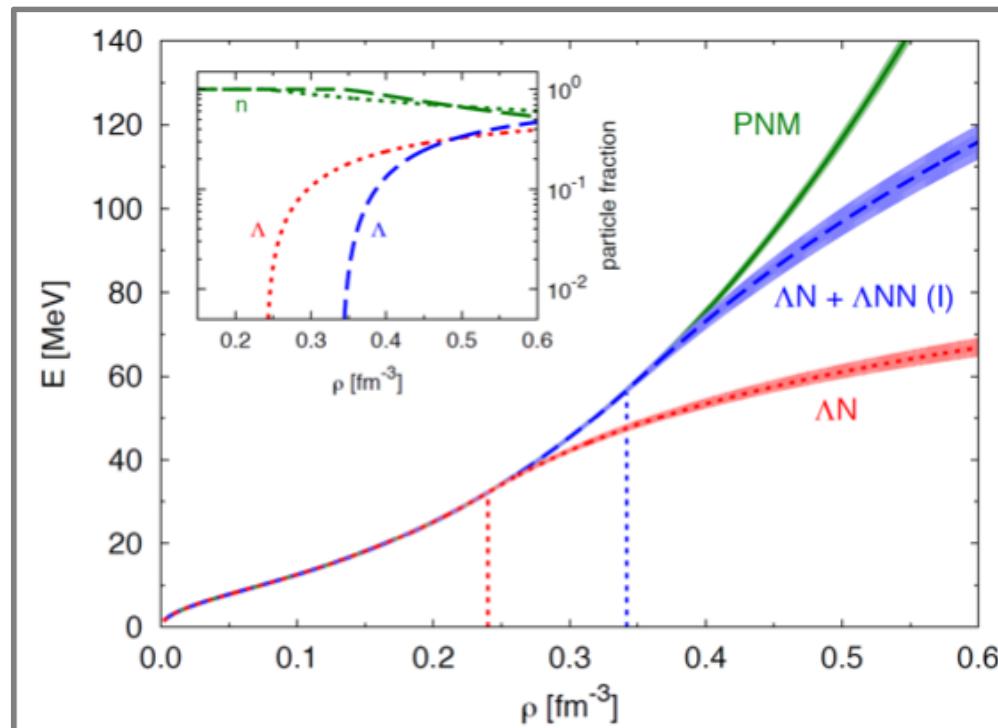
PHYSIK-DEPARTMENT - TECHNISCHE UNIVERSITÄT MÜNCHEN

QNP SATELLITE WORKSHOP – TOKAI CAMPUS 11 NOVEMBER 2018

1. Hyperon-Nucleon and Hyperon-Hyperon interactions in Neutron Stars
2. Femtoscopy at ALICE experiment
3. Results pp 13 TeV and p-Pb 5.02 TeV
  - $\Lambda$ - $\Lambda$  correlation function and the existence of H di-baryon
  - First observation of the p- $\Xi^-$  strong attractive potential
4. Summary and Outlooks

# Hyperon Puzzle in Neutron Stars

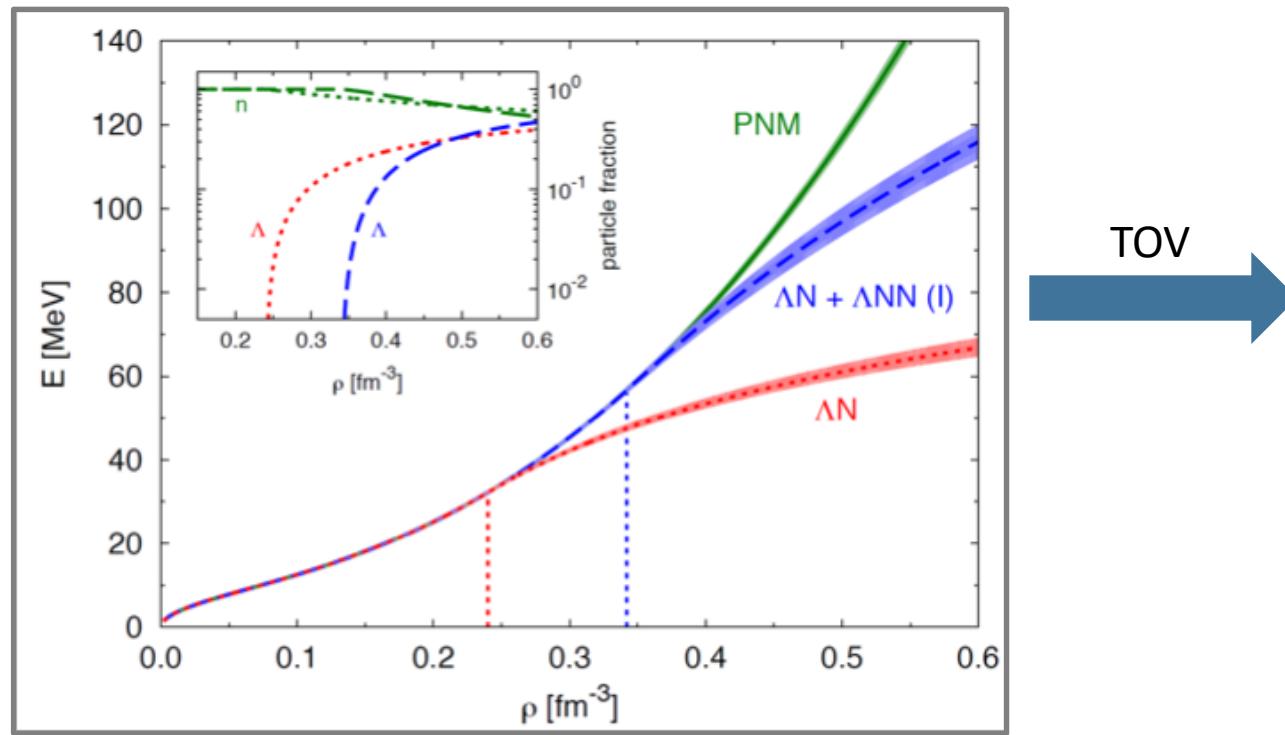
- Hyperons production becomes **energetically favorable at moderate-large densities in dense neutron-rich matter**
- Threshold **strongly depends on the Y-N interaction**
- The appearance of Hyperons softens the EoS



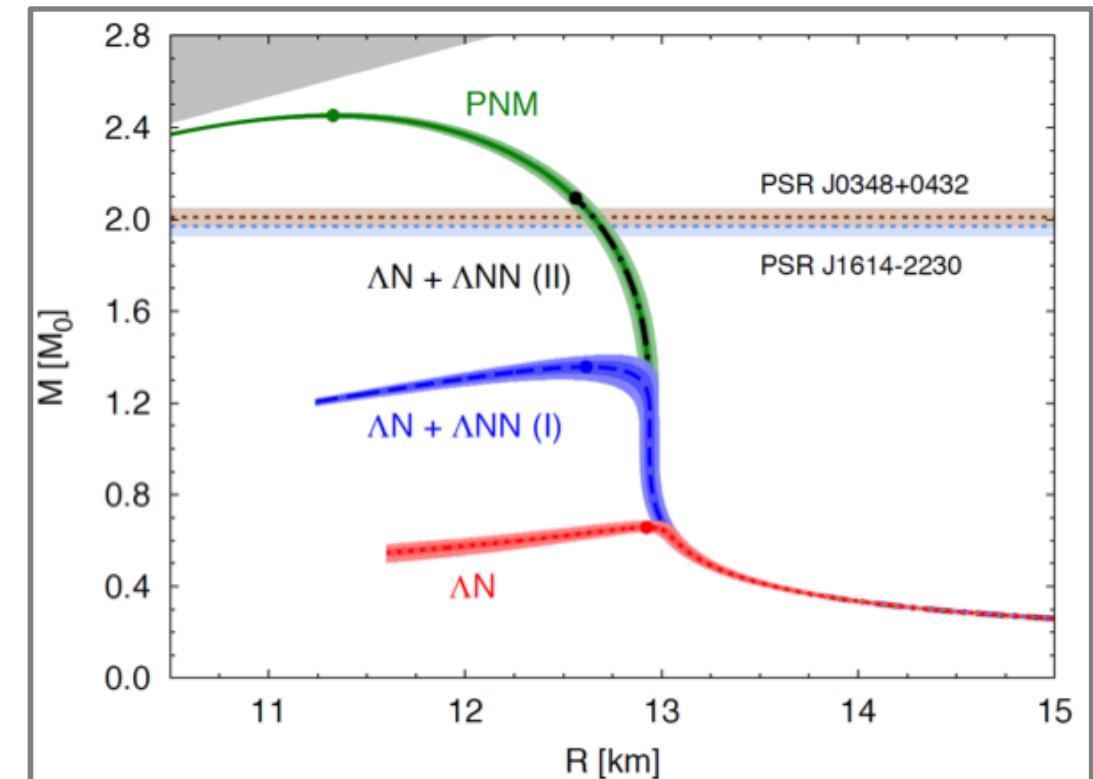
D. Lonardoni, A. Lovato, S. Gandolfi, F. Pederiva Phys. Rev. Lett. 114, 092301 (2015)

# Hyperon Puzzle in Neutron Stars

- Hyperons production becomes **energetically favorable at moderate-large densities in dense neutron-rich matter**
- Threshold **strongly depends on the Y-N interaction**
- The appearance of Hyperons softens the EoS
- **Maximum NS masses get smaller**  $\Rightarrow$  **HYPERON PUZZLE**



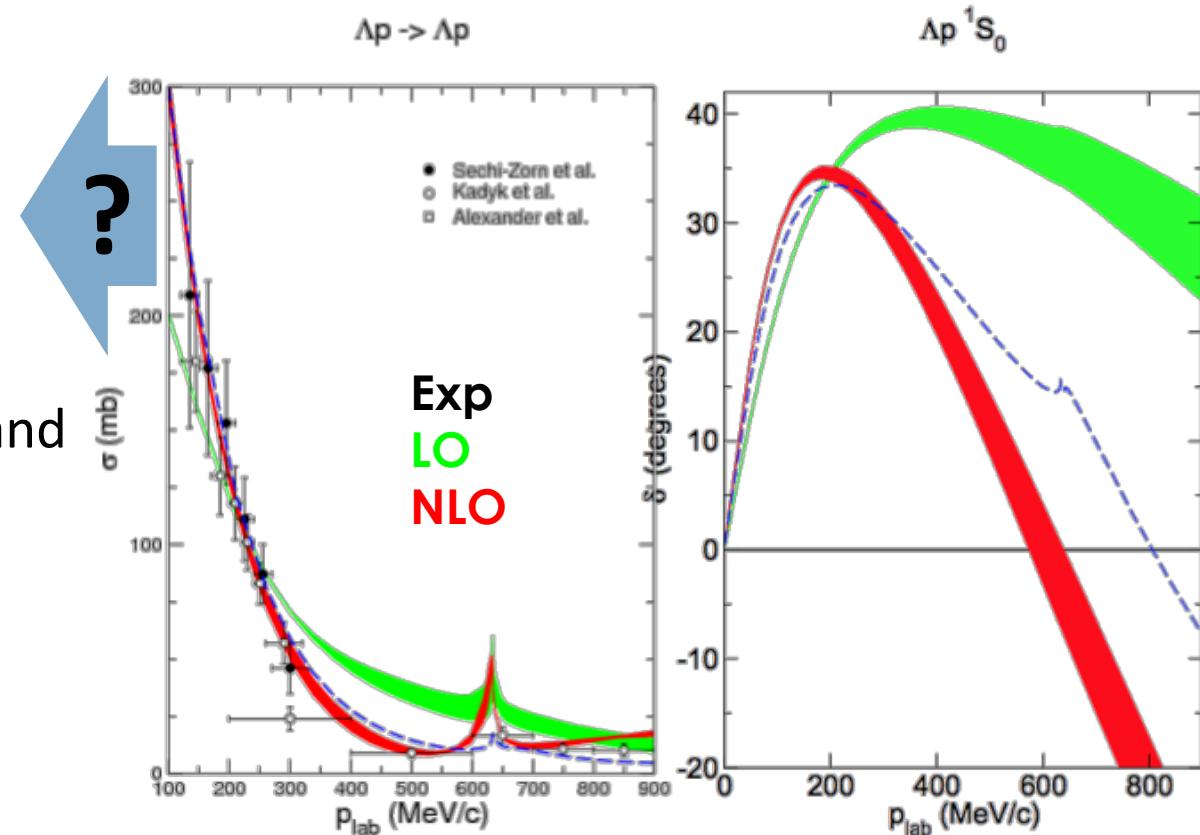
TOV  $\rightarrow$



D. Lonardoni, A. Lovato, S. Gandolfi, F. Pederiva Phys. Rev. Lett. 114, 092301 (2015)

# YN and YY interactions: the traditional ways

- Scattering experiments / Hypernuclei data
- **SCATTERING DATA:**
- Constraint only on  $p\Lambda/p\Sigma$  interaction
  - Data from scattering experiments from 1968 and 1971 in bubble chambers
  - Production threshold:  $p_{LAB} \gtrsim 100$  MeV



LO: H. Polinder, J.H., U. Meiñner, NPA 779 (2006) 244  
 NLO: J. Haidenbauer., N. Kaiser, et al., NPA 915 (2013) 24

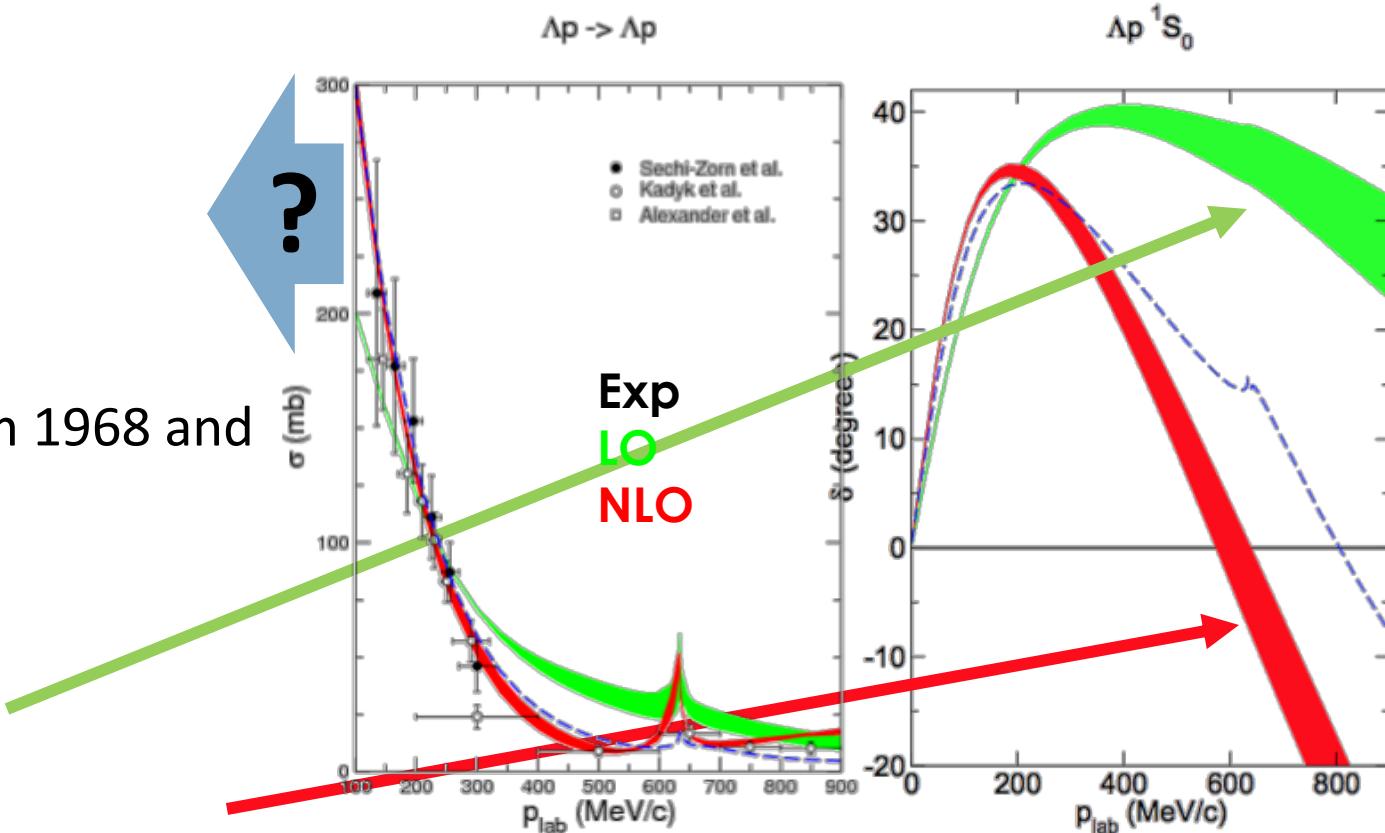
# YN and YY interactions: scattering experiments

- Scattering experiments / Hypernuclei data

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- Constraint only on  $p\Lambda/p\Sigma$  interaction

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- Production threshold:  $p_{LAB} \gtrsim 100$  MeV
- Cannot probe low momentum region



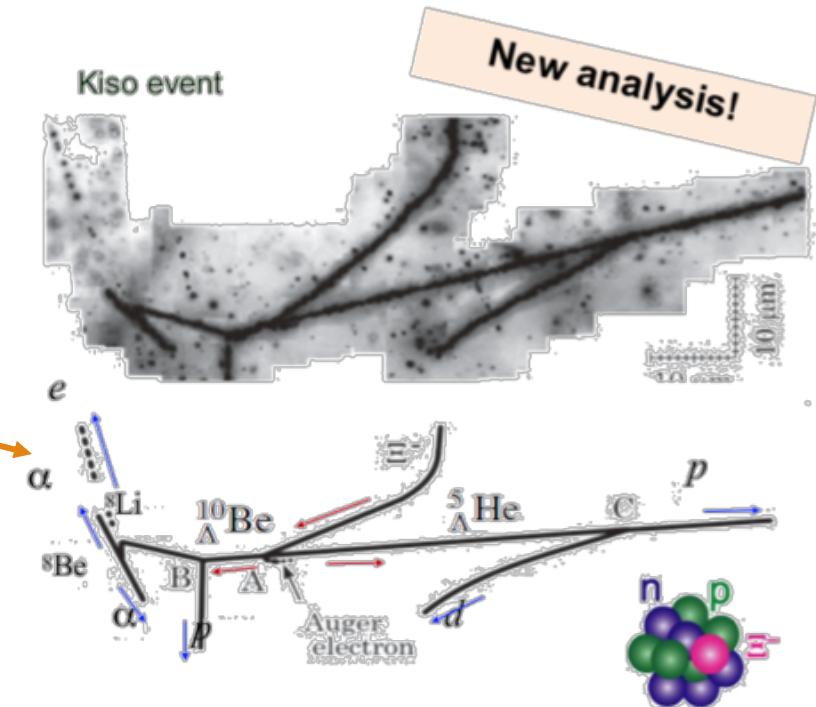
**LO: ALWAYS ATTRACTIVE**

**NLO: BECOMES REPULSIVE  $\Rightarrow$  REPULSIVE CORE**

LO: H. Polinder, J.H., U. Mei $\beta$ nner, NPA 779 (2006) 244  
 NLO: J. Haidenbauer., N. Kaiser, et al., NPA 915 (2013) 24

# YN and YY interactions: hypernuclei

Courtesy H. Tamura, Bormio Winter Meeting 2018



The first clear  $\Xi$  hypernucleus

$$B_{\Xi^-} = 4.38 \pm 0.25 \text{ MeV}, \\ - 1.11 \pm 0.25 \text{ MeV}$$

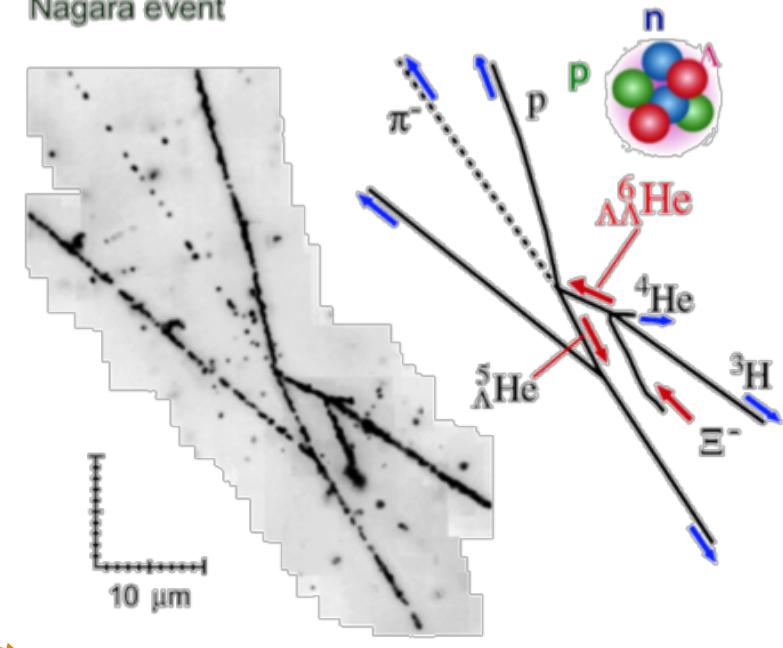
K. Nakazawa et al. PTEP 2015, 033D02

- **HYPERNUCLEI DATA:**
  - $\Lambda$  hypernuclei data:  $B_{\Lambda\text{nucleus}} = 30 \text{ MeV}$
  - $\Xi$  hypernuclei data show **shallow attractive** interaction [Kiso Evt]
  - $\Lambda\Lambda$  hypernuclei exist as well showing **weakly attraction** [Nagara Evt]
  - $\Sigma$  hypernuclei: nothing known so far!

