

Equation of state for hyperonic nuclear matter and its application to compact astrophysical objects

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Outline

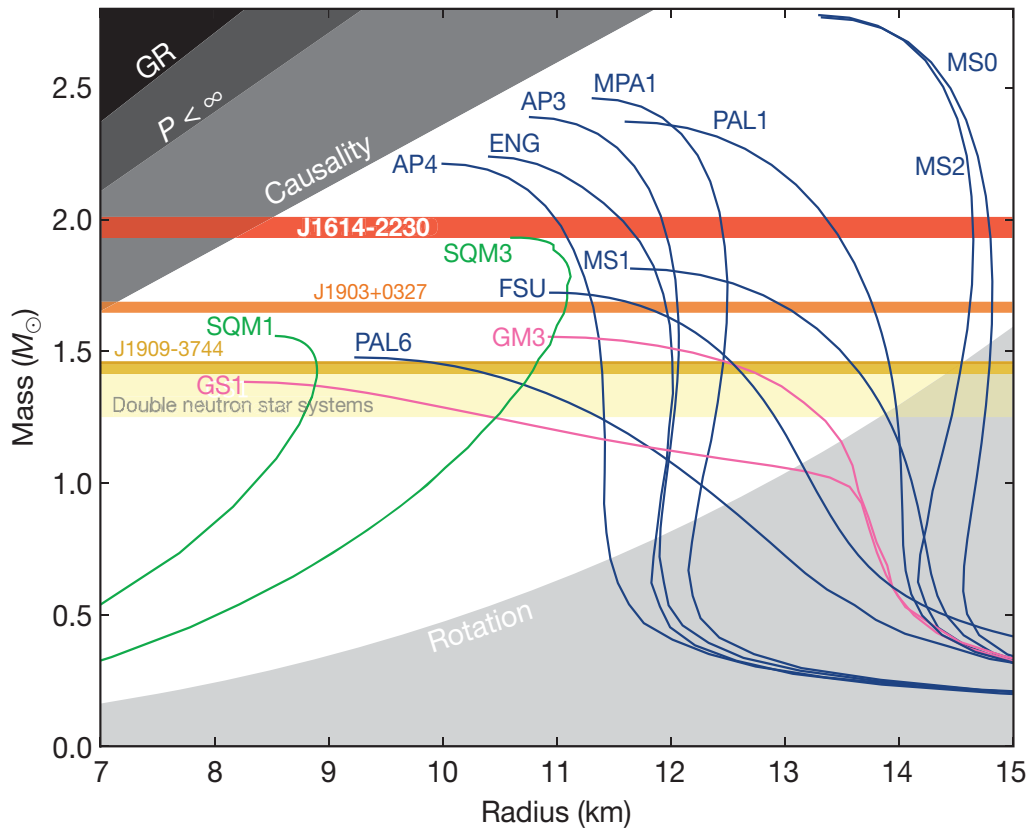
- 1 : Introduction
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International workshop on “Hadron structure and interaction in dense matter”
@ KEK Tokai campus, November 12, 2018

1. Introduction

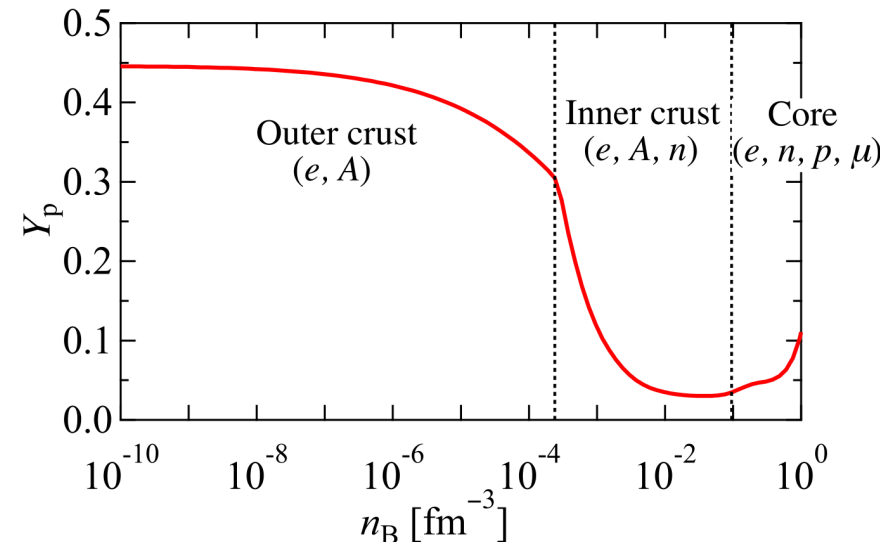
*Neutron star structure is governed by
the nuclear equation of state (EOS) at zero temperature.*