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Nagara event



$$\Delta B_{\Lambda\Lambda} = 0.67 \pm 0.17 \text{ MeV}$$

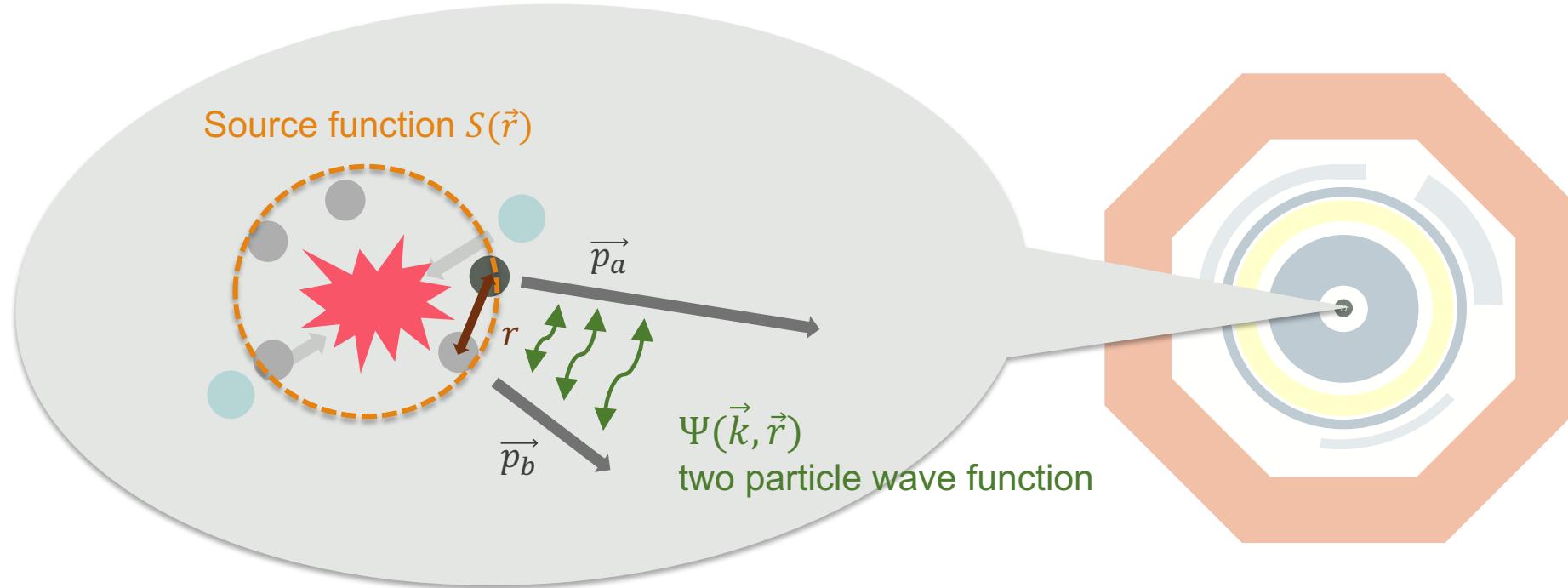
*H. Takahashi et al., PRL 87 (2001) 212502*

**$\Lambda\Lambda$  is weakly attractive**

# Take home message

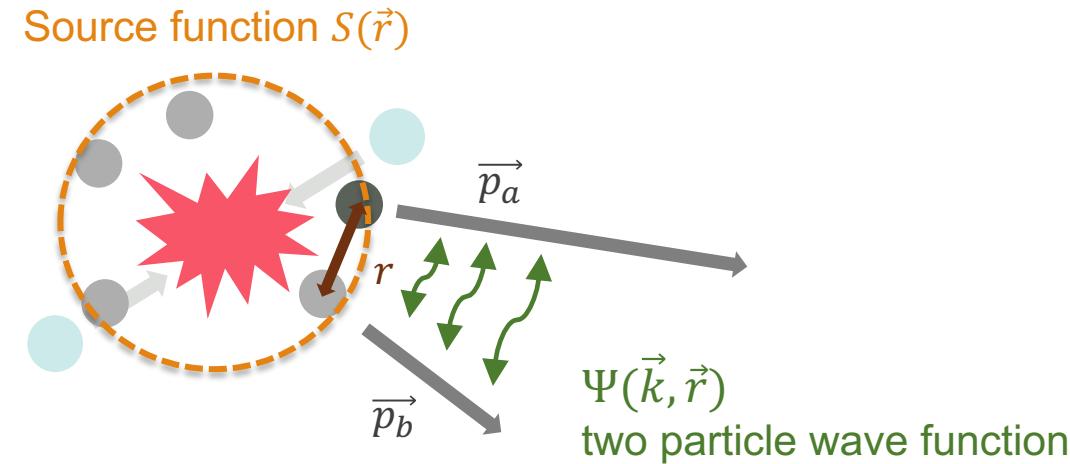
- Hyperon-nucleon and hyperon-hyperon interactions are fundamental ingredients for the EoS of NSs but unfortunately already in vacuum they are still not well fully understood
- Scattering experiments cannot be applied to all hyperon species and cannot access the low-momentum range
- Hypernuclei experiments are based on single events and their interpretation is model dependent
- How can we do the next step?  
**Femtoscopy**

# Basics of Femtoscopy



- Look for particle pairs
- Measure the relative momentum  $k$  distribution in your experimental set up
- Search for correlations

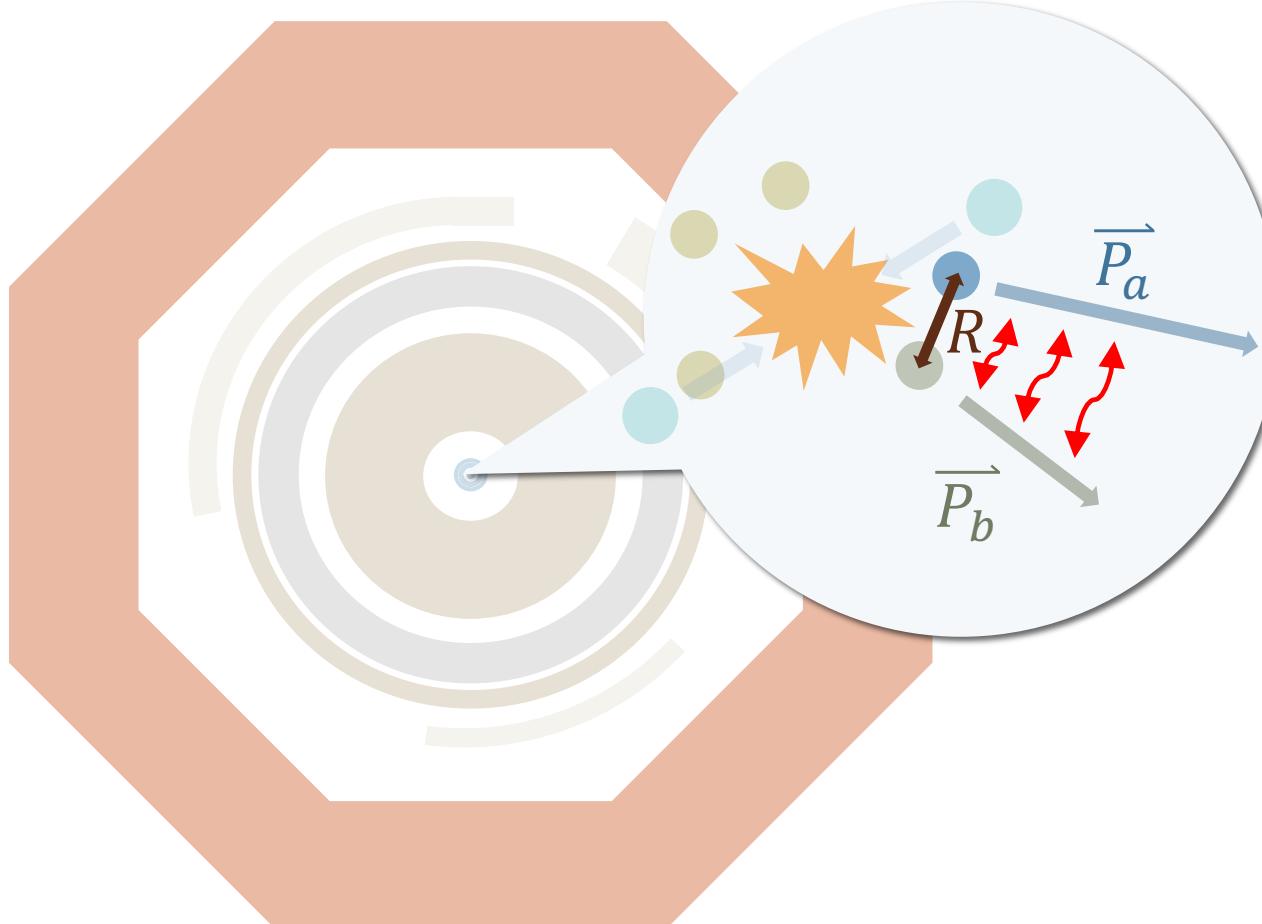
# Key observable: the correlation function



<u>Statistical definition</u>	<u>Experimental definition</u>	<u>Theoretical definition</u>
$C(k) = \frac{\mathcal{P}(\vec{p}_a, \vec{p}_b)}{\mathcal{P}(\vec{p}_a)\mathcal{P}(\vec{p}_b)}$	$= \mathcal{N} \frac{N_{\text{Same}}(k)}{N_{\text{Mixed}}(k)}$	$= \int S(\vec{r})  \Psi(\vec{k}, \vec{r}) ^2 d^3 \vec{r} \xrightarrow{k \rightarrow \infty} 1$
Single-particle momenta	Relative distance / reduced momentum in the rest frame of the pair	

Lisa, Pratt, Wiedemann,Solz,  
Ann.Rev.Nucl.Part.Sci. 55 (2005) 357-402

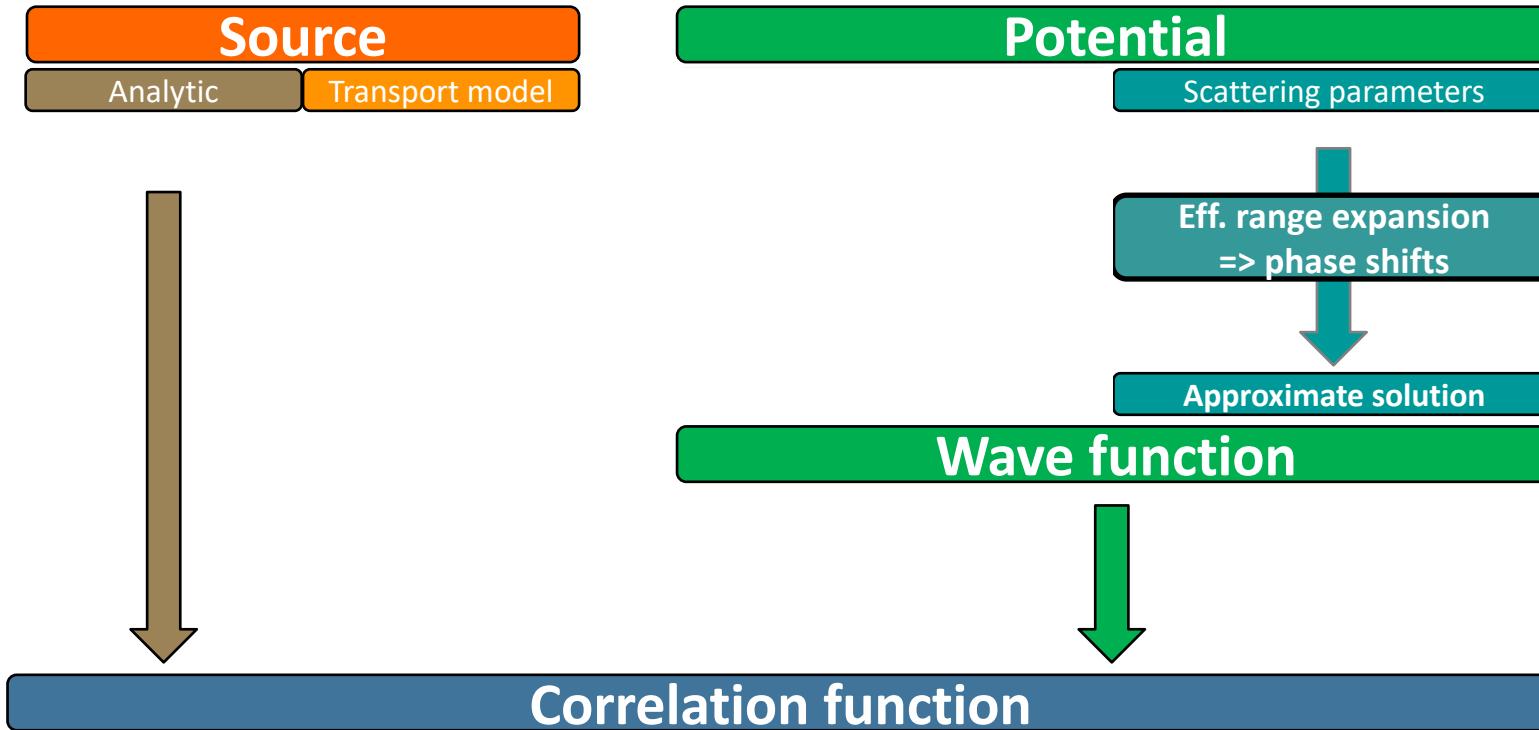
# Femtoscopy at A Large Ion Collider Experiment



$$C(k) = \int S(\vec{r}) |\Psi(\vec{k}, \vec{r})|^2 d^3\vec{r} \xrightarrow{k \rightarrow \infty} 1$$

- We measure **p-p, p- $\Lambda$ ,  $\Lambda$ - $\Lambda$ , p- $\Xi$**
- Proton identification with TPC and TOF
- Reconstruction of hyperons
  - $\Lambda \rightarrow p\pi^-$  (BR  $\sim 64\%$ )
  - $\Xi^- \rightarrow \Lambda\pi^-$  (BR  $\sim 100\%$ )
- Datasets:
 

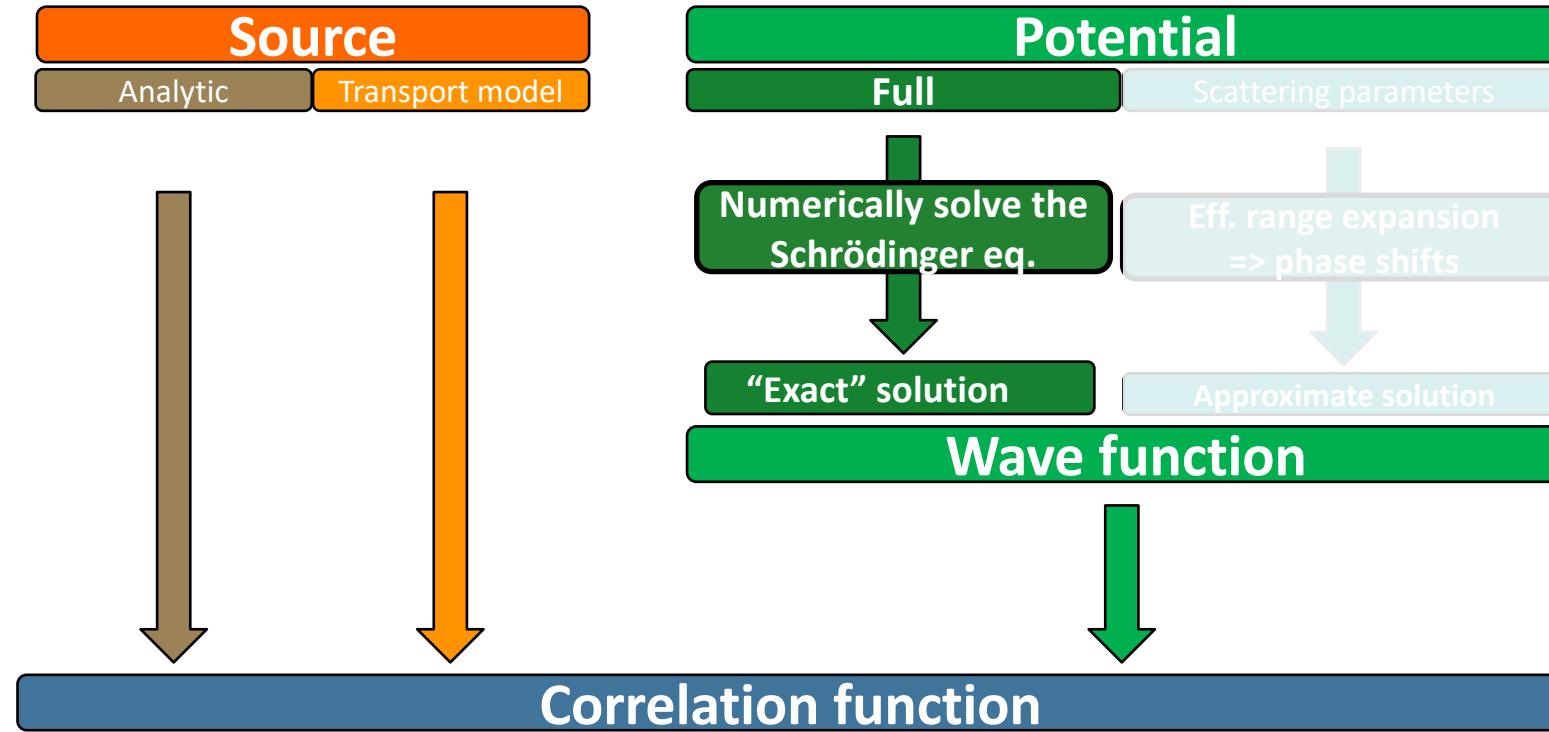
• <b>pp 7 TeV:</b>	<b><math>3.4 \cdot 10^8</math> Events</b>
• <b>pp 13 TeV:</b>	<b><math>10 \cdot 10^8</math> Events</b>
• <b>p-Pb 5.02 TeV:</b>	<b><math>6.0 \cdot 10^8</math> Events</b>
- Assumption of a **common Gaussian source**  
 ⇒ fixing the radius from **p-p**  
 ⇒ **p- $\Lambda$ , p- $\Xi$  and  $\Lambda$ - $\Lambda$**  Correlation Function to study the interaction



$$C(k) = 1 + \sum_S \rho_S \left[ \frac{1}{2} \left| \frac{f^S(k)}{R_G^{\Lambda p}} \right|^2 \left( 1 - \frac{d_0^S}{2\sqrt{\pi} R_G^{\Lambda p}} \right) + 2 \frac{\mathcal{R} f^S(k)}{\sqrt{\pi} R_G^{\Lambda p}} F_1(Q R_G^{\Lambda p}) - \frac{\mathcal{I} f^S(k)}{R_G^{\Lambda p}} F_2(Q R_G^{\Lambda p}) \right]$$