Neutron Stars: Stiffness (EOS at 0 MeV) \Leftrightarrow Self-gravity



Mass-radius relation of cold neutron stars

P. B. Demorest et al., NATURE 467 (2010)

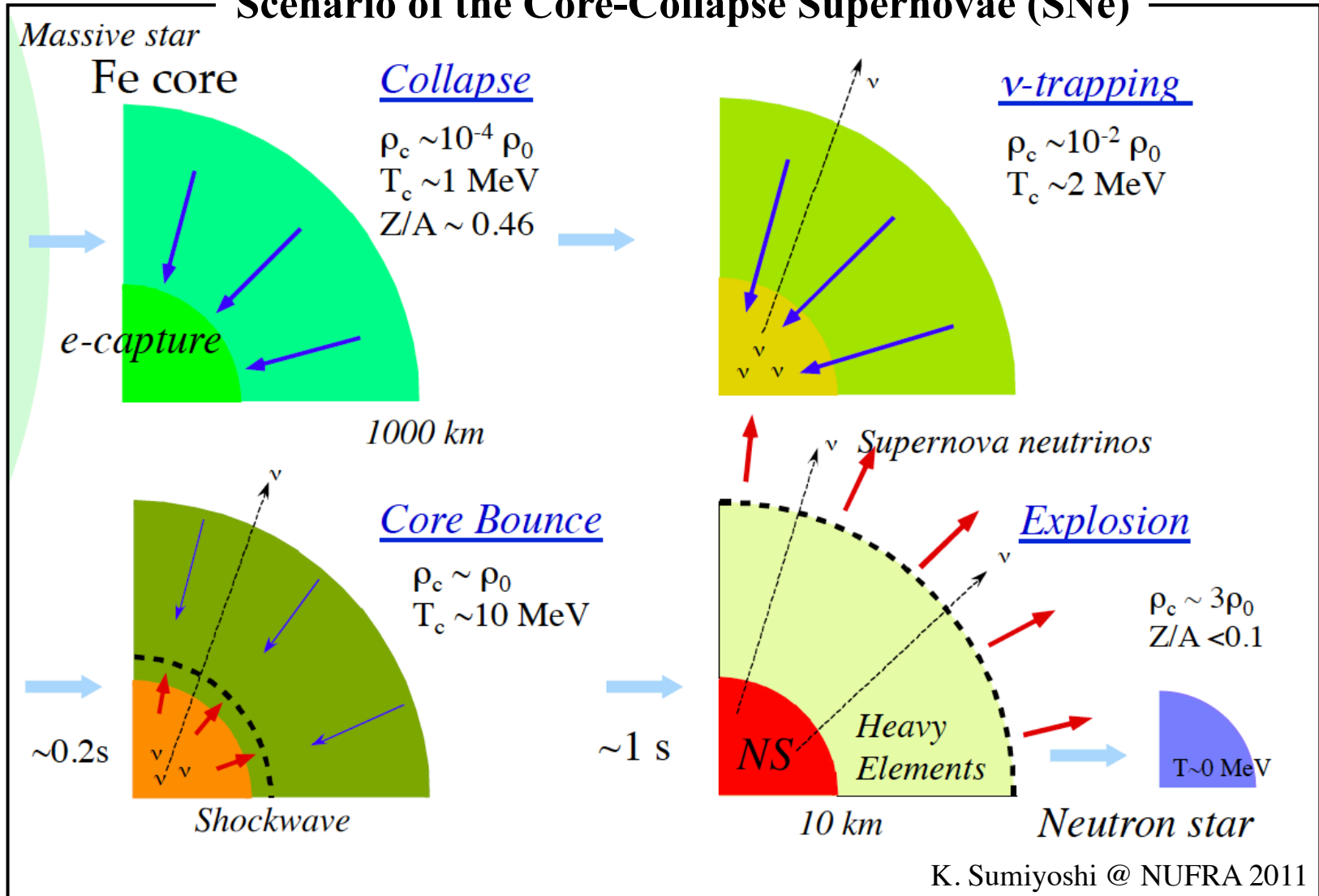


Phase diagram of cold nuclear matter

Nuclear EOS and Core-Collapse Supernovae

Nuclear EOS at finite temperature is one of the crucial ingredients
for the numerical simulations of *Core-Collapse Supernovae*.

Scenario of the Core-Collapse Supernovae (SNe)

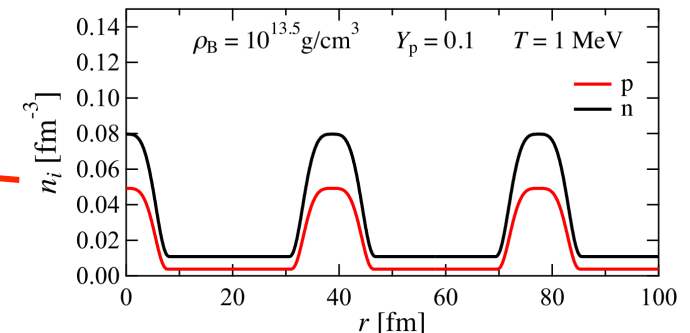
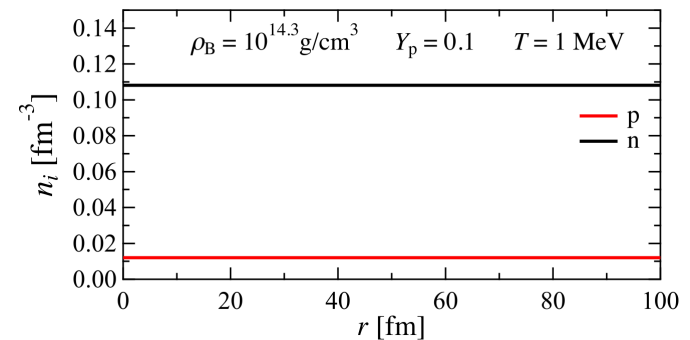
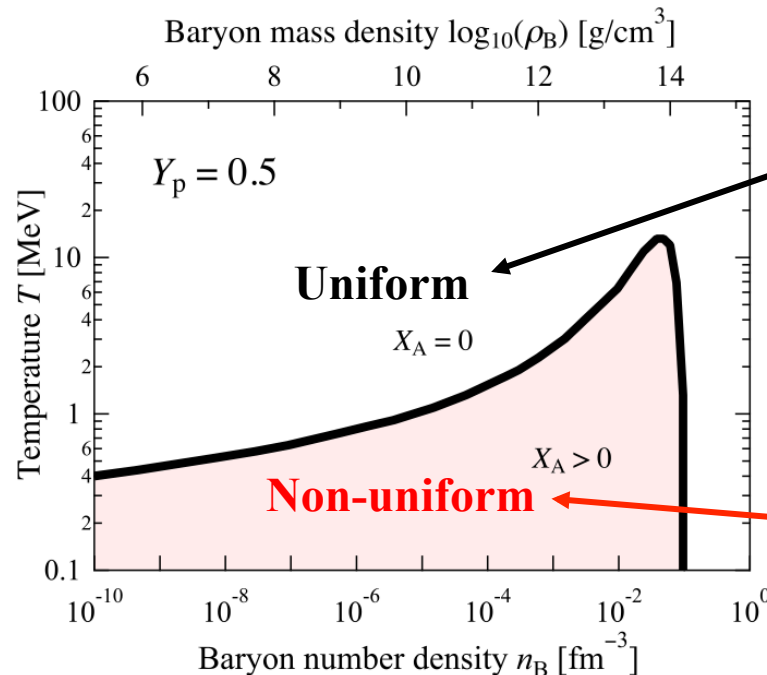


Nuclear EOS for supernova simulations

- SN-EOS should provide thermodynamic quantities in the wide ranges.

- Temperature T : $0 \leq T \leq 100$ MeV
- Density ρ : $10^{5.1} \leq \rho_B \leq 10^{16.0} \text{ g/cm}^3$
- Proton fraction Y_p : $0 \leq Y_p \leq 0.65$

- SN matter contains uniform and non-uniform phases.



Phase diagram of nuclear matter [based on HT *et al.*, NPA 961 (2017) 78]

Current status of SN-EOS with hyperons

Nuclear Interaction	n_{sat} (fm ⁻³)	BE/A (MeV)	K (MeV)	Q ($\frac{\text{MeV}}{\text{fm}^3}$)	J (MeV)	L (MeV)	type of int.	used in
SKa	0.155	16.0	263	-300	32.9	74.6	Skyrme	H&W
LS180	0.155	16.0	180	-451	28.6	73.8	Skyrme	LS180
LS220	0.155	16.0	220	-411	28.6	73.8	Skyrme	LS220, LS220A, LS220 π
LS375	0.155	16.0	375	176	28.6	73.8	Skyrme	LS375
TMA	0.147	16.0	318	-572	30.7	90.1	RMF	HS(TMA)
NL3	0.148	16.2	272	203	37.3	118.2	RMF	SHT, HS(NL3)
FSUgold	0.148	16.3	230	-524	32.6	60.5	RMF	SHO(FSU1.7), HS(FSUgold)
FSUgold2.1	0.148	16.3	230	-524	32.6	60.5	RMF	SHO(FSU2.1)
IUFSU	0.155	16.4	231	-290	31.3	47.2	RMF	HS(IUFSU)
DD2	0.149	16.0	243	169	31.7	55.0	RMF	HS(DD2), BHBA, BHBA ϕ
SFH ϕ	0.158	16.2	245	-468	31.6	47.1	RMF	SFH ϕ
SFH χ	0.160	16.2	239	-457	28.7	23.2	RMF	SFH χ
TM1	0.145	16.3	281	-285	36.9	110.8	RMF	STOS, FYSS, HS(TM1), STOSA, STOSY, STOSY π , STOS π , STOS π Q, STOSQ, STOSB139, STOSB145, STOSB155, STOSB162, STOSB165

SN-EOS list by M. Hempel

Hyperon EOS

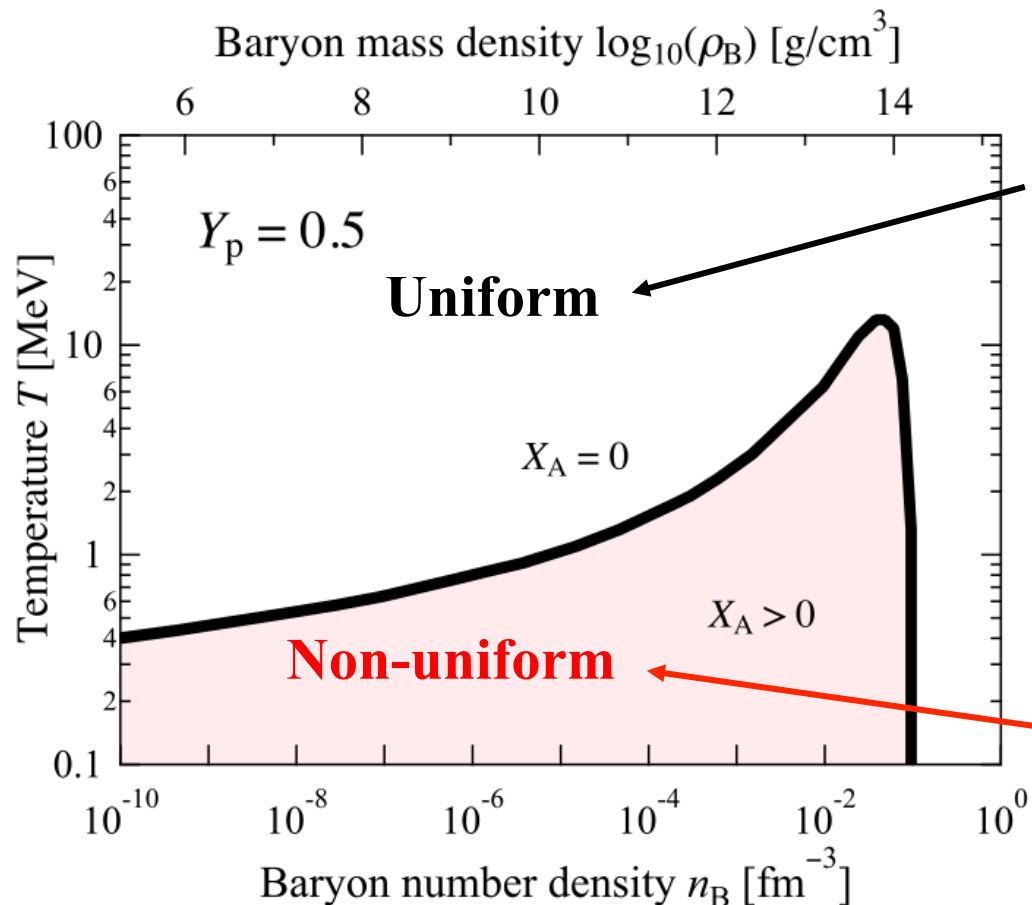
There is no SN-EOSs based on the microscopic many-body theory.

- Shen EOS with Λ , Σ , Ξ [$M_{\text{max}} = 1.67 M_{\odot}$] (C. Ishizuka et al., JPG 35 (2008) 085201)
- Shen EOS with Λ [$M_{\text{max}} = 1.75 M_{\odot}$] (H. Shen et al., APJS 197 (2011) 20)
- LS EOS with Λ [$M_{\text{max}} = 1.91 M_{\odot}$] (M. Oertel et al., PRC 85 (2012) 055806)
- DD2 EOS with Λ [$M_{\text{max}} = 2.11 M_{\odot}$] (S. Banik et al., APJS 214 (2014) 22)
- DD2 EOS with Λ , Σ , Ξ [$M_{\text{max}} = 2.04 M_{\odot}$] (M. Marques et al., PRC 96 (2017) 045806)

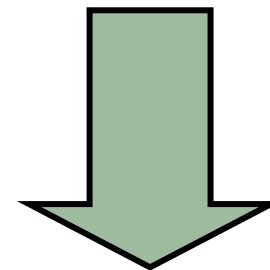
New EOS table for core-collapse simulations

<http://www.np.phys.waseda.ac.jp/EOS/>

(HT, K. Nakazato, Y. Takehara, S. Yamamuro, H. Suzuki, M. Takano, NPA961 (2017) 78)



1: Cluster variational method
with AV18 + UIX potentials



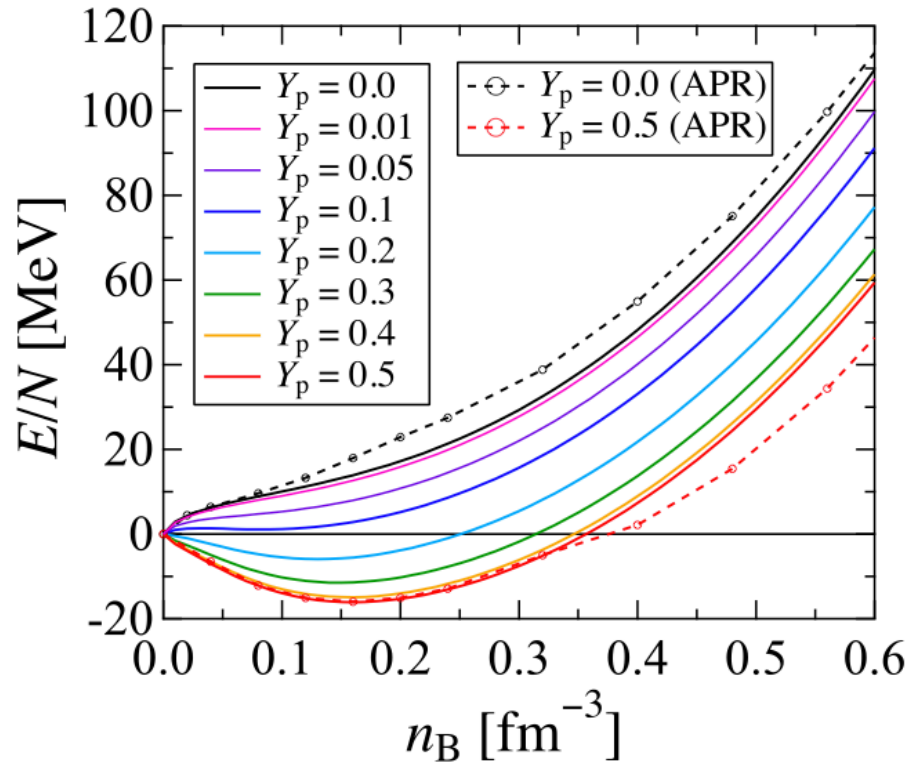
2: Thomas-Fermi calculation
for non-uniform matter