Depends on scattering parameters, might locally break down for small sources

R. Lednicky and V. L. Lyuboshits, Sov. J. Nucl. Phys. 35, 770 (1982), [Yad. Fiz. 35, 1316 (1981)].

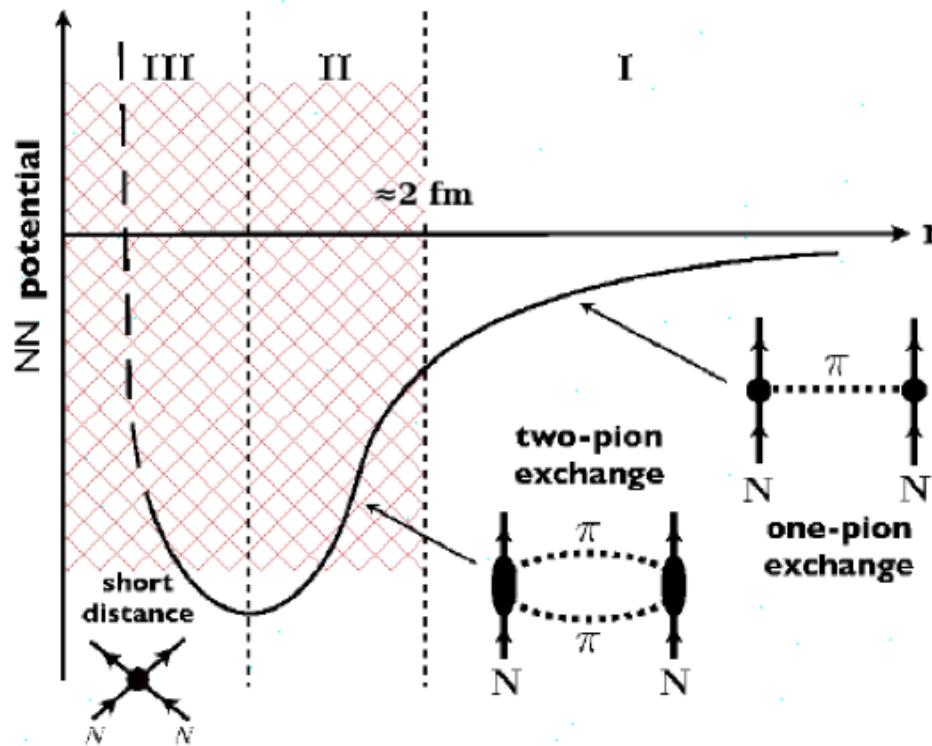


$$C(k) = \int S(\vec{r}, k) |\psi(\vec{r}, k)|^2 d\vec{r} \xrightarrow{k \rightarrow \infty} 1$$

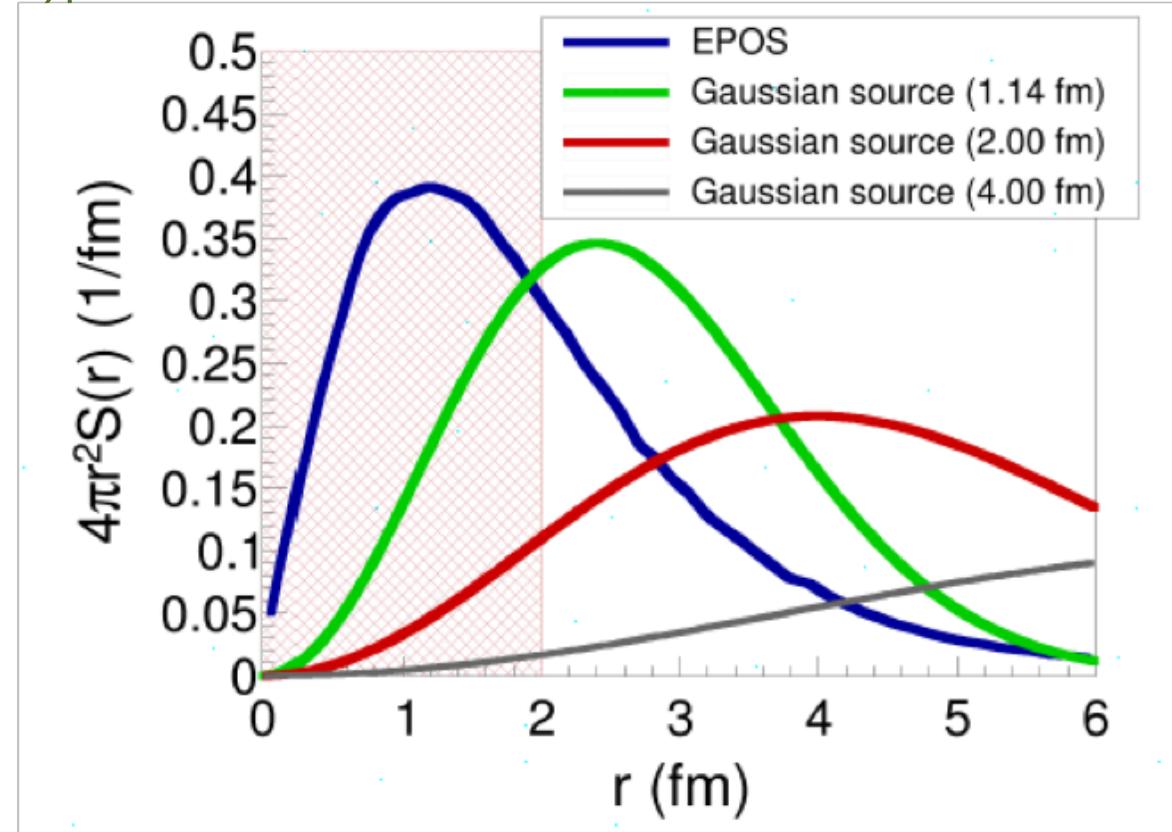
(D.L.Mihaylov et al. Eur.Phys.J. C78 (2018) no.5,394)

# Length scales involved in Femtoscopy

$$C(k) = \int S(\vec{r}) |\Psi(\vec{k}, \vec{r})|^2 d^3\vec{r} \xrightarrow{k \rightarrow \infty} 1$$



Based on Holt, Kaiser and Weise,  
Prog.Part.Nucl.Phys. 73 (2013) 35-83



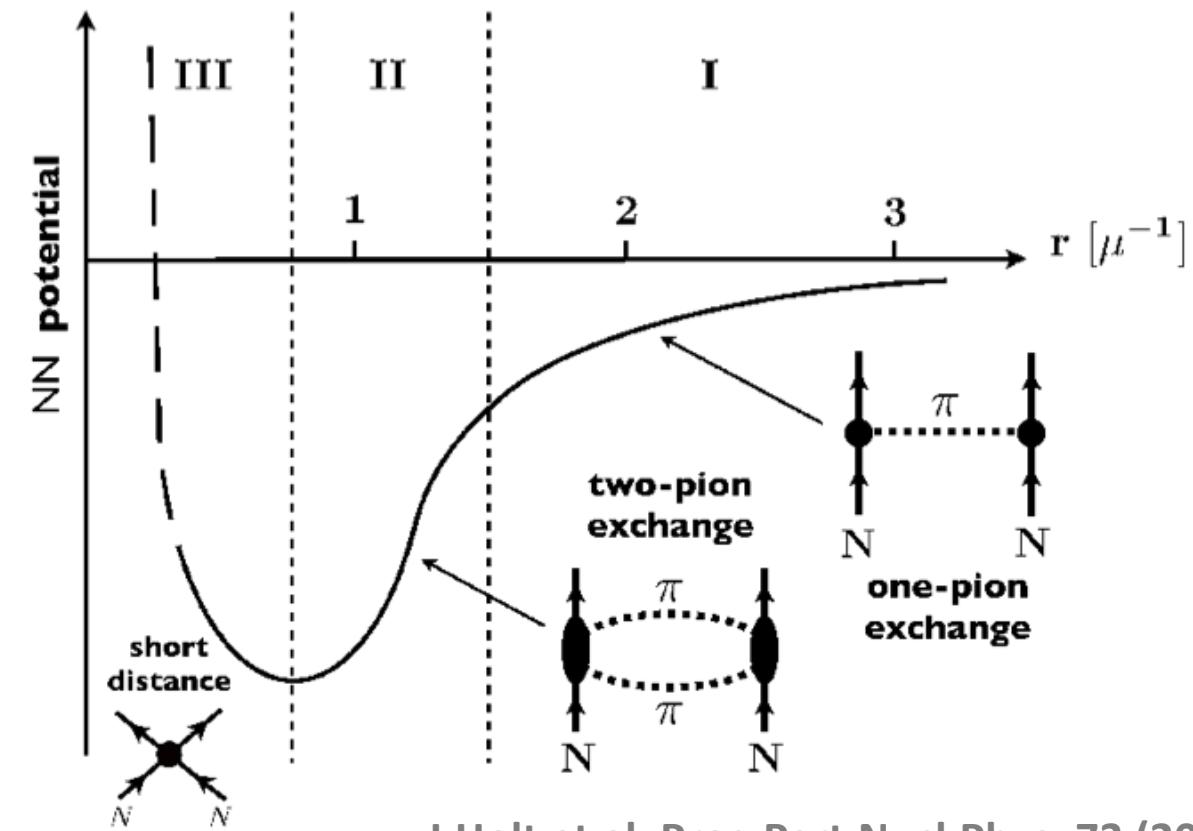
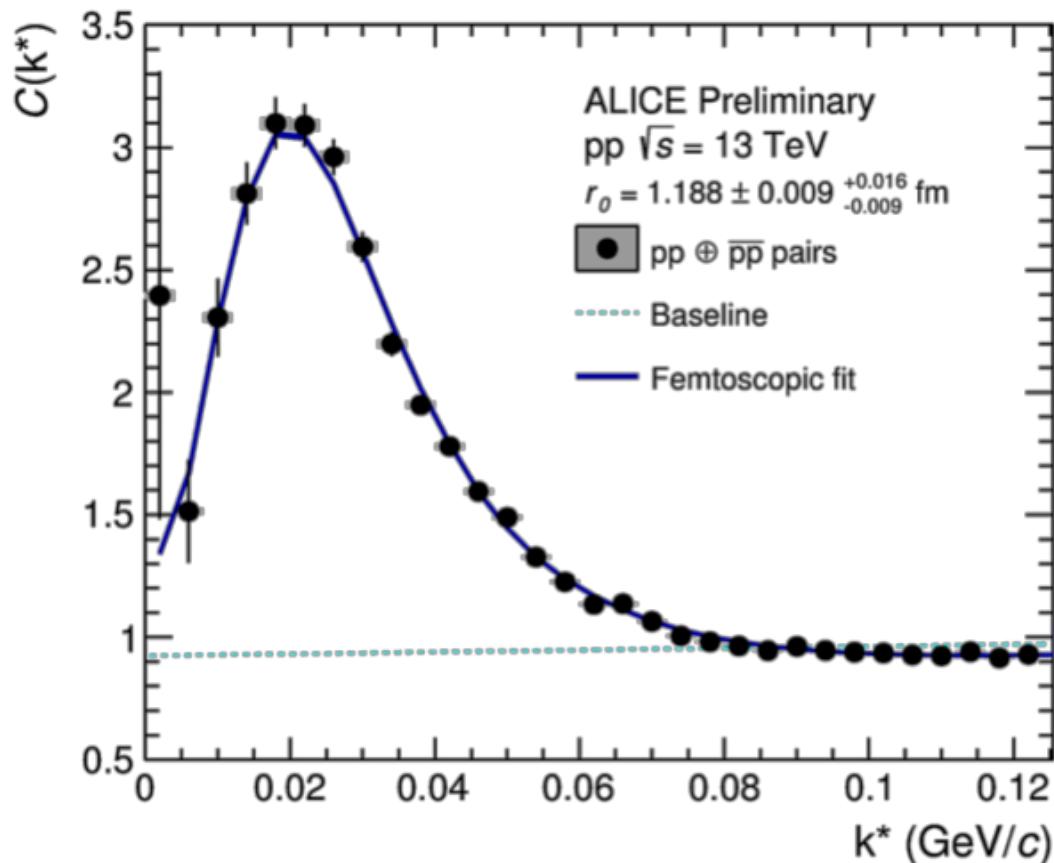
- A-A collisions:  $R_{\text{Gauß}} \sim 4$  fm
- p-A collisions:  $R_{\text{Gauß}} \sim 2$  fm
- p-p collisions:  $R_{\text{Gauß}} \sim 1$  fm

# Correlation function and Interactions

$> 1$  attraction

$C(k^*) \rightarrow = 1$  no interaction

$< 1$  repulsion



J.Holt et al, Prog.Part.Nucl.Phys. 73 (2013)

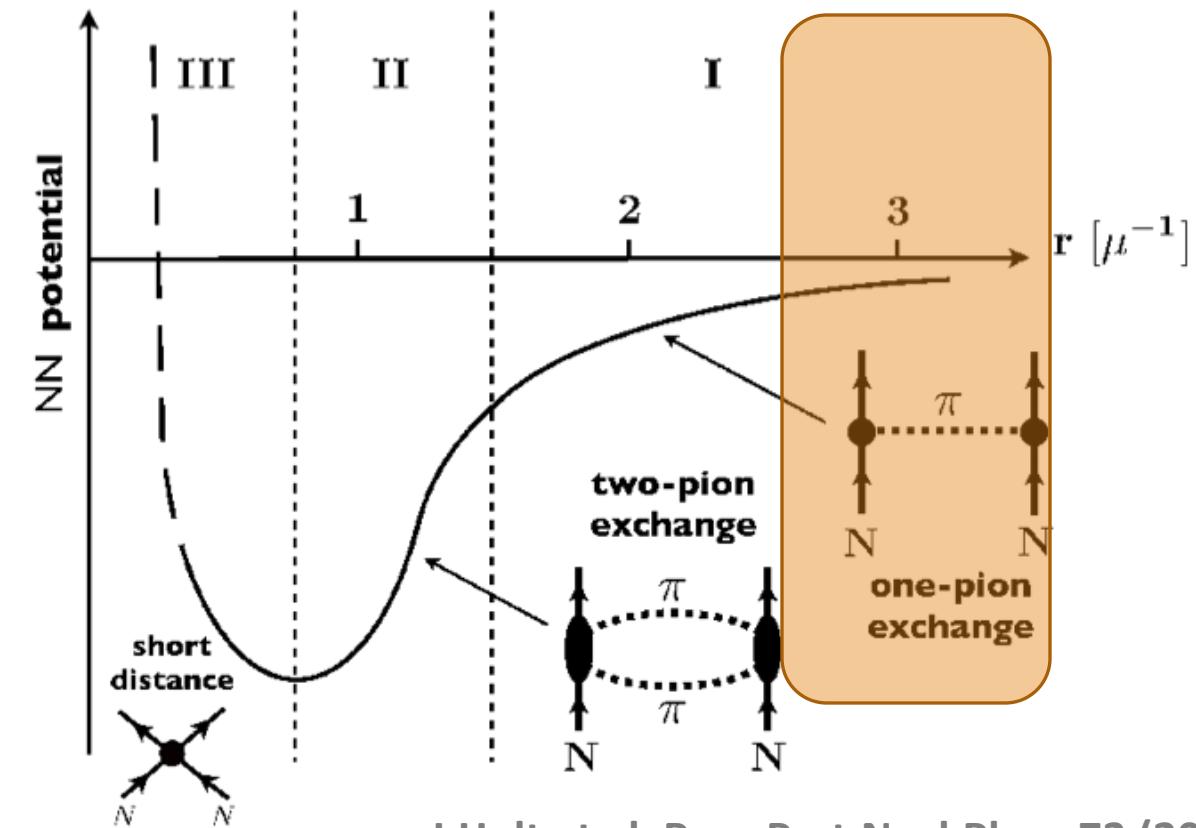
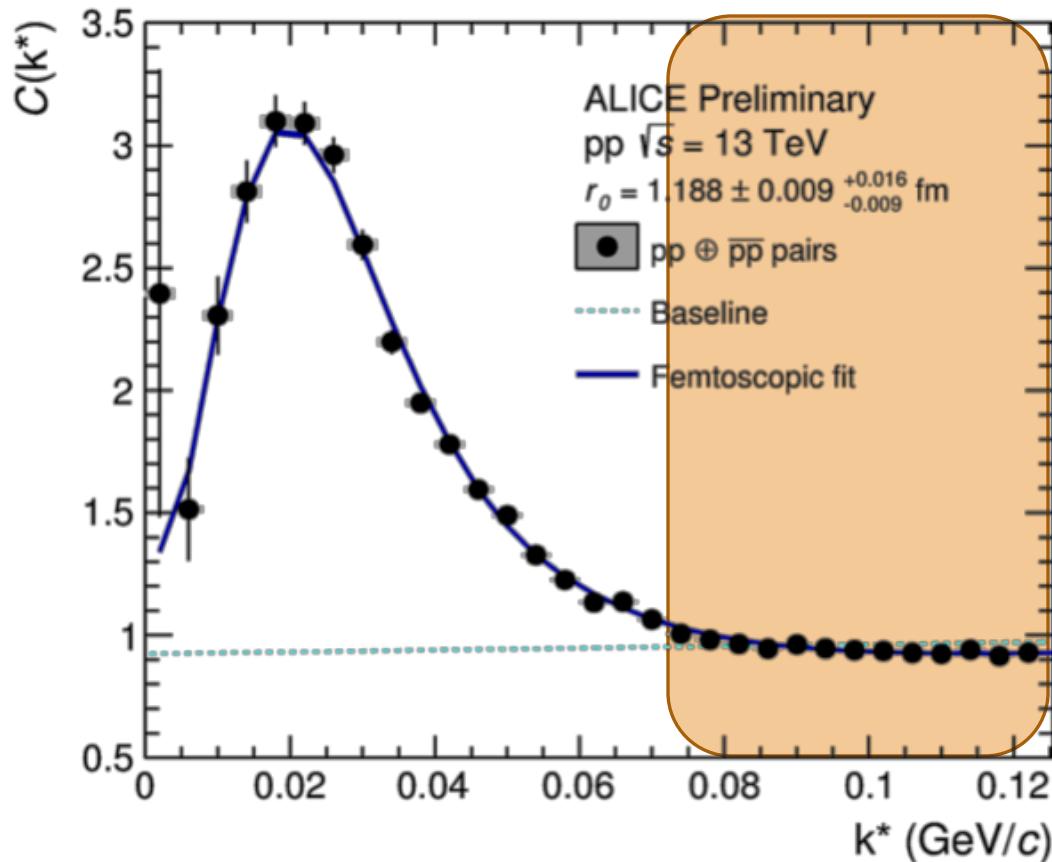
ALI-PREL-144793