*We aim to extend the microscopic EOS table
to consider **Λ hyperon mixing**.*

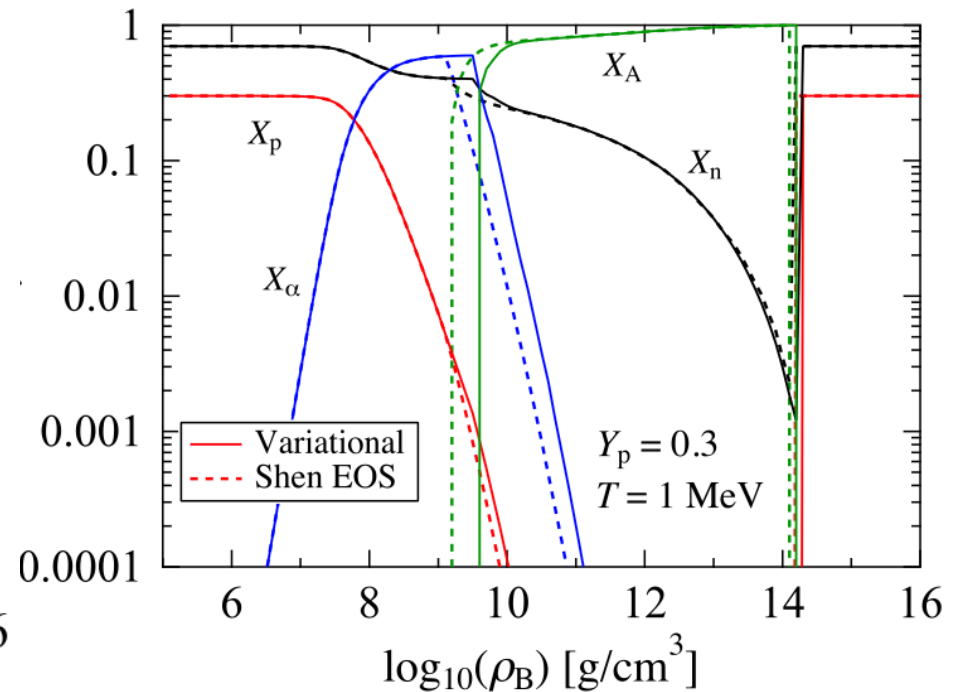
2. Supernova EOS with realistic nuclear forces

Uniform EOS: Cluster variational method with AV18 + UIX potentials

Non-uniform EOS: Thomas-Fermi method



Energy per nucleon for uniform matter



Particle fractions for non-uniform matter

$n_0[\text{fm}^{-3}]$	$E_0[\text{MeV}]$	$K[\text{MeV}]$	$E_{\text{sym}}[\text{MeV}]$
0.16	-16.1	245	30.0

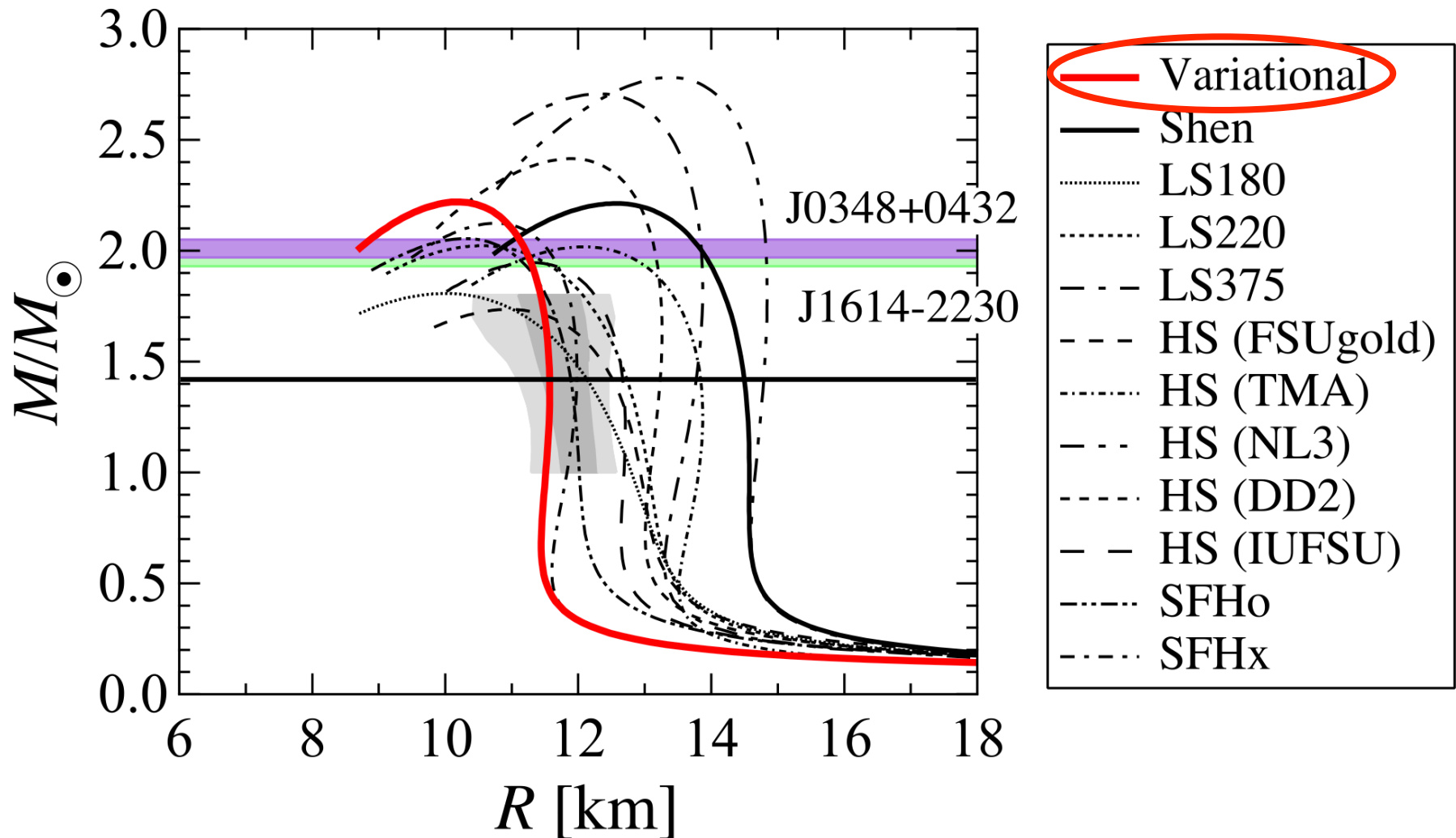
Our EOS : HT and M. Takano, NPA 902 (2013) 53

APR : A. Akmal, V. R. Pandharipande, D. G. Ravenhall,
PRC 58 (1998) 1804

Shen EOS : APJS 197 (2011) 20

Application to Neutron Star

Mass-Radius relation of neutron stars



J0348+0432: Science 340 (2013) 1233232

J1614-2230: Nature 467 (2010) 1081

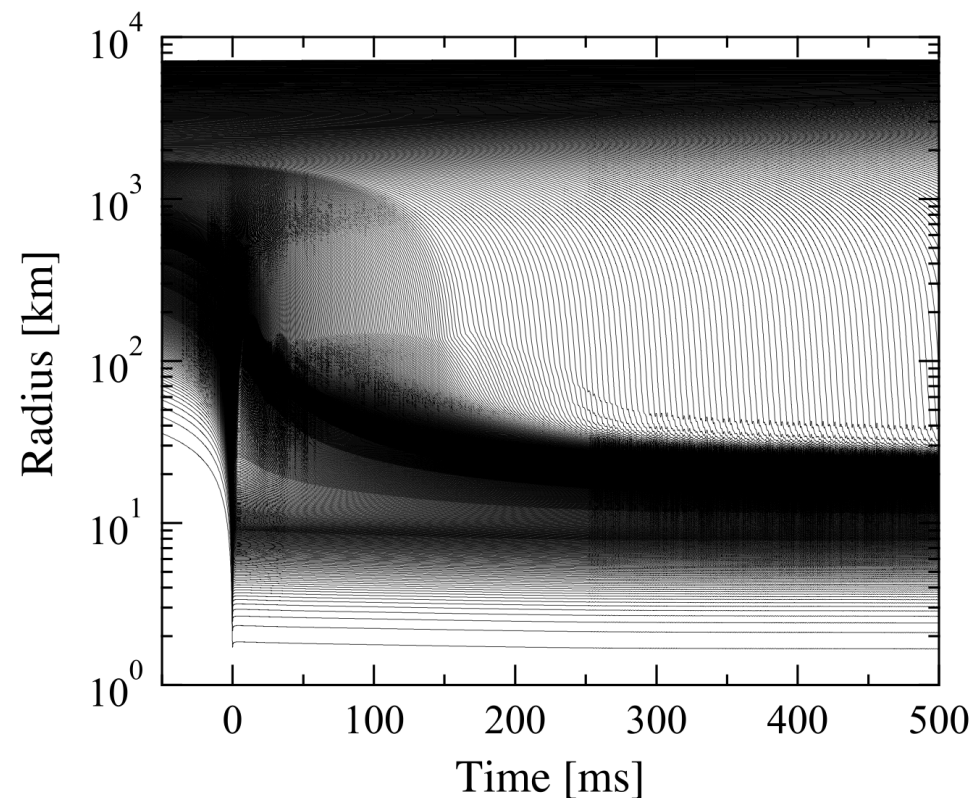
Shaded region is the observationally suggested region by Steiner et al.

Application to Core-Collapse Supernovae

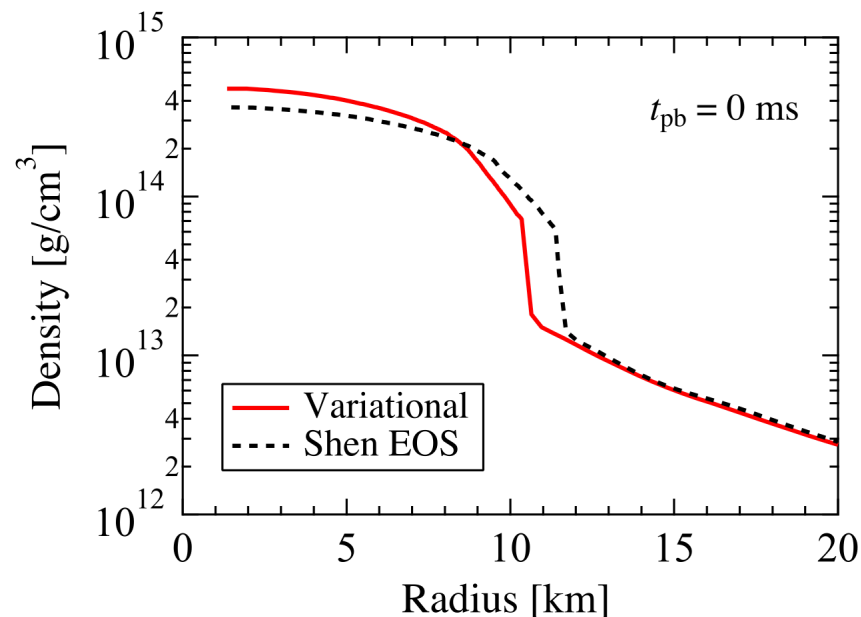
1D neutrino-radiation hydrodynamics simulations

Progenitor: Woosley Weaver 1995, $15M_{\odot}$ *Astrophys. J. Suppl.* 101 (1995) 181