# Correlation function and Interactions

$C(k^*) \rightarrow$ 
  
 > 1 attraction

$= 1$  no interaction

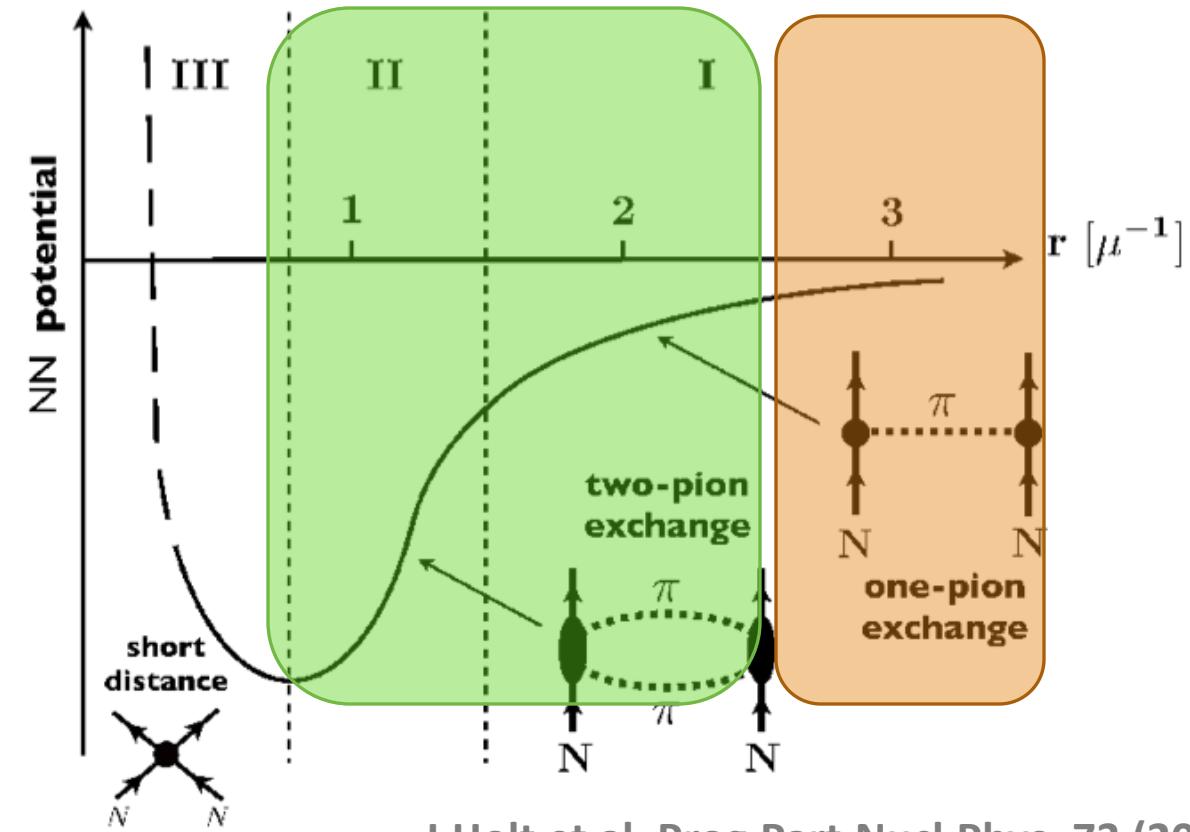
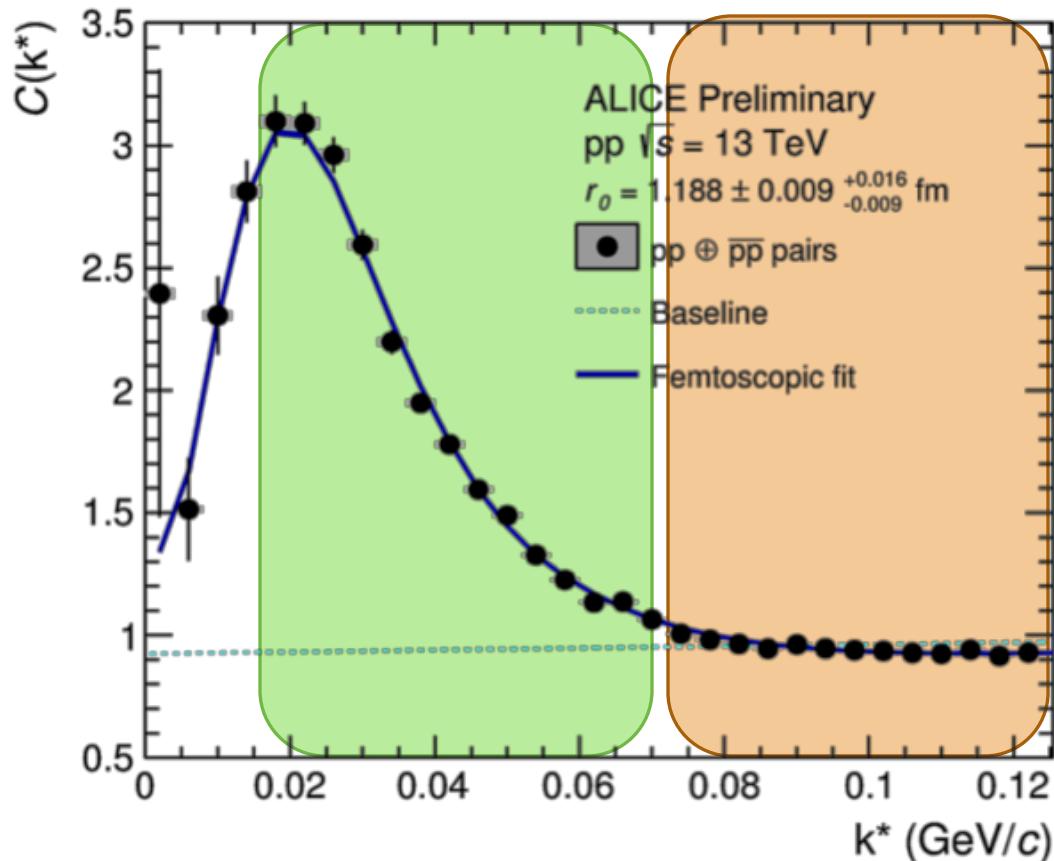
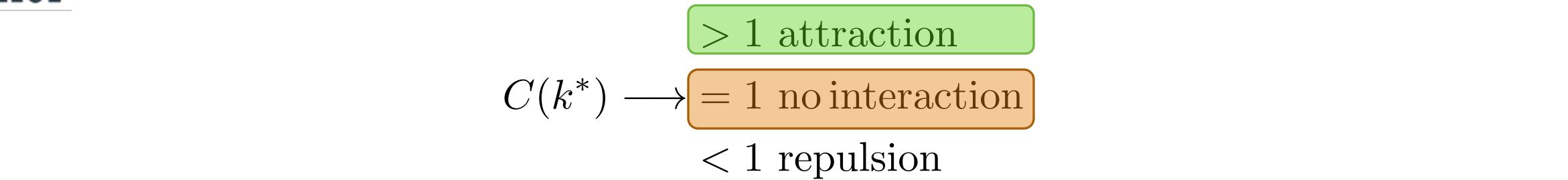
< 1 repulsion



J.Holt et al, Prog.Part.Nucl.Phys. 73 (2013)

ALI-PREL-144793

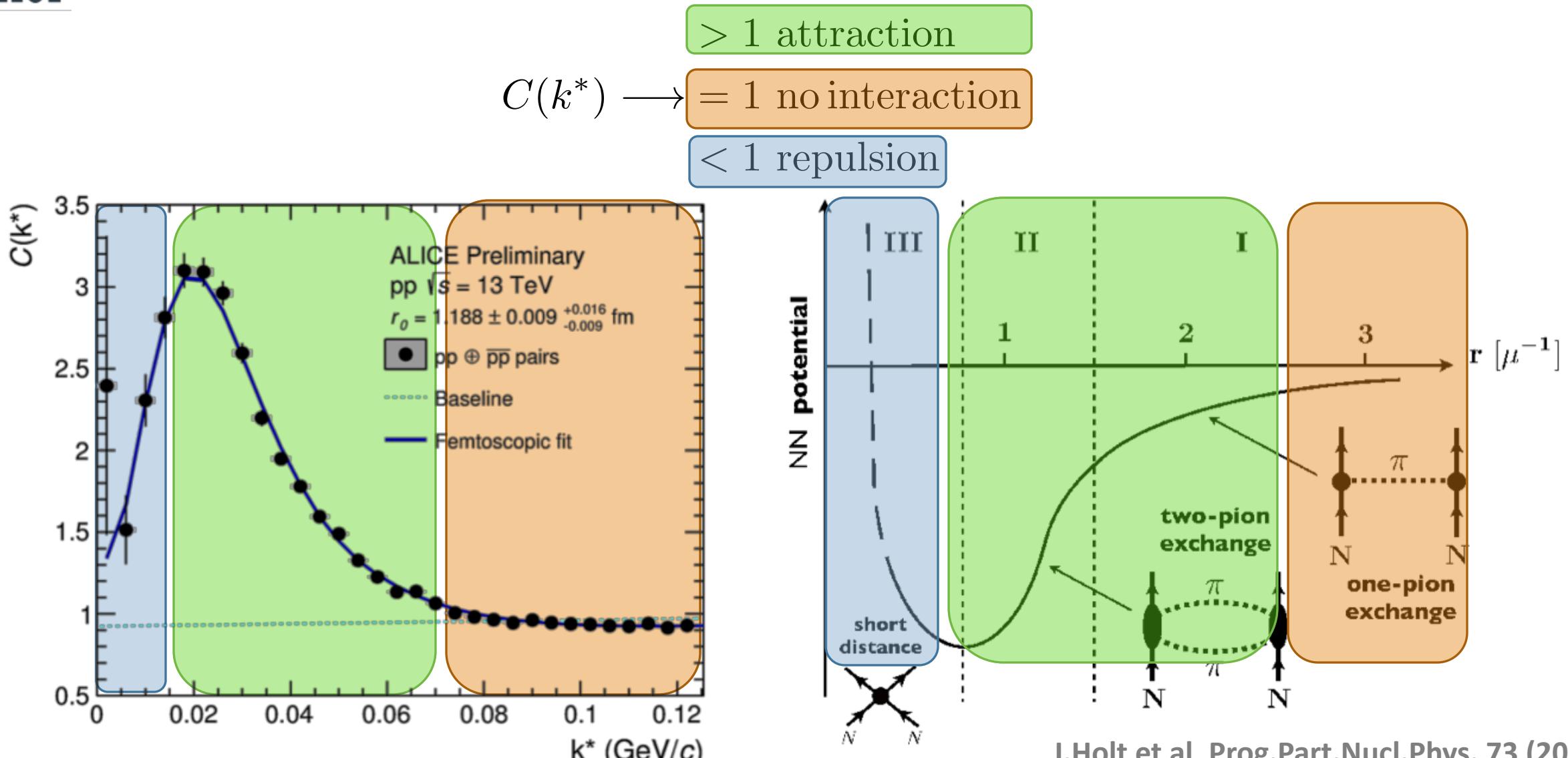
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J. Holt et al, Prog.Part.Nucl.Phys. 73 (2013)

ALI-PREL-144793

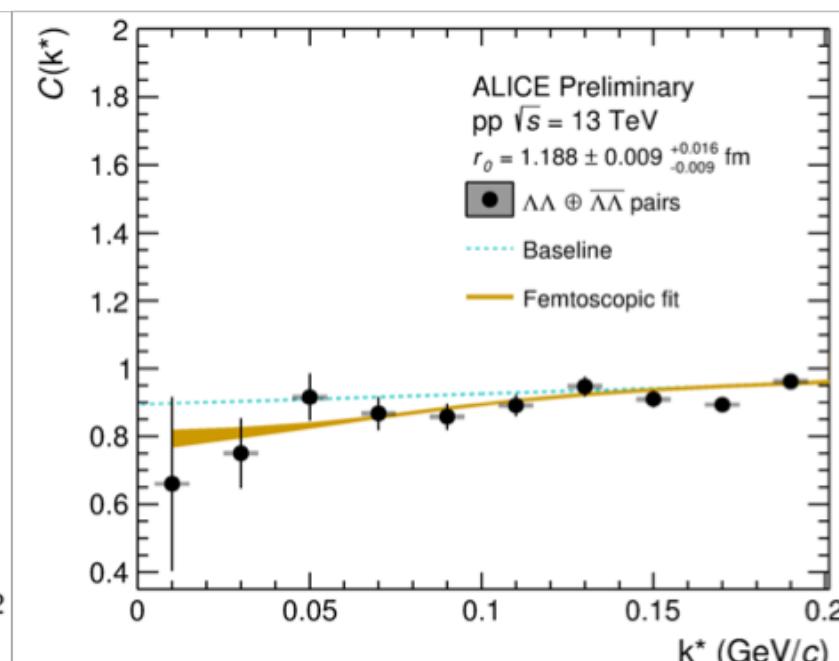
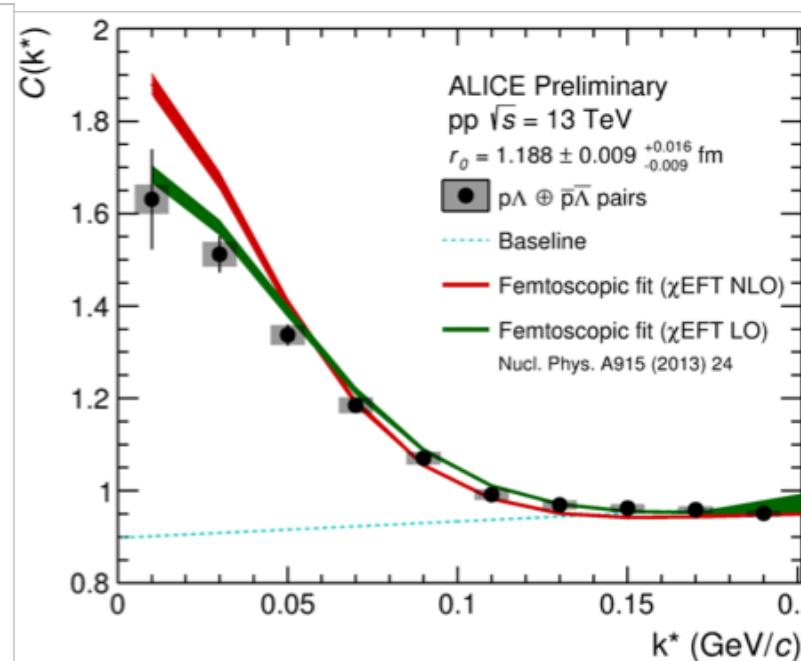
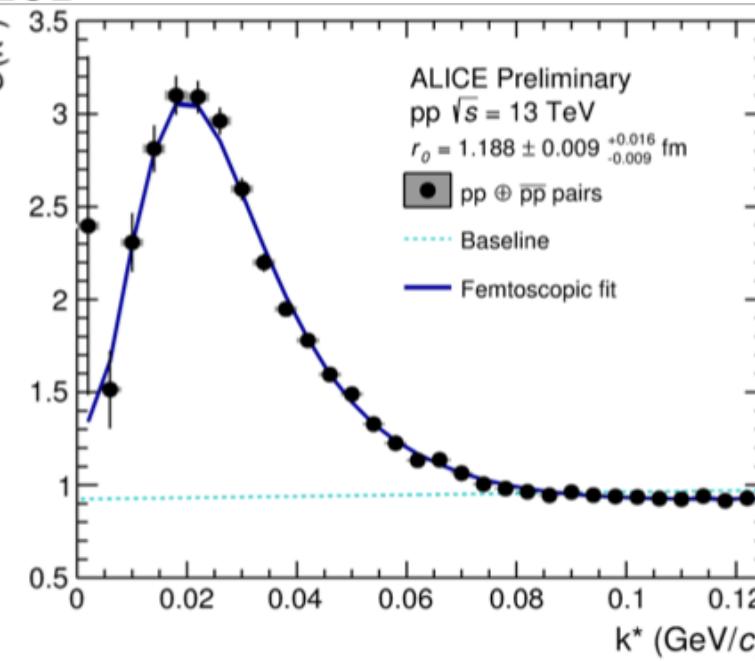
# Correlation function and Interactions





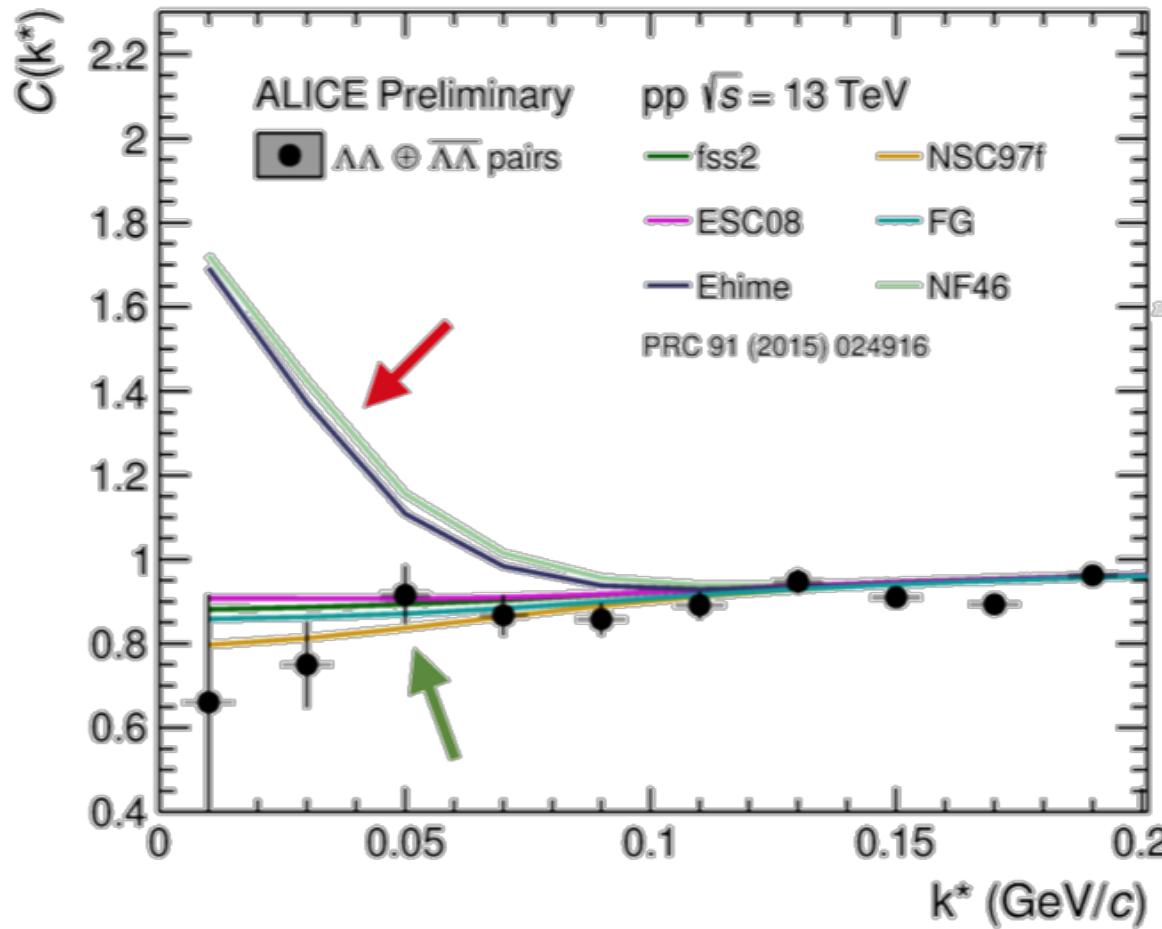
ALICE

# Results - pp collisions $\sqrt{s} = 13$ TeV



- Gaussian source and Argonne  $v_{18}$  potential describes the p-p correlation function
  - Source size of the **pp (7 TeV)** system  $r_0=1.14$  fm (**ALICE Coll. arXiv:1805.12455**)
  - Source size of the **pp (13 TeV)** system  $r_0=1.19$  fm
  - Source size of the **p-Pb (5.02 TeV)** system  $r_0=1.38$  fm
- p- $\Lambda$  correlation  $\Rightarrow$  strong sensitivity to the source  $\Rightarrow$  more investigations on the effect of collective phenomena and resonance decays are ongoing

# $\Lambda-\bar{\Lambda}$ interaction potentials



- Different model predictions for  $\Lambda-\bar{\Lambda}$  correlation function
- $\chi^2$  evaluation with all combined datasets and evaluate the agreement model/data in number of  $\sigma$  deviations

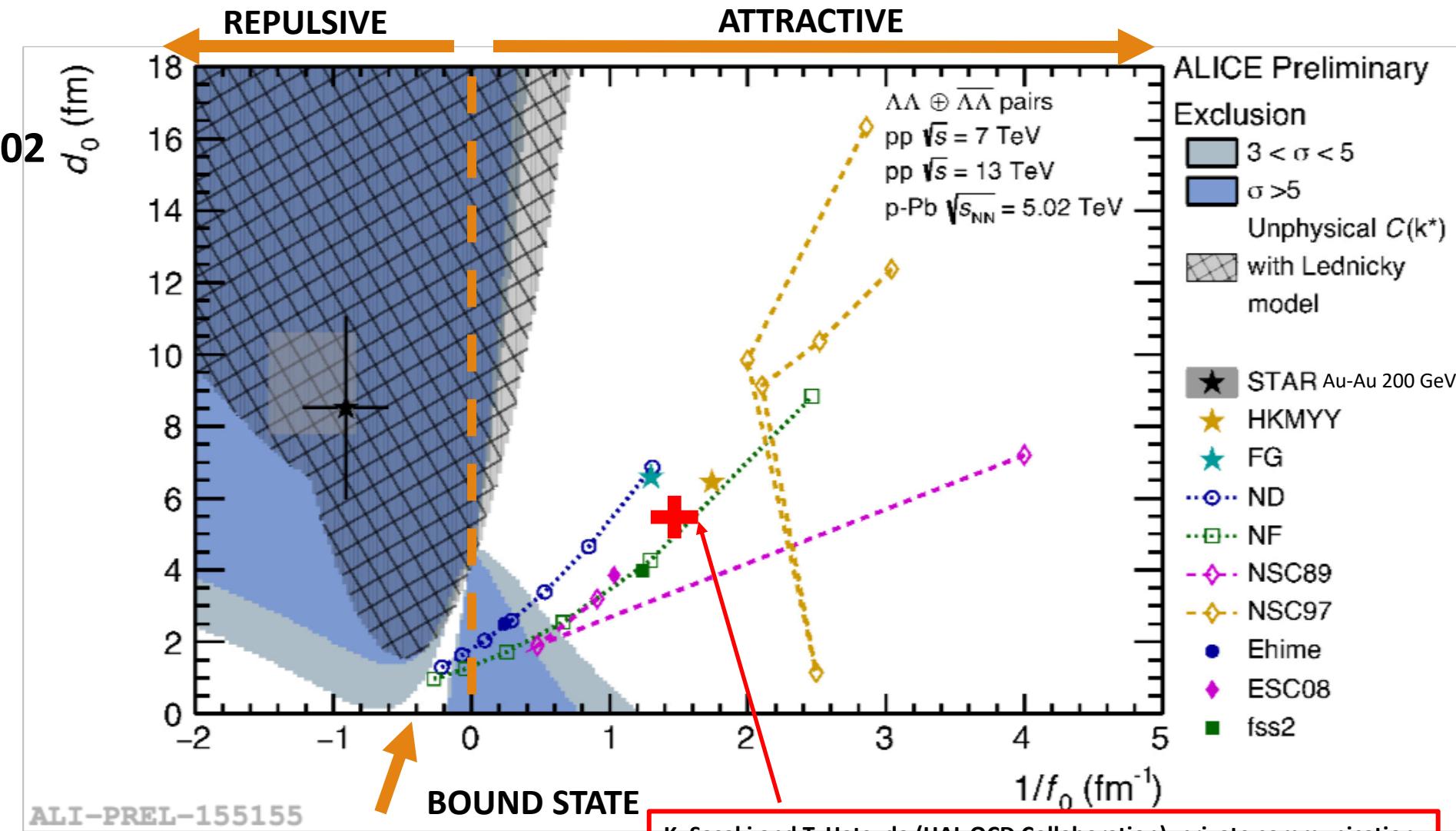


# $\Lambda-\bar{\Lambda}$ Correlations: Combined Exclusion Plot

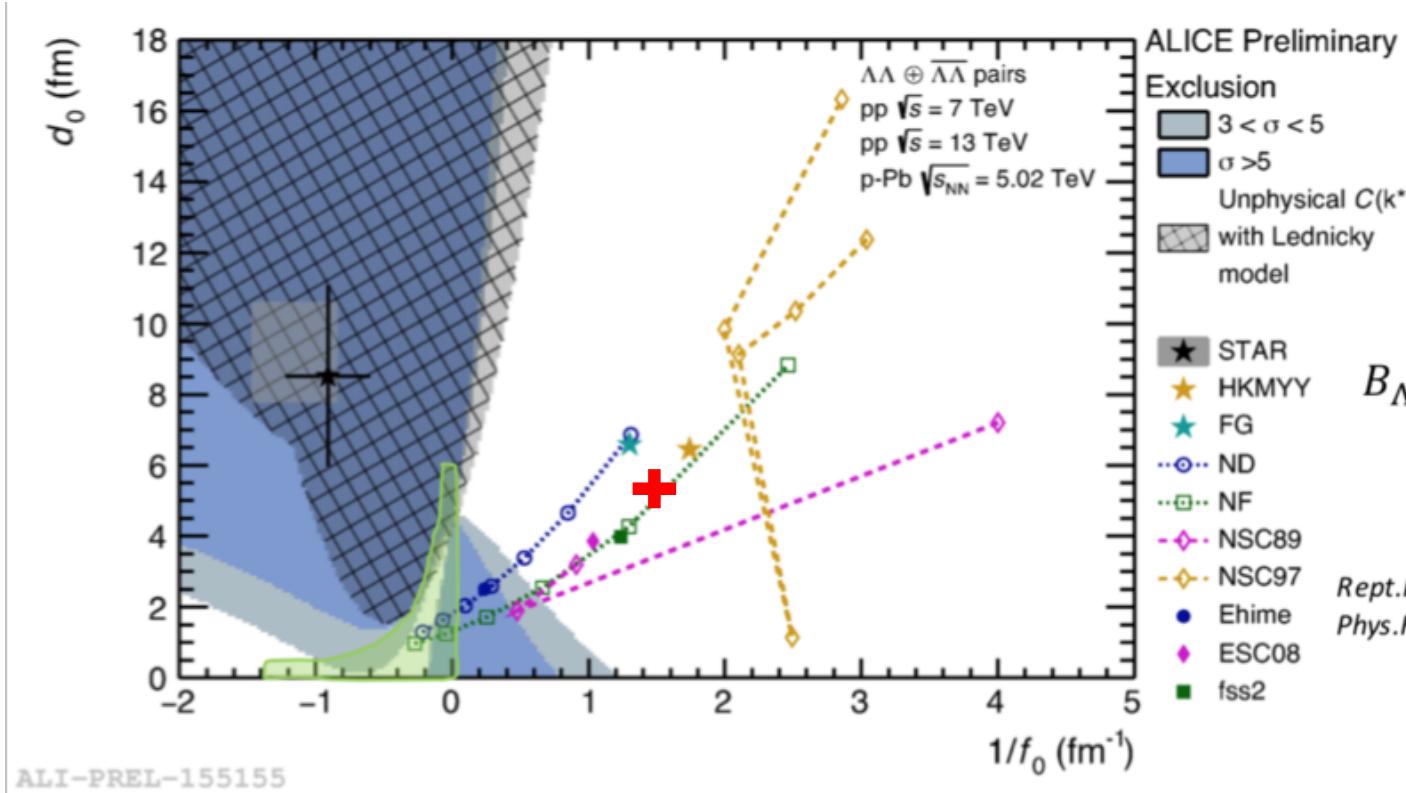


ALICE

- Combination of all available datasets: pp 7 TeV, pp 13 TeV, p-Pb 5.02 TeV
- Full scan** of scattering parameters space with the **Lednicky model**
- Test of the **agreement between data and the prediction by the Lednicky model** by  $n\sigma$



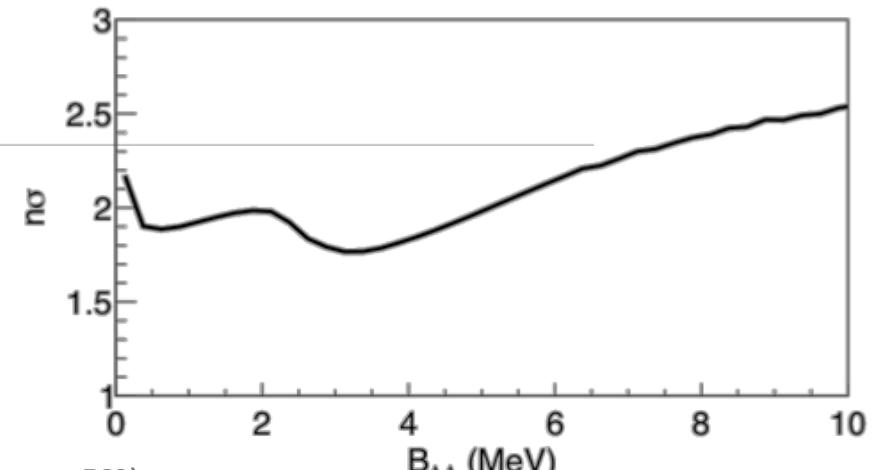
# H-dibaryon binding energies



$$B_{\Lambda\Lambda} = \frac{1}{m_\Lambda d_0^2} \left( 1 - \sqrt{1 + \frac{2d_0}{f_0}} \right)^2$$

Femto sign convention!

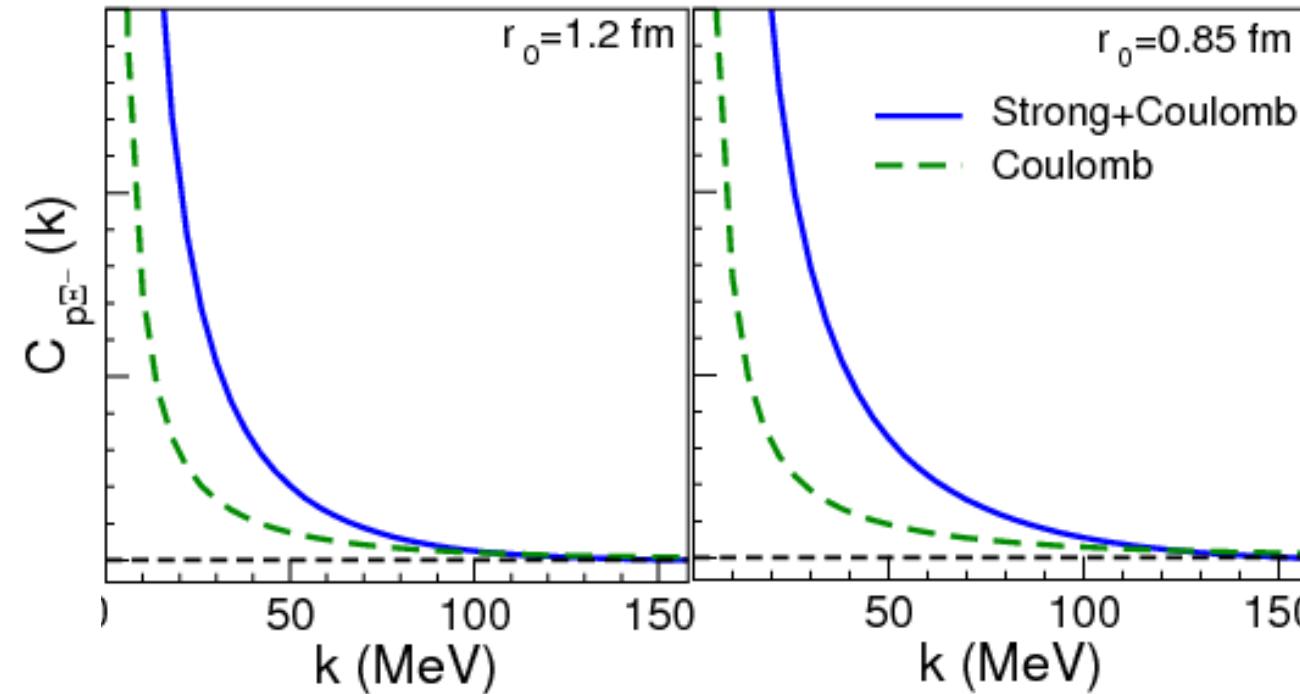
*Rept. Prog. Phys. 80 (2017) no.5, 056001*  
*Phys. Rev. Lett. 120 (2018) 212001*





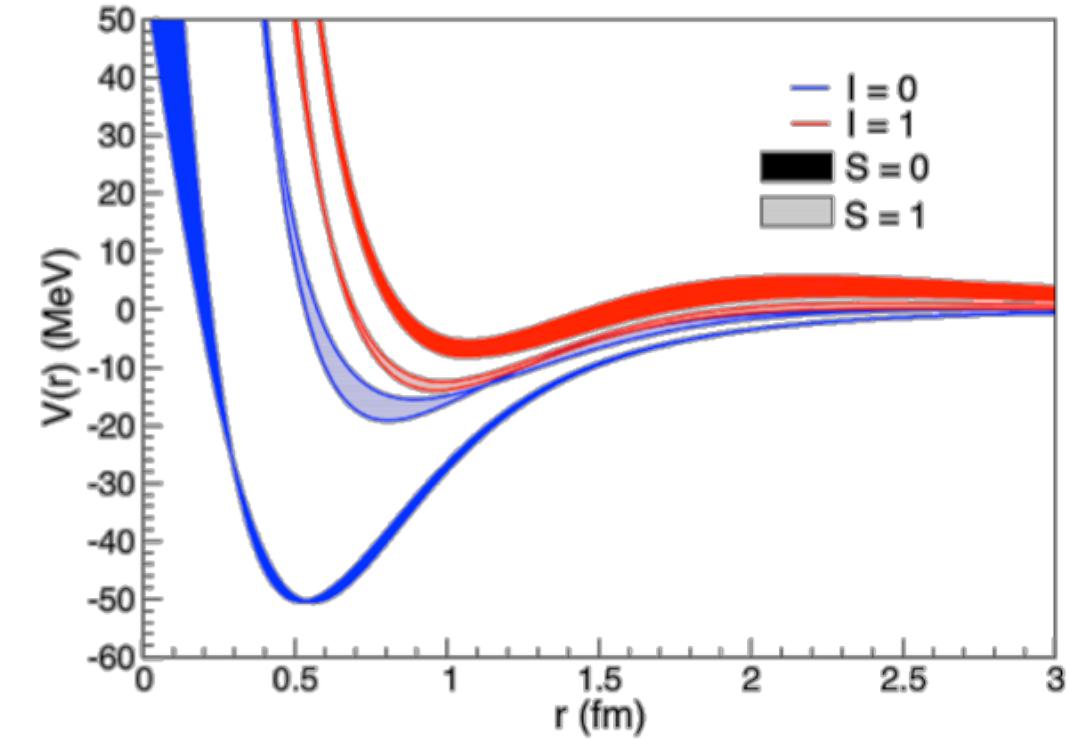
ALICE

# Theoretical p- $\Xi^-$ Correlation function



D.L.Mihaylov, V.M.S, O.W.Arnold, L.Fabbietti, B.Hohlweger,  
A.M.Mathis, Eur.Phys.J. C78 (2018) no.5, 394

- Comparison **Coulomb-only/Coulomb+Strong**
- Strong potential from preliminary calculations from HAL-QCD collaboration



Potential from Hatsuda et al., NPA967 (2017) 856, PoS  
Lattice2016 (2017) 116

$$C(k^*) = \frac{1}{8} (C_{I=0}^{S=0} + C_{I=1}^{S=0}) + \frac{3}{8} (C_{I=0}^{S=1} + C_{I=1}^{S=1})$$

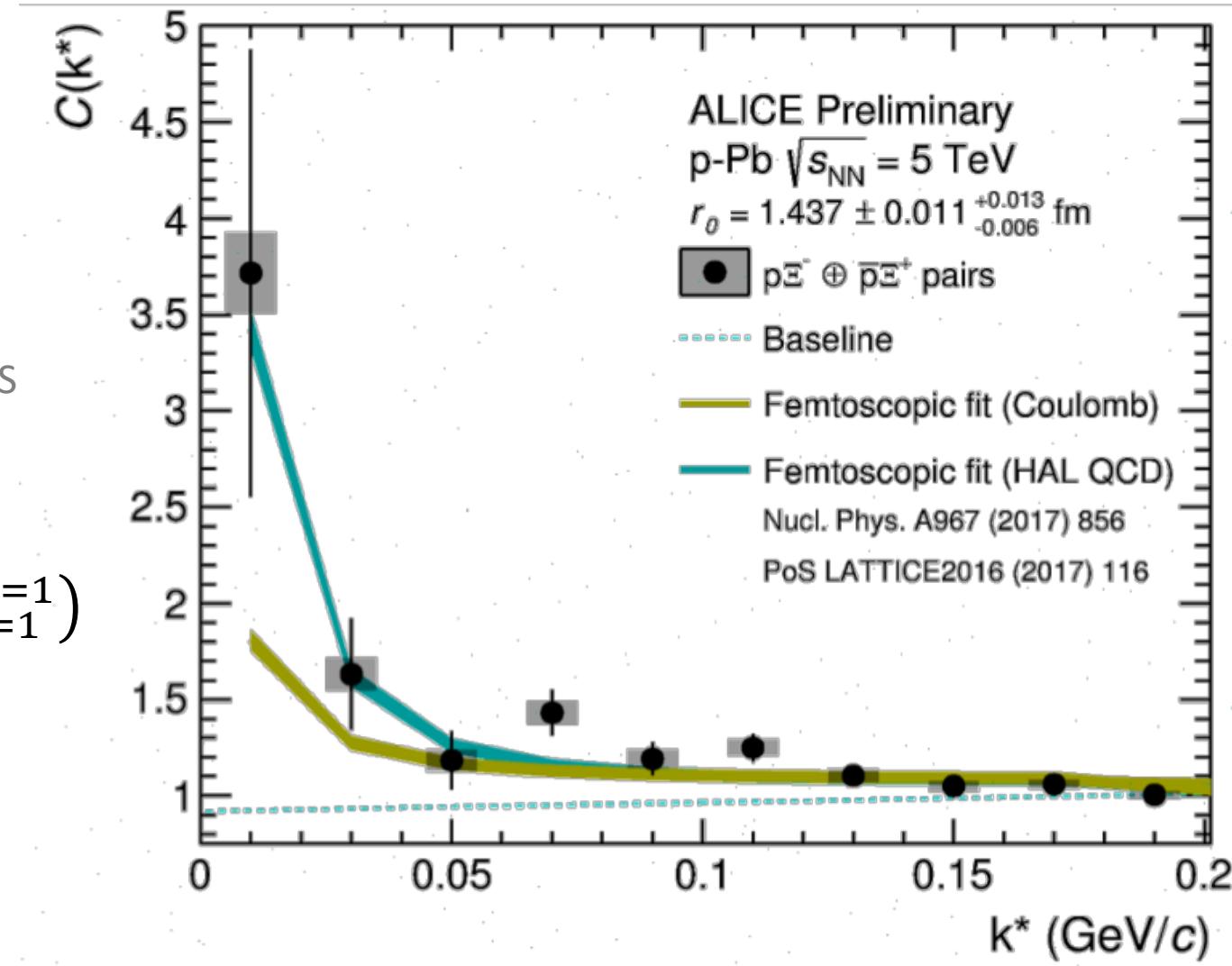
# p- $\Xi^-$ Correlation Function in p-Pb 5.02 TeV