SN simulation numerical code: K. Sumiyoshi, et al., *NPA* 730 (2004) 227



Radial trajectories of mass elements



Central density: 0.30 fm^{-3}

Temperature: $\sim 10 \text{ MeV}$

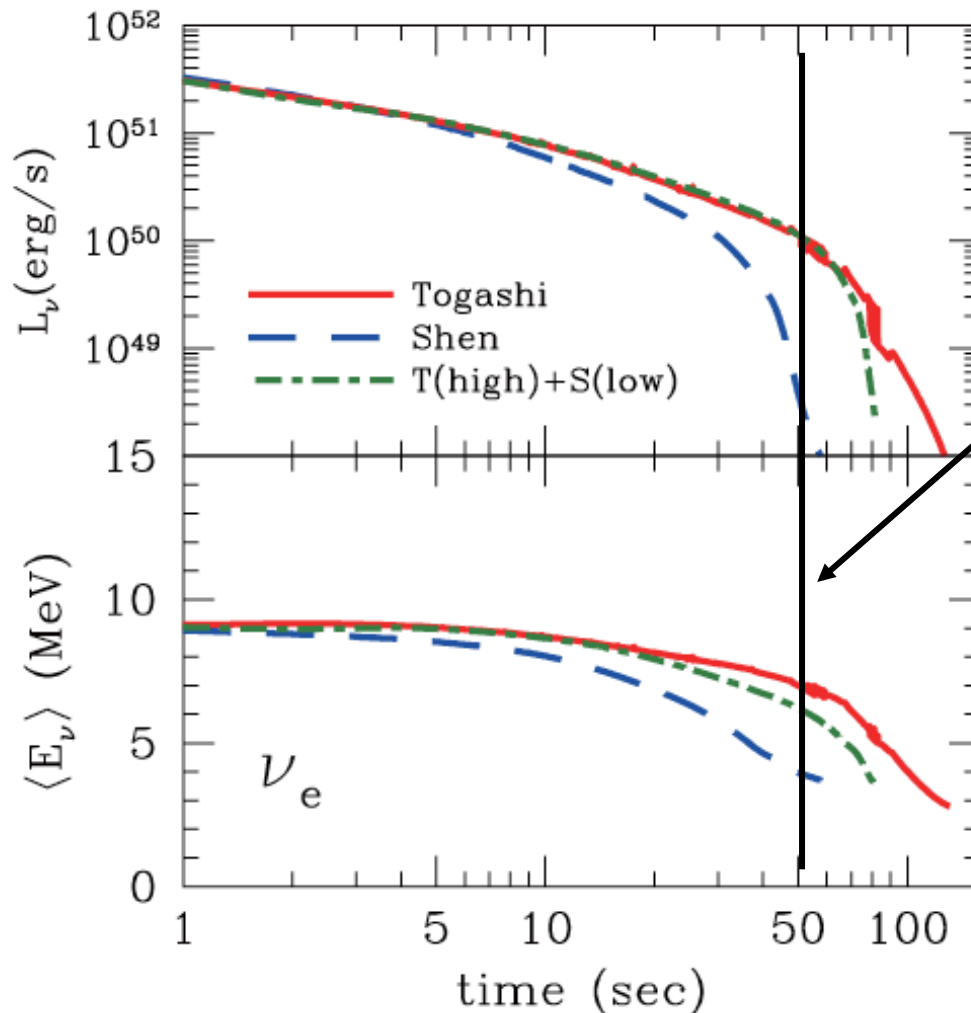
Proton fraction: ~ 0.3

Application to Proto-Neutron Star Cooling

K. Nakazato, H. Suzuki, and HT, Phys. Rev. C 97 (2018) 035804

1D neutrino-radiation hydrodynamics simulations (until 300 ms)

→ Quasi-static evolutionary calculation of PNS cooling



Central density: 0.47 fm^{-3}
Temperature: $\sim 10 \text{ MeV}$
Proton fraction: ~ 0.1

3. Supernova EOS with Λ hyperons

Two-body Hamiltonian

$$H_2 = -\sum_i \frac{\hbar^2}{2m_i} \nabla_i^2 + \sum_{i < j} [V_{ij}^{\text{NN}} + V_{ij}^{\Lambda\text{N}} + V_{ij}^{\Lambda\Lambda}]$$

V_{ij}^{NN} : Argonne v18 (AV18) potential

Hyperon Two-Body Central Potentials

$V_{ij}^{\Lambda\text{N}}$: Λ -Nucleon (N) potential (E. Hiyama et al., PRC 74 (2006) 054312)

- Constructed so as to reproduce the experimental binding energies of *light Λ hypernuclei with the Gaussian expansion method.*

$V_{ij}^{\Lambda\Lambda}$: Λ - Λ potential (E. Hiyama et al., PRC 66 (2002) 024007)

- the experimental double- Λ binding energy from ${}_{\Lambda\Lambda}^6\text{He}$ (*NAGARA event*)

E_2 : Expectation value of H_2 in *the two-body cluster approximation*

Expectation value of the Hamiltonian

Jastrow wave function

$$\Psi = \text{Sym} \left[\prod_{i < j} f_{ij} \right] \Phi_F$$

Φ_F : The Fermi-gas wave function

Correlation function:

$$f_{ij} = \sum_{\mu, p, s} [f_{Cps}^{\mu}(r_{ij}) + s f_{Tp}^{\mu}(r_{ij}) S_{Tij} + s f_{SOp}^{\mu}(r_{ij}) (\mathbf{L}_{ij} \cdot \mathbf{s})] P_{psij}^{\mu}$$

p : parity

s : two-particle total spin

μ : particle pair

Cluster-expansion

$$\frac{\langle H_2 \rangle}{N} = \frac{1}{N} \frac{\langle \Psi | H_2 | \Psi \rangle}{\langle \Psi | \Psi \rangle} = \frac{\langle H_2 \rangle_2}{N} + \frac{\langle H_2 \rangle_3}{N} + \dots$$

E_2 is the expectation value of H_2 in *the two-body cluster approximation* with the Jastrow wave function.

Three-Body Energy

Nuclear EOS (previous study)

$$\frac{E_3}{N} = \underbrace{\frac{1}{N} \left\langle \sum_{i < j < k}^N [\alpha V_{ijk}^R + \beta V_{ijk}^{2\pi}] \right\rangle_F}_{\text{Modified expectation value of } H_3 \text{ with } \Phi_F} + \underbrace{\gamma n_B^2 e^{-\delta n_B} [1 - (1 - 2Y_p)^2]}_{\text{Correction term}}$$

Repulsive part of UIX pot. is extended to the potential for Hyperon TBF



$$V_{ijk}^R = \sum_{\mu} \alpha^{\mu} V_{ijk}^R P_{ijk}^{\mu} \quad (\mu = \text{NNN}, \Lambda\text{NN}, \Lambda\Lambda\text{N}, \Lambda\Lambda\Lambda)$$

P_{ijk}^{μ} : Three-particle projection operator

α^{NNN} : we use the value in the nuclear EOS

(Saturation properties + Thomas-Fermi calculation for atomic nuclei)