- **First observation of strong attractive interaction in p- $\Xi^-$**
- modeled with preliminary QCD strong potential by the HAL QCD collaboration

Potential from Hatsuda et al., NPA967 (2017) 856, PoS Lattice2016 (2017) 116

$$C(k^*) = \frac{1}{8} (C_{I=0}^{S=0} + C_{I=1}^{S=0}) + \frac{3}{8} (C_{I=0}^{S=1} + C_{I=1}^{S=1})$$

**COULOMB-ONLY  
HYPOTHESIS EXCLUDED  
AROUND  $3\sigma$**



# Summary and Outlook:

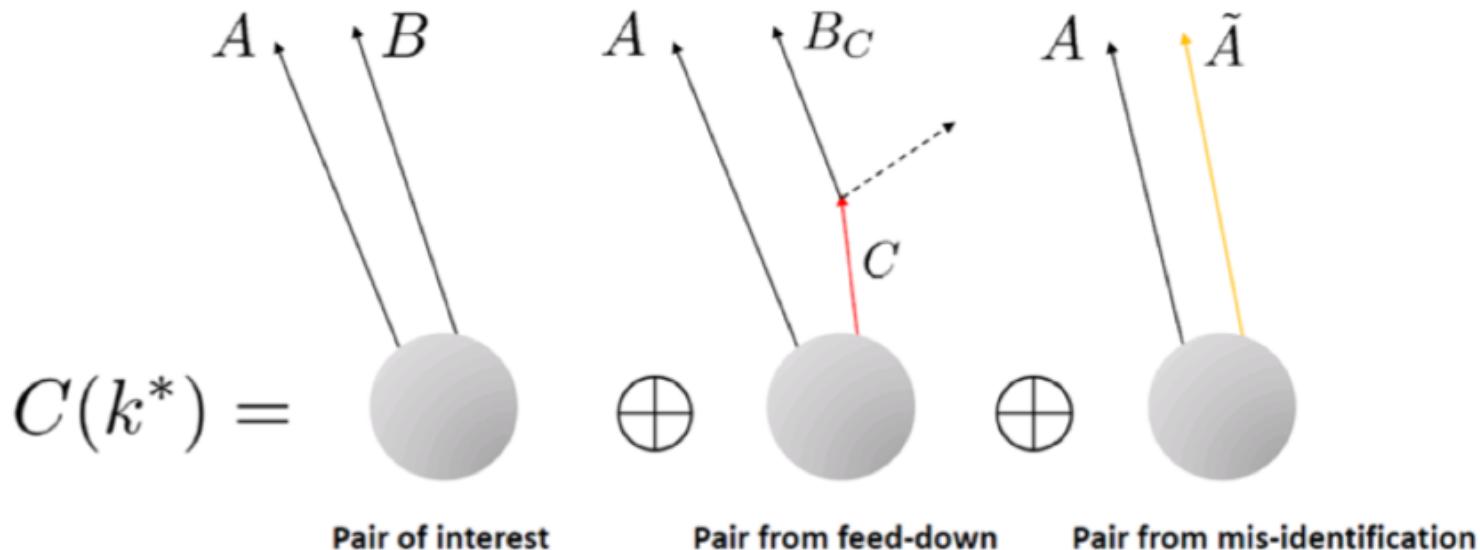
- **Femtoscopy** is an excellent tool to study **interactions of particle pairs**
  - Significant sensitivity to the interaction potentials
  - For hyperons, accesses novel regions not constrained by scattering experiments
- **$\Lambda$ - $\Lambda$  analysis** strongly constrains the parameter space for the  $\Lambda$ - $\Lambda$  interaction  $\Rightarrow$  the existence of H-dibaryons seems to be disfavored (ALICE coll., Phys. Lett. B 752)
- Observation of **attractive p- $\Xi^-$  interaction for the first time**  $\Rightarrow$  set constraints on the average potential of  $\Xi$  hyperons at finite density for NS EoS
- Extending the study to more hyperons-nucleon pairs (p- $\Omega$ , p- $\Sigma^0$ ) and in the future to three-body interactions
- In RUN3 ( from 2021 on) we expect factor 100 in statistics!

# Backup Slides

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# The experimental correlation function

$$C(\vec{p}_a, \vec{p}_b) = \frac{P(\vec{p}_a, \vec{p}_b)}{P(\vec{p}_a)P(\vec{p}_b)} = \mathcal{N} \frac{N_{SE}(k)}{N_{ME}(k)}$$



# Decomposition of the p-p correlation function

$$\{pp\} = pp + p_{\Lambda}p + p_{\Lambda} + p_{\Lambda} + p_{\Sigma^+}p + p_{\Sigma^+}p_{\Sigma^+} + p_{\Lambda}p_{\Sigma^+} + \bar{p}p + \bar{p}p_{\Lambda} + \bar{p}p_{\Sigma^+} + \bar{p}\bar{p},$$

- Purity from MC (Pythia 8)
- Feed-down fractions from MC template fits to the  $DCA_{xy}$  distribution

Pair	p-p $\lambda$ [%]
pp	75.19
$p_{\Lambda}p$	15.06
$p_{\Lambda}p_{\Lambda}$	0.75
$p_{\Sigma^+}p$	6.46
$p_{\Sigma^+}p_{\Sigma^+}$	0.14
$p_{\Lambda}p_{\Sigma^+}$	0.65
$\bar{p}p$	1.52
$\bar{p}p_{\Lambda}$	0.15
$\bar{p}p_{\Sigma^+}$	0.07
$\bar{p}\bar{p}$	0.01

# Decomposition of the p- $\Lambda$ correlation function

$$\begin{aligned} \{p\Lambda\} = & p\Lambda + p\Lambda_{\Xi^-} + p\Lambda_{\Xi^0} + p\Lambda_{\Sigma^0} + p_\Lambda\Lambda + p_\Lambda\Lambda_{\Xi^-} + p_\Lambda\Lambda_{\Xi^0} + p_\Lambda\Lambda_{\Sigma^0} \\ & + p_{\Sigma^+}\Lambda + p_{\Sigma^+}\Lambda_{\Xi^-} + p_{\Sigma^+}\Lambda_{\Xi^0} + p_{\Sigma^+}\Lambda_{\Sigma^0} + \tilde{p}\Lambda + \tilde{p}\Lambda_{\Xi^-} + \tilde{p}\Lambda_{\Xi^0} + \tilde{p}\Lambda_{\Sigma^0} \\ & + p\tilde{\Lambda} + p_\Lambda\tilde{\Lambda} + p_{\Sigma^+}\tilde{\Lambda} + \tilde{p}\tilde{\Lambda}. \end{aligned}$$

- Purity from fits to the invariant mass distribution
- Feed-down fractions from MC template fits to the  $\cos\alpha$  distribution

Pair	p- $\Lambda$	Pair	p- $\Lambda$
	$\lambda$ [%]		$\lambda$ [%]
p $\Lambda$	52.42	$\tilde{p}\Lambda$	0.53
p $\Lambda_{\Xi^-}$	6.94	$\tilde{p}\Lambda_{\Xi^-}$	0.07
p $\Lambda_{\Xi^0}$	6.94	$\tilde{p}\Lambda_{\Xi^0}$	0.07
p $\Lambda_{\Sigma^0}$	17.47	$\tilde{p}\Lambda_{\Sigma^0}$	0.18
p $_\Lambda\Lambda$	5.25	p $\tilde{\Lambda}$	2.95
p $_\Lambda\Lambda_{\Xi^-}$	0.69	p $_\Lambda\tilde{\Lambda}$	0.30
p $_\Lambda\Lambda_{\Xi^0}$	0.69	p $_{\Sigma^+}\tilde{\Lambda}$	0.13
p $_\Lambda\Lambda_{\Sigma^0}$	1.75	$\tilde{p}\tilde{\Lambda}$	0.03
p $_{\Sigma^+}\Lambda$	2.25		
p $_{\Sigma^+}\Lambda_{\Xi^-}$	0.30		
p $_{\Sigma^+}\Lambda_{\Xi^0}$	0.30		
p $_{\Sigma^+}\Lambda_{\Sigma^0}$	0.75		

# Decomposition of the $\Lambda$ - $\Lambda$ correlation function

$$\begin{aligned} \{\Lambda\Lambda\} = & \Lambda\Lambda + \Lambda\Lambda_{\Sigma^0} + \Lambda_{\Sigma^0}\Lambda_{\Sigma^0} + \Lambda\Lambda_{\Xi^0} + \Lambda_{\Xi^0}\Lambda_{\Xi^0} + \Lambda\Lambda_{\Xi^-} \\ & + \Lambda_{\Xi^-}\Lambda_{\Xi^-} + \Lambda_{\Sigma^0}\Lambda_{\Xi^0} + \Lambda_{\Sigma^0}\Lambda_{\Xi^-} + \Lambda_{\Xi^0}\Lambda_{\Xi^-} \\ & + \tilde{\Lambda}\Lambda + \tilde{\Lambda}\Lambda_{\Sigma^0} + \tilde{\Lambda}\Lambda_{\Xi^-} + \tilde{\Lambda}\Lambda_{\Xi^0} + \tilde{\Lambda}\tilde{\Lambda}. \end{aligned}$$

Lambda properties obtained from the  $\Lambda$  purity and the  $\cos\alpha$  template fits

Pair	$\Lambda$ - $\Lambda$
Pair	$\lambda$ [%]
$\Lambda\Lambda$	36.54
$\Lambda\Lambda_{\Sigma^0}$	24.36
$\Lambda_{\Sigma^0}\Lambda_{\Sigma^0}$	4.06
$\Lambda\Lambda_{\Xi^0}$	9.67
$\Lambda_{\Xi^0}\Lambda_{\Xi^0}$	0.64
$\Lambda\Lambda_{\Xi^-}$	9.67
$\Lambda_{\Xi^-}\Lambda_{\Xi^-}$	0.64
$\Lambda_{\Sigma^0}\Lambda_{\Xi^0}$	3.22
$\Lambda_{\Sigma^0}\Lambda_{\Xi^-}$	3.22
$\Lambda_{\Xi^0}\Lambda_{\Xi^-}$	1.28

Pair	$\Lambda$ - $\Lambda$
Pair	$\lambda$ [%]
$\tilde{\Lambda}\Lambda$	4.11
$\tilde{\Lambda}\Lambda_{\Sigma^0}$	1.37
$\tilde{\Lambda}\Lambda_{\Xi^0}$	0.54
$\tilde{\Lambda}\Lambda_{\Xi^-}$	0.54
$\tilde{\Lambda}\tilde{\Lambda}$	0.12

# Decomposition of the p- $\Xi$ correlation function

$$\begin{aligned} \{p\Xi^-\} = & p\Xi^- + p\Xi_{\Xi^-(1530)}^- + p\Xi_{\Xi^0(1530)}^- + p\Xi_\Omega^- + p_\Lambda\Xi^- + p_\Lambda\Xi_{\Xi^-(1530)}^- \\ & + p_\Lambda\Xi_{\Xi^0(1530)}^- + p_\Lambda\Xi_\Omega^- + p_\Sigma^+\Xi^- + p_\Sigma^+\Xi_{\Xi^-(1530)}^- + p_\Sigma^+\Xi_{\Xi^0(1530)}^- + p_\Sigma^+\Xi_\Omega^- \\ & + \tilde{p}\Xi^- + \tilde{p}\Xi_{\Xi^-(1530)}^- + \tilde{p}\Xi_{\Xi^0(1530)}^- + \tilde{p}\Xi_\Omega^- + p\tilde{\Xi}^- + p_\Lambda\tilde{\Xi}^- + p_\Sigma^+\tilde{\Xi}^- + \tilde{p}\Xi^-. \end{aligned}$$

Feeding from

- $\Omega$  (BR very small)
- $\Xi^0(1530)$  and  $\Xi^-(1530)$ 
  - Isospin partners: assume to be produced in the same amount
  - $\Xi(1530)/\Xi^- = 0.32$   
(<https://doi.org/10.1140/epjc/s10052-014-3191-x>)
  - $BR(\Xi^0(1530) \rightarrow \Xi^-) = 2/3$
  - $BR(\Xi^-(1530) \rightarrow \Xi^-) = 1/3$

p- $\Xi$		p- $\Xi$	
Pair	$\lambda$ [%]	Pair	$\lambda$ [%]
p $\Xi^-$	52.40	$\tilde{p}\Xi^-$	0.53
p $\Xi_{\Xi^-(1530)}^-$	8.32	$\tilde{p}\Xi_{\Xi^-(1530)}^-$	0.08
p $\Xi_{\Xi^0(1530)}^-$	16.65	$\tilde{p}\Xi_{\Xi^0(1530)}^-$	0.17
p $\Xi_\Omega^-$	0.67	$\tilde{p}\Xi_\Omega^-$	0.01
p $\Lambda\Xi^-$	5.25	p $\Xi^-$	8.67
p $\Lambda\Xi_{\Xi^-(1530)}^-$	0.83	p $\Lambda\tilde{\Xi}^-$	0.87
p $\Lambda\Xi_{\Xi^0(1530)}^-$	1.67	p $\Sigma^+\tilde{\Xi}^-$	2.25
p $\Lambda\Xi_\Omega^-$	0.07	$\tilde{p}\Xi^-$	0.09
p $\Sigma^+\Xi^-$	2.25		
p $\Sigma^+\Xi_{\Xi^-(1530)}^-$	0.36		
p $\Sigma^+\Xi_{\Xi^0(1530)}^-$	0.71		
p $\Sigma^+\Xi_\Omega^-$	0.03		

# Some numbers

**pp@13 TeV**

Particle	# baryons
p	$113.7 \times 10^6$
$\bar{p}$	$97.4 \times 10^6$
$\Lambda$	$22.3 \times 10^6$
$\bar{\Lambda}$	$21.0 \times 10^6$
$\Xi^-$	$509 \times 10^3$
$\Xi^+$	$527 \times 10^3$

Pair	# of pairs $k^* < 200$ MeV/c
$p - p$	$190 \times 10^3$
$\bar{p} - \bar{p}$	$140 \times 10^3$
$p - \Lambda$	$62 \times 10^3$
$\bar{p} - \bar{\Lambda}$	$49 \times 10^3$
$\Lambda - \Lambda$	5659
$\bar{\Lambda} - \bar{\Lambda}$	5243
$p - \Xi^-$	407
$\bar{p} - \Xi^+$	364

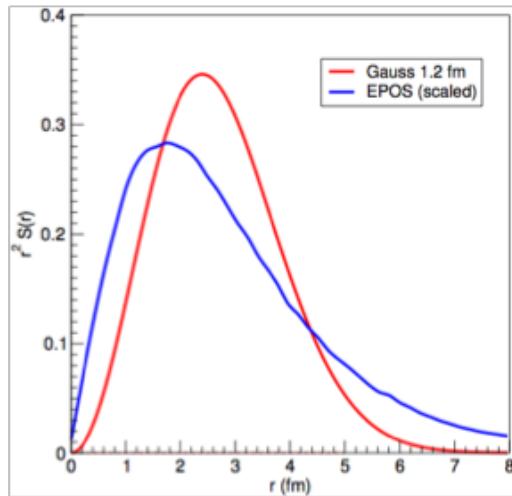
Particle	Candidates
p	$155 \times 10^6$
$\bar{p}$	$133 \times 10^6$
$\Lambda$	$26 \times 10^6$
$\bar{\Lambda}$	$24 \times 10^6$
$\Xi$	$0.9 \times 10^6$
$\bar{\Xi}$	$0.9 \times 10^6$

**P-Pb @ 5.02 TeV**

Combinations	Pairs ( $k^* < 200$ MeV/c)
$p - p$	$517 \times 10^3$
$\bar{p} - \bar{p}$	$370 \times 10^3$
$p - \Lambda$	$127 \times 10^3$
$\bar{p} - \bar{\Lambda}$	$62 \times 10^3$
$\Lambda - \Lambda$	$13 \times 10^3$
$\bar{\Lambda} - \bar{\Lambda}$	$12 \times 10^3$
$p - \Xi$	$1.8 \times 10^3$
$\bar{p} - \bar{\Xi}$	$1.3 \times 10^3$

# Epos Source

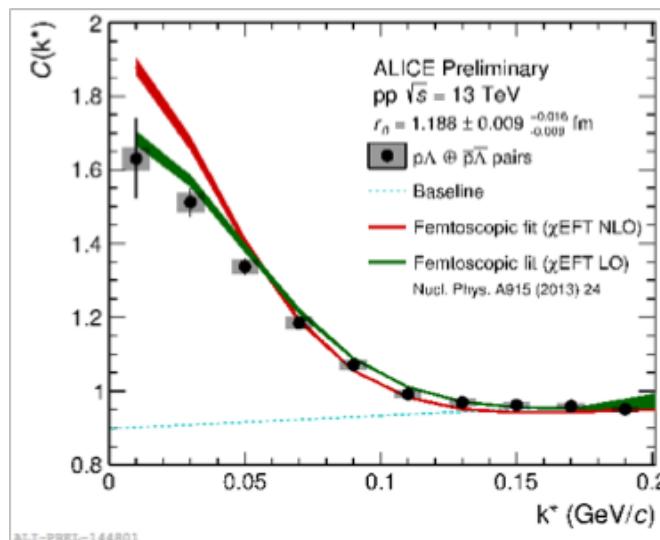
(D.L.Mihaylov et al. Eur.Phys.J. C78 (2018) no.5,394)



Gaussian Source

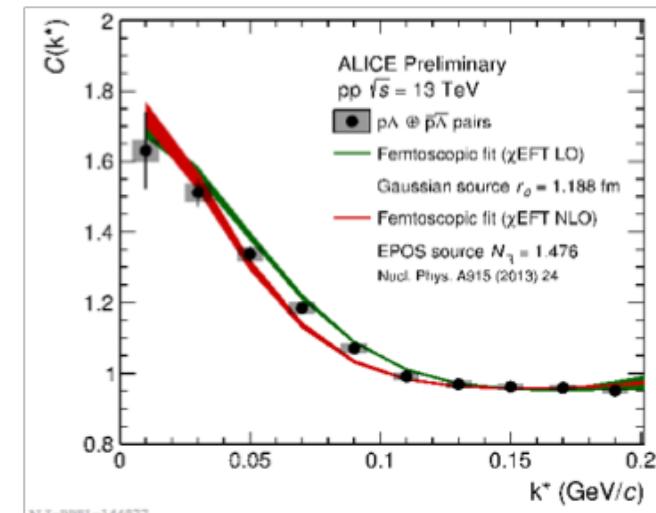
Source with a non-Gaussian shape  
Maximum at smaller R and long tail