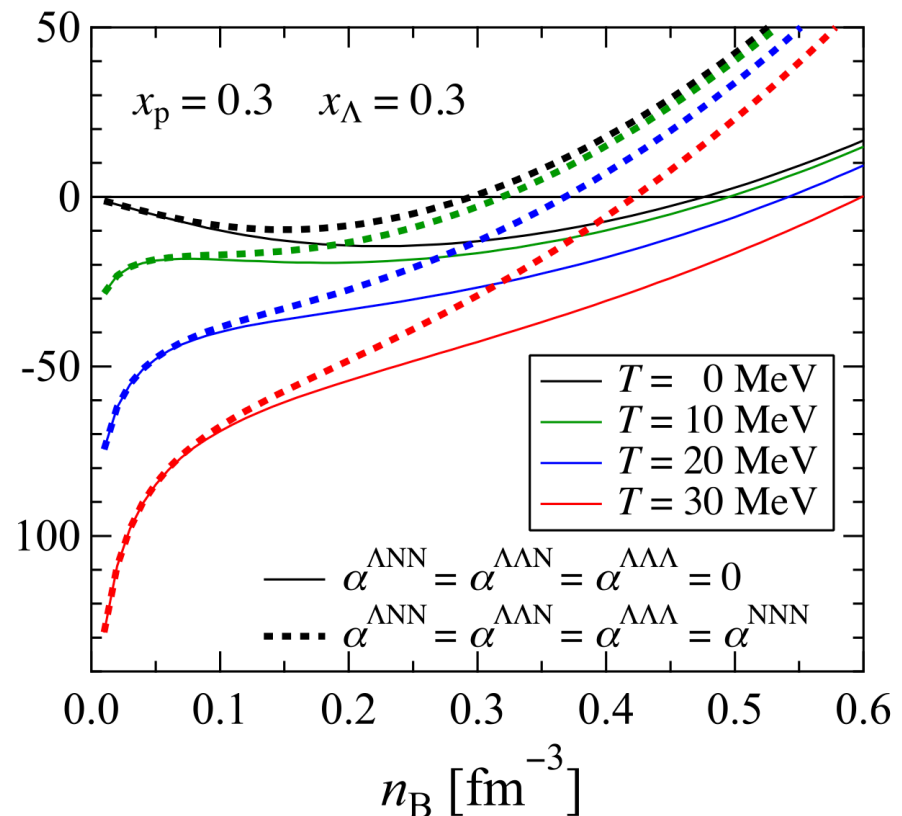
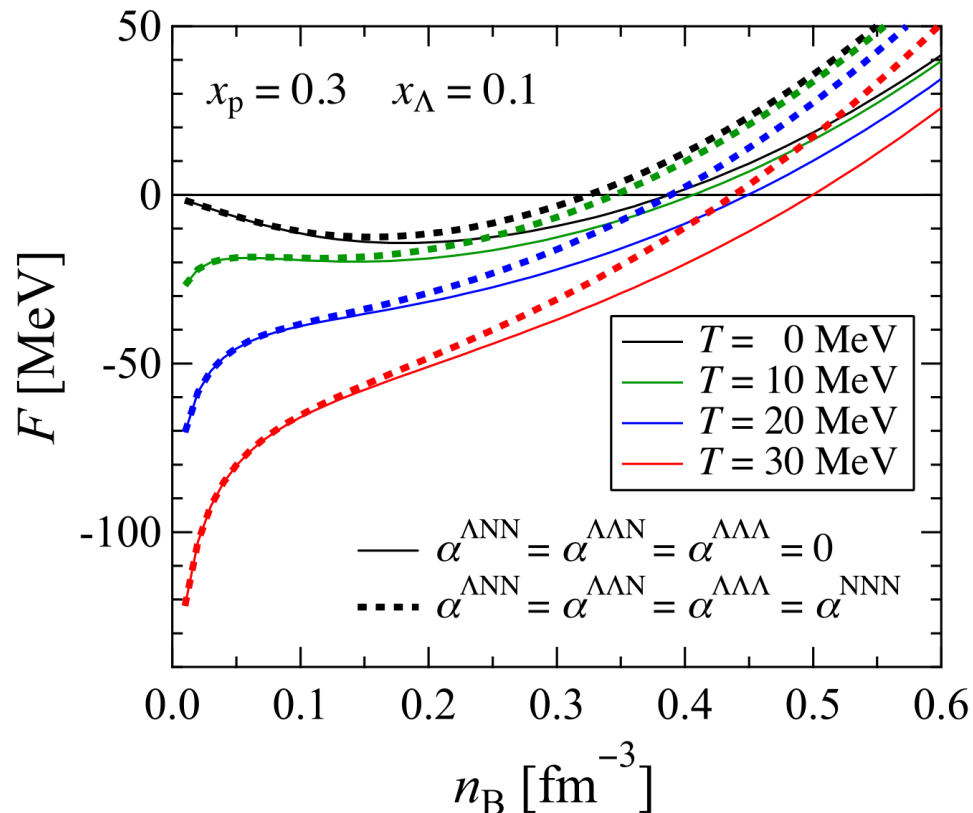
$\alpha^{\Lambda\text{NN}}, \alpha^{\Lambda\Lambda\text{N}}, \alpha^{\Lambda\Lambda\Lambda}$: Free parameters ($0 \leq \alpha^{\mu} \leq \alpha^{\text{NNN}}$)

Free energy for Λ hyperon matter

Total energy per baryon: $E(n_B, x_p, x_\Lambda) = E_2 + E_3$

Baryon number density : $n_B = n_p + n_n + n_\Lambda$ Particle fraction: $x_i = n_i/n_B$ ($i = p, \Lambda$)

The prescription by Schmidt and Pandharipande is employed
to obtain the free energy *at finite temperature*.