Test with 'expanded' EPOS for pp 13 TeV

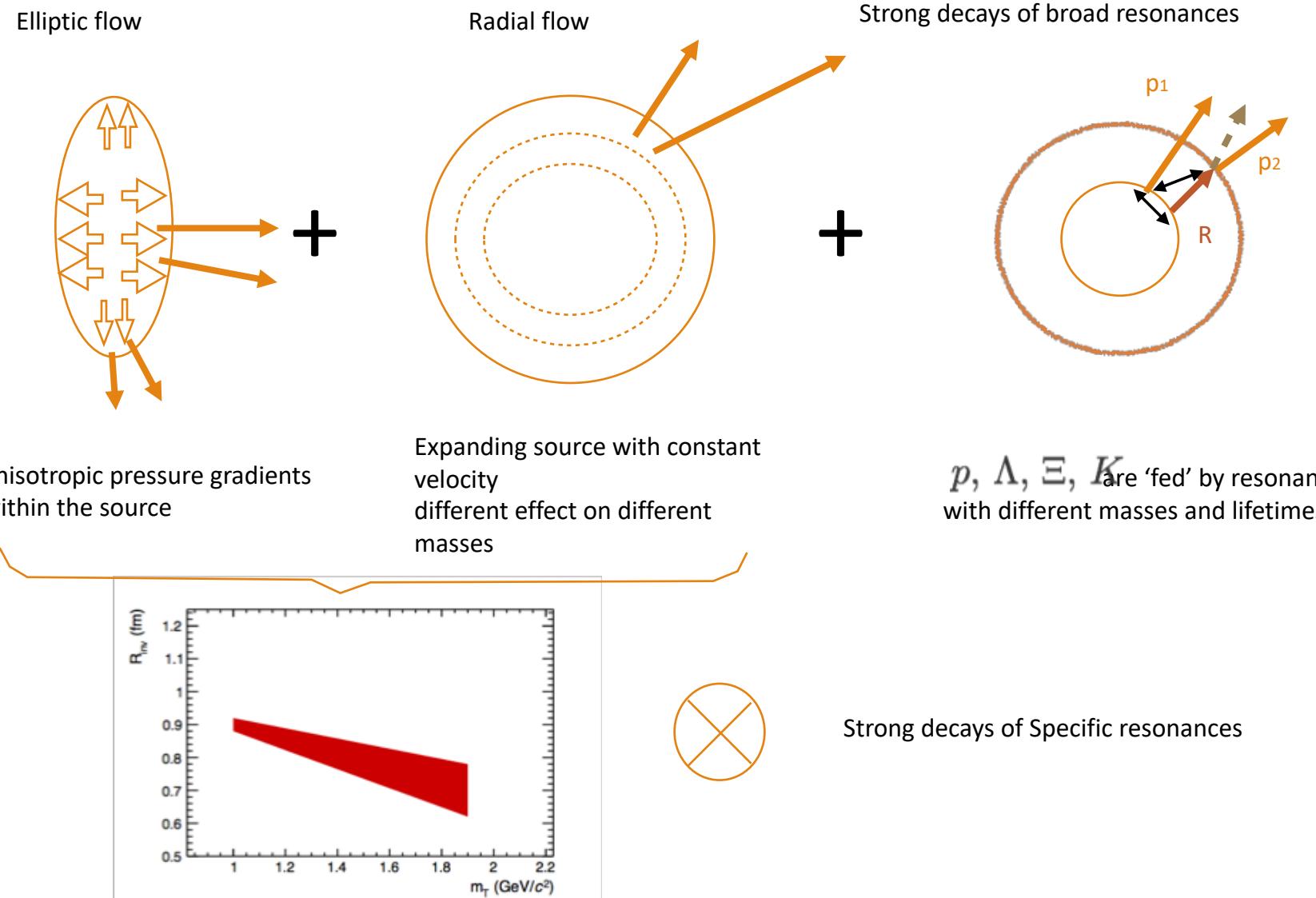


Better Agreement with 'expanded' EPOS for NLO for p-Lambda correlations

EPOS Source

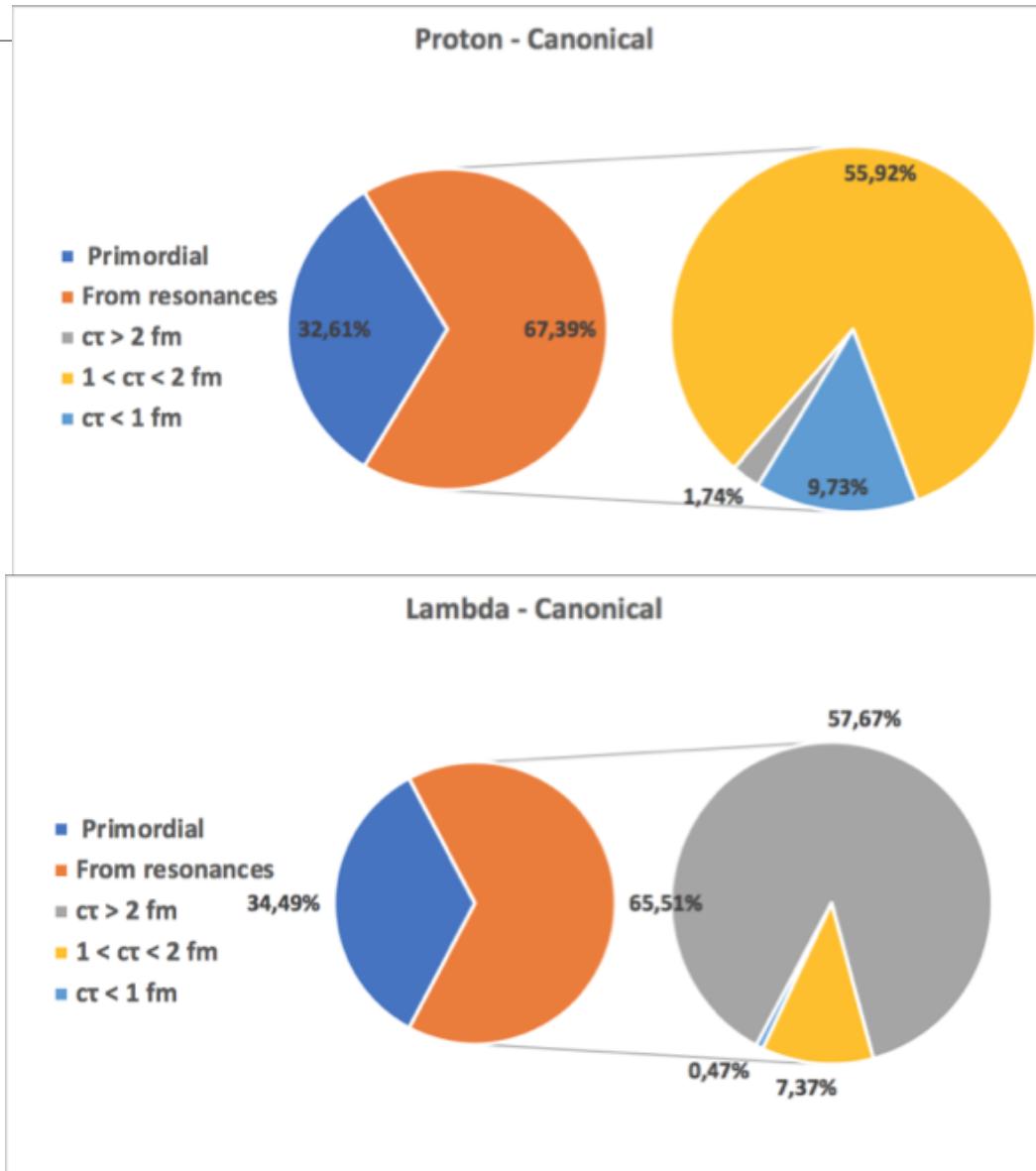


# Collective Effects and Strong Decays



# Which Resonances?

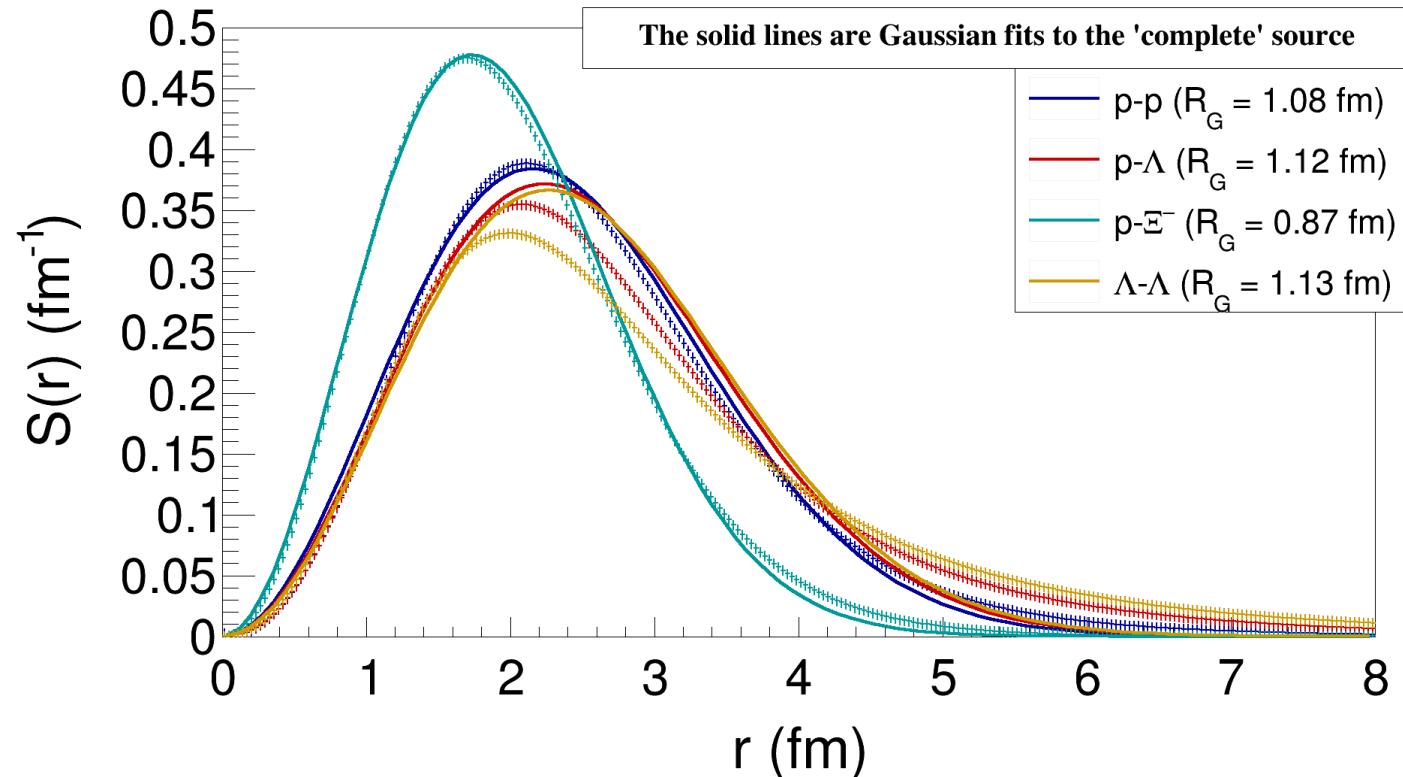
Courtesy F. Becattini



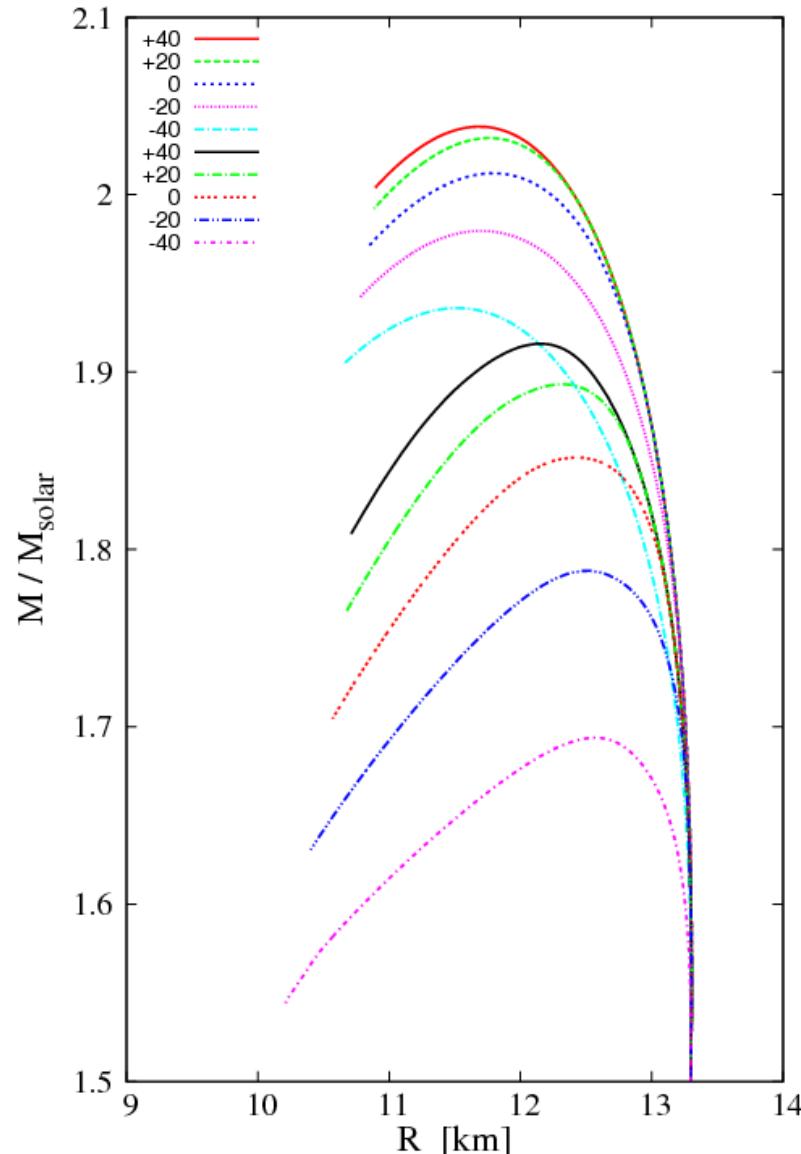
$c\tau > 2 \text{ fm} (2.5\text{-}12 \text{ fm})$	$\Sigma(1670), \Lambda(1520), \dots$
$1 < c\tau < 2 \text{ fm}$	$\Lambda(1600), \Delta(1232), \dots$
$c\tau < 1 \text{ fm}$	$\Delta(1600), N(1440), \dots$
$M_{\text{eff}} [\text{MeV}]$	$\tau_{\text{eff}} [\text{fm}]$
1368.92	1.6

$c\tau > 2 \text{ fm} (\text{mainly } 5 \text{ fm})$	$\Sigma(1385) \text{ mainly.}$
$1 < c\tau < 2 \text{ fm}$	$N(1710), \Sigma(1915), \dots$
$c\tau < 1 \text{ fm}$	$N(1720) \text{ only}$
$M_{\text{eff}} [\text{MeV}]$	$\tau_{\text{eff}} [\text{fm}]$
1461.46	4.78

# Source Distributions



# p- $\Xi^-$ interaction and Neutron Stars

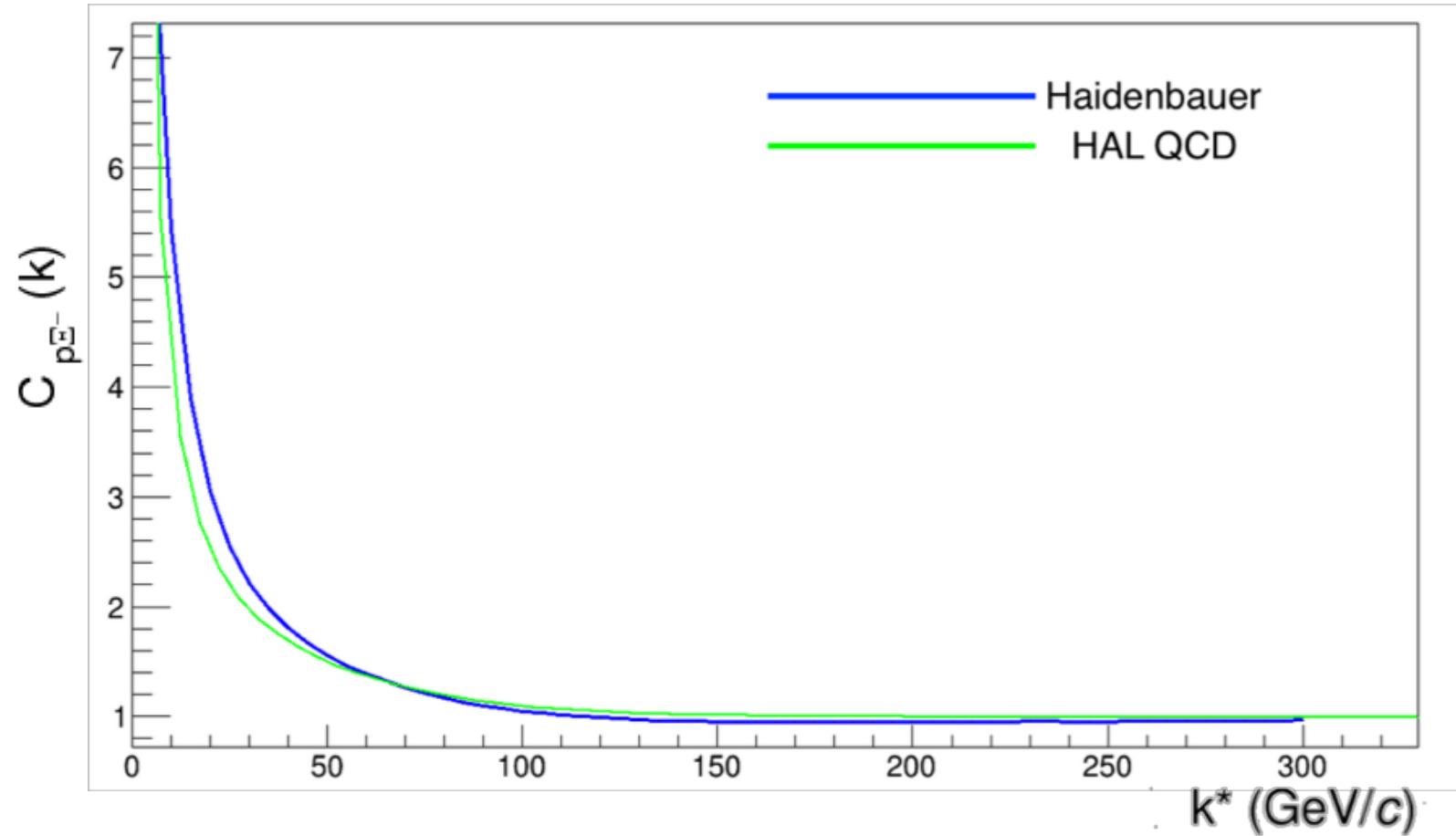


- If the in-medium N- $\Xi$  potential is not attractive, the EoS becomes stiffer
- Ultimately this drives the allowed mass of neutron stars up (even above 2 solar masses)
- A complication: in the case of neutron stars only the n- $\Xi$  interaction is important, but for us neutrons are not detectable

	$I=0$	$I=1$	Detectable
$n-\Xi^-$	X	✓	No
$p-\Xi^0$	X	✓	Difficult
$p-\Xi^-$	✓	✓	Yes

- Experimentally challenging to isolate the  $I=1$  channel

# Theoretical p- $\Xi^-$ Correlation function

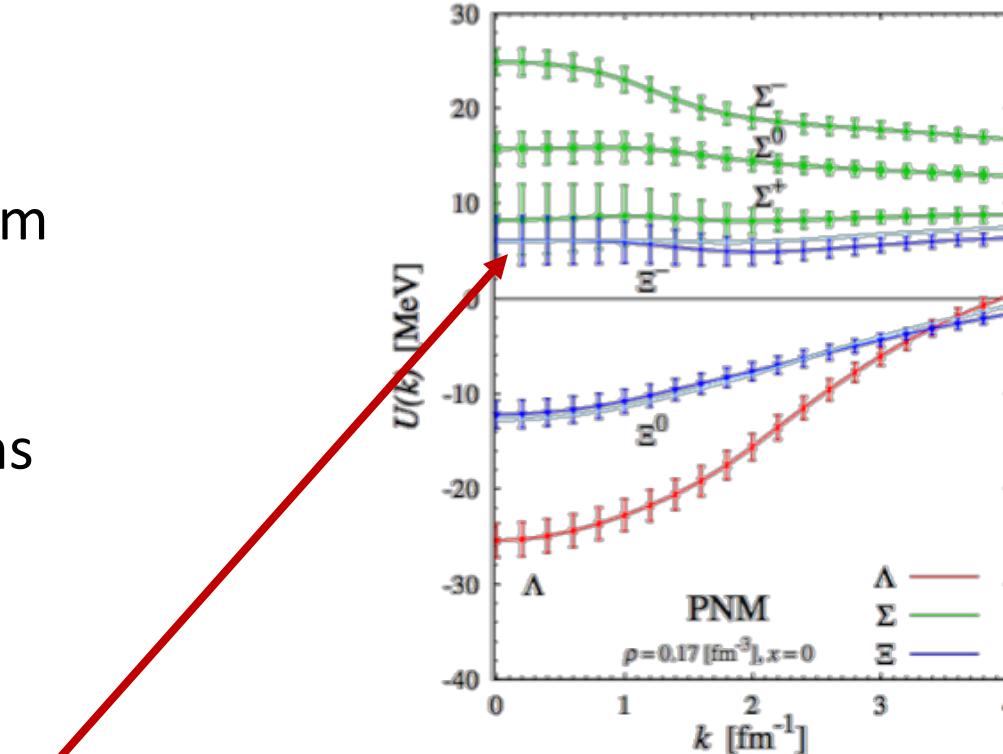


- Comparison HAL-QCD/ $\chi$ EFT from recent work by Haidenbauer and Meissner ([arxiv:1810.04883](https://arxiv.org/abs/1810.04883))

# Single particle potentials from Lattice

- NS environment  $\Rightarrow$  Pure Neutron Matter
- Tested HAL-QCD potential in vacuum with ALICE  $\Rightarrow$  Brueckner-Hartree-Fock many-body calculations  $\Rightarrow U_Y$  single-particle potential of hyperons in nucleonic matter
- At saturation density in PNM:

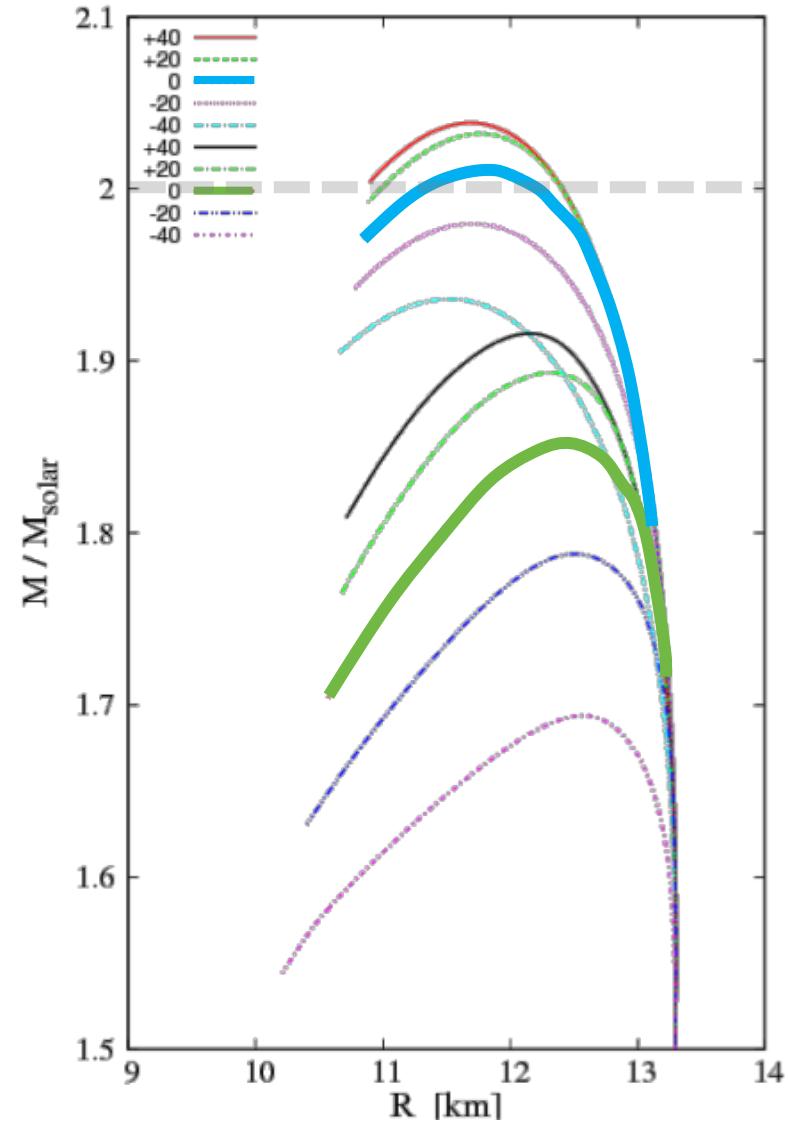
**$U_{\Xi}$  slightly repulsive**



HAL-QCD Collaboration, arXiv:1809.08932 (2018)

# Consequences for Neutron Stars

- Single-hyperon potential fundamental ingredient in EoS for NSs
- **Repulsive interaction  $\Rightarrow$  Production of  $\Xi^-$  pushed to higher densities  $\Rightarrow$  stiffer EoS, higher masses**



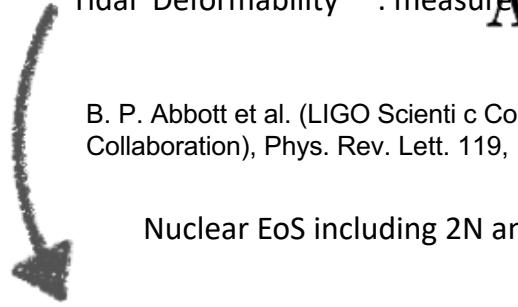
Weissborn et al., NPA881 (2012) 62-77

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# Recent Constraints from Gravitational Waves

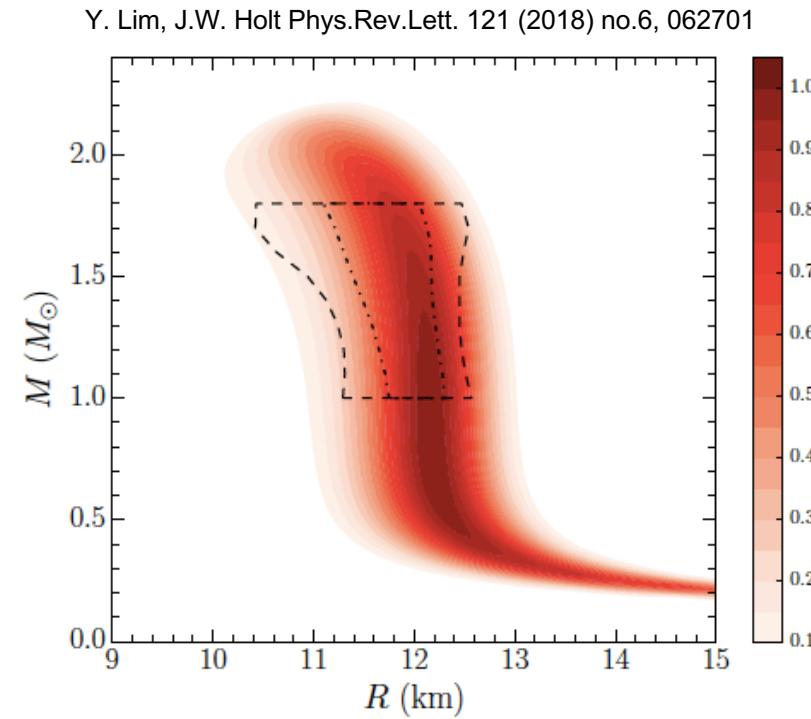
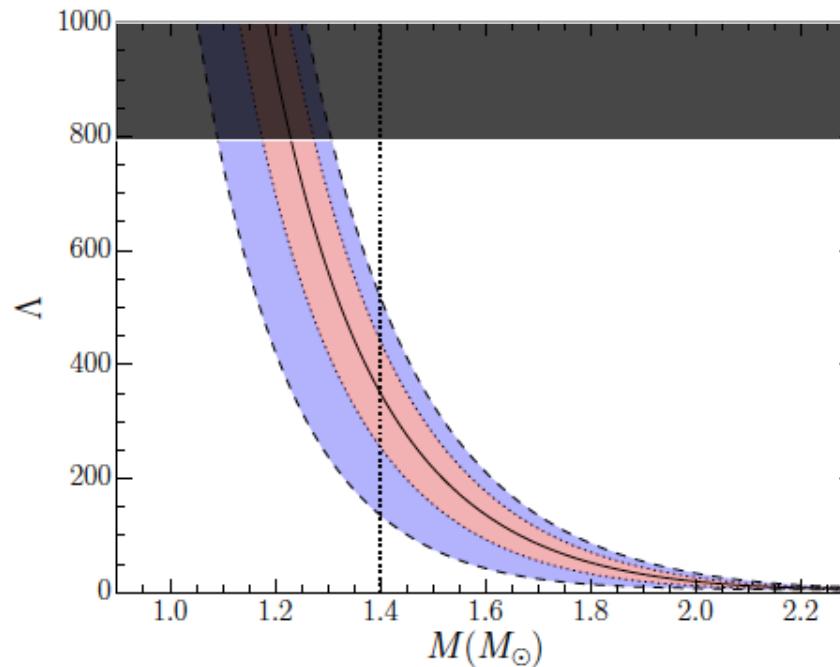
not constraining really very much so far...

Tidal Deformability : measurement of the NS deformability, upper limit extracted from GW170817

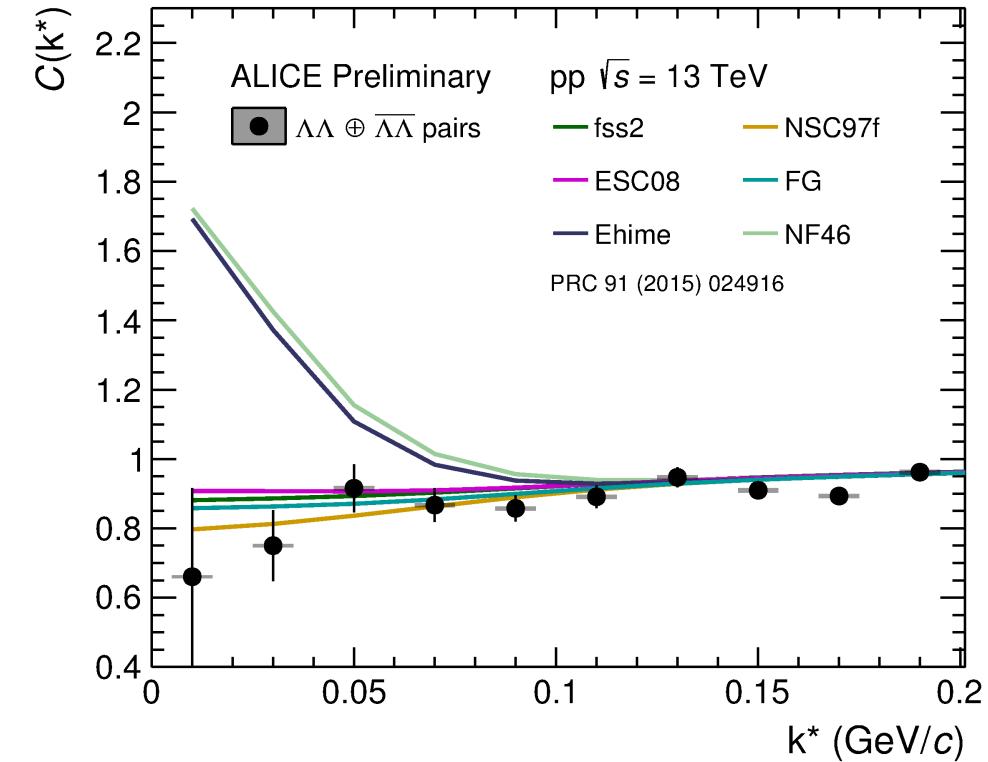
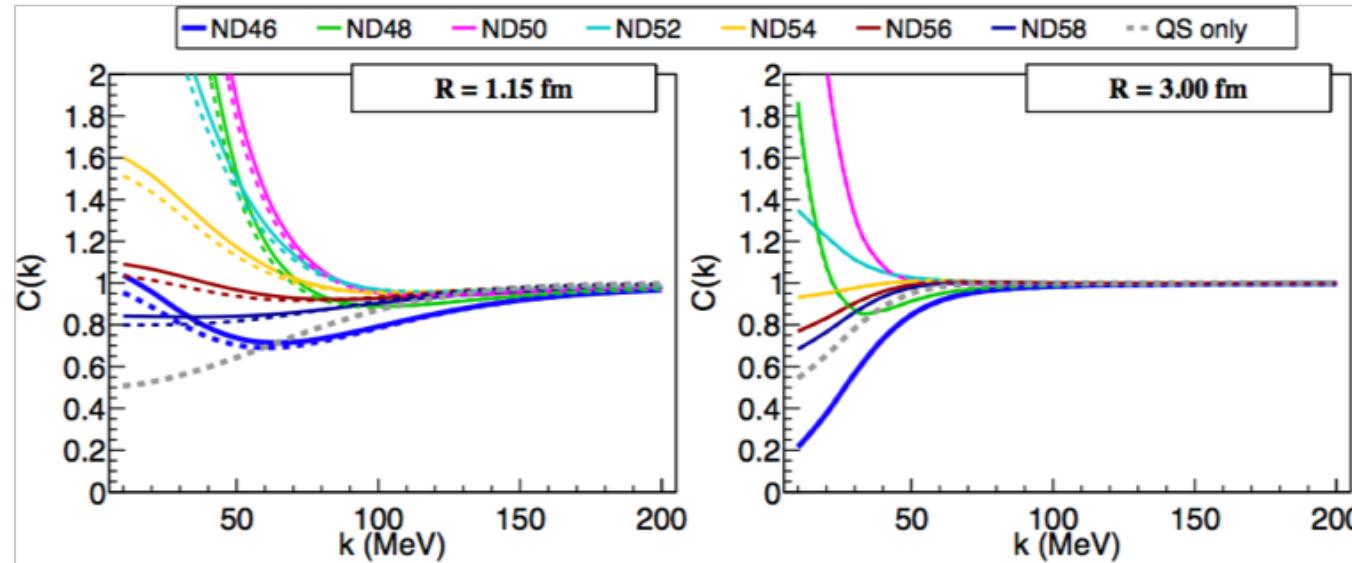


B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration), Phys. Rev. Lett. 119, 161101 (2017).

Nuclear EoS including 2N and 3N interactions are consistent with these constraints



# $\Lambda$ - $\Lambda$ Correlations: Predictions with Lednicky

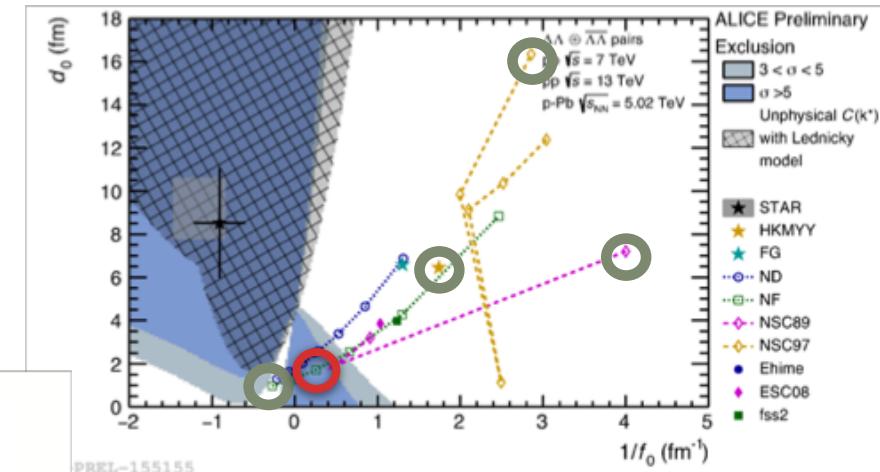
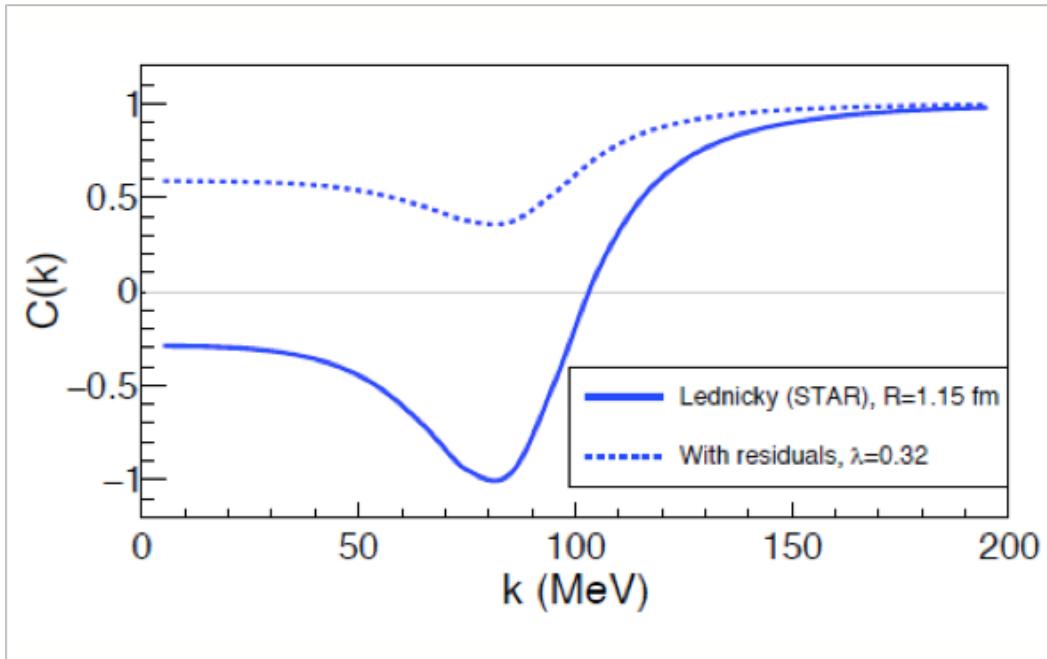


ALICE-PREL-144881

- Curves represent different points in the  $\Lambda$ - $\Lambda$  exclusion plot
- For scattering parameters in the region  $a_0 > 0$  the correlation function is not sensitive

# Non Physical Region and STAR parameters

- Lednicky model breaks down for small radii, large effective ranges and negative scattering parameters
- It shows although at most 10% deviation from CATS in the region of positive scattering parameters



# The unique opportunity of small sources