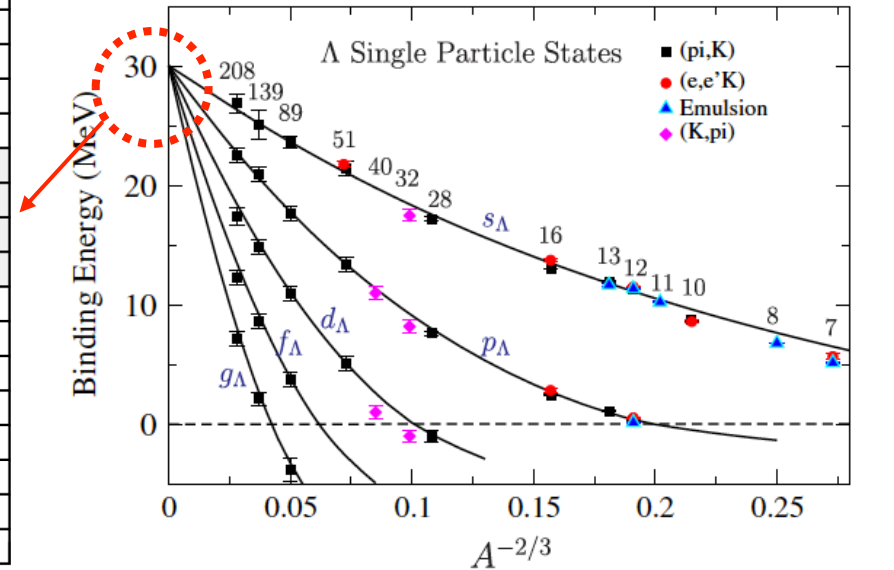
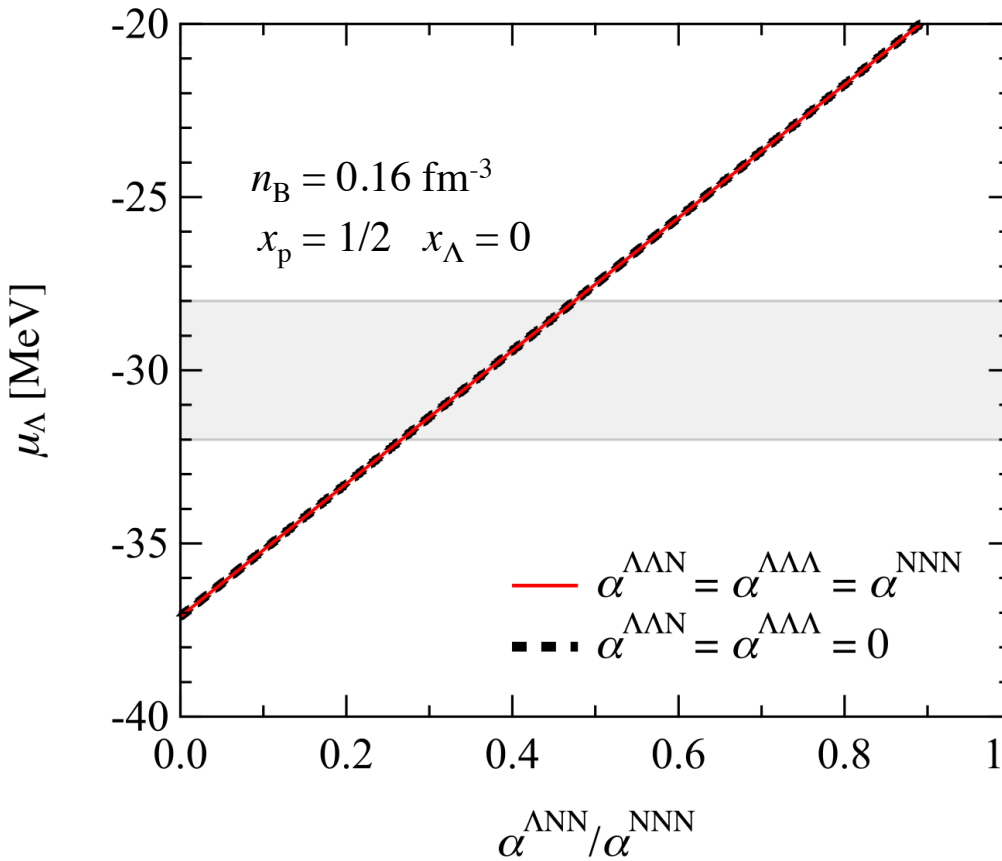
Single particle potential for Λ particle

Total energy per baryon: $E(n_B, x_p, x_\Lambda) = E_2 + E_3$

$$\mu_\Lambda(n_B, x_p, x_\Lambda) = [\partial(n_B E)/\partial n_\Lambda]_{n_n, n_p}$$

Baryon number density : $n_B = n_p + n_n + n_\Lambda$

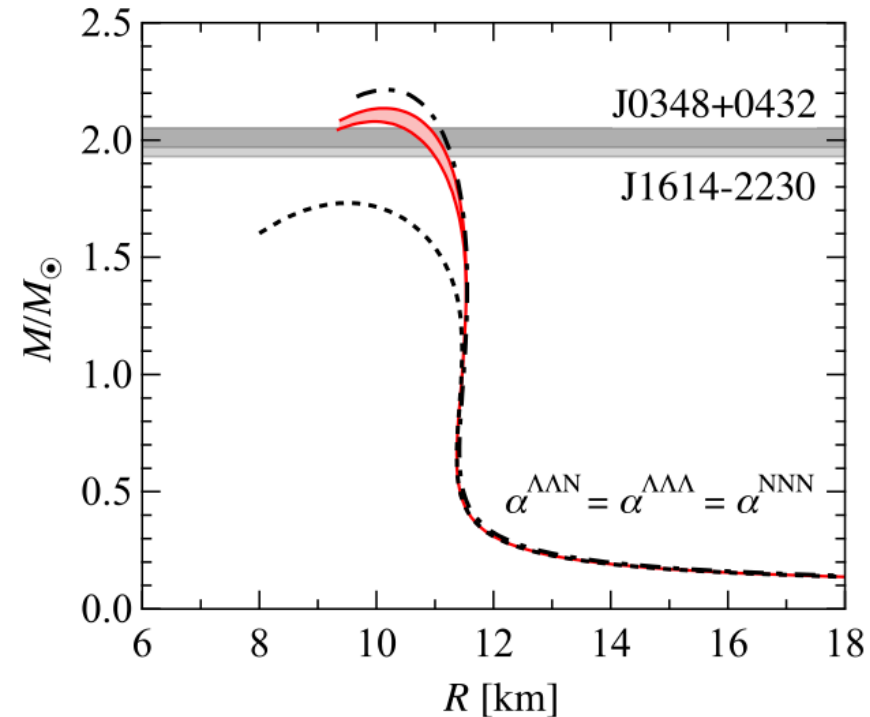
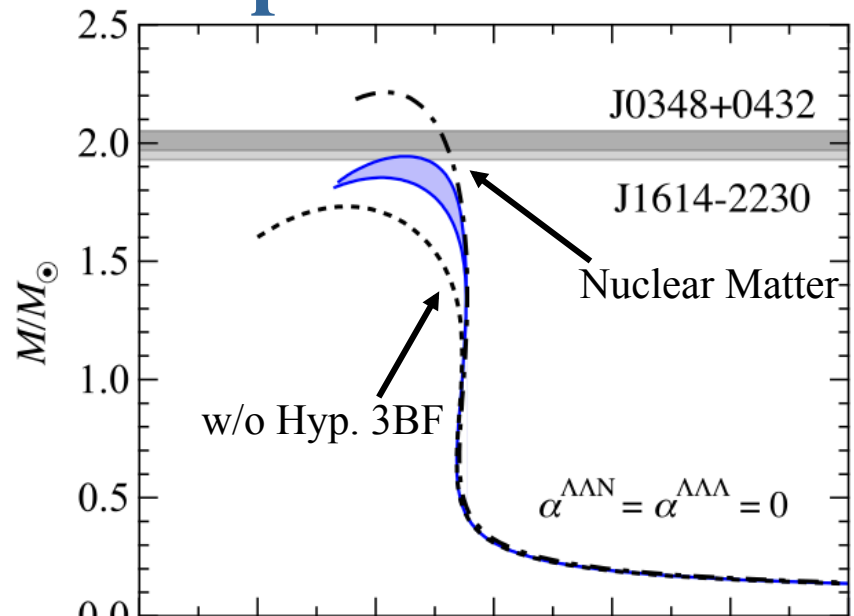
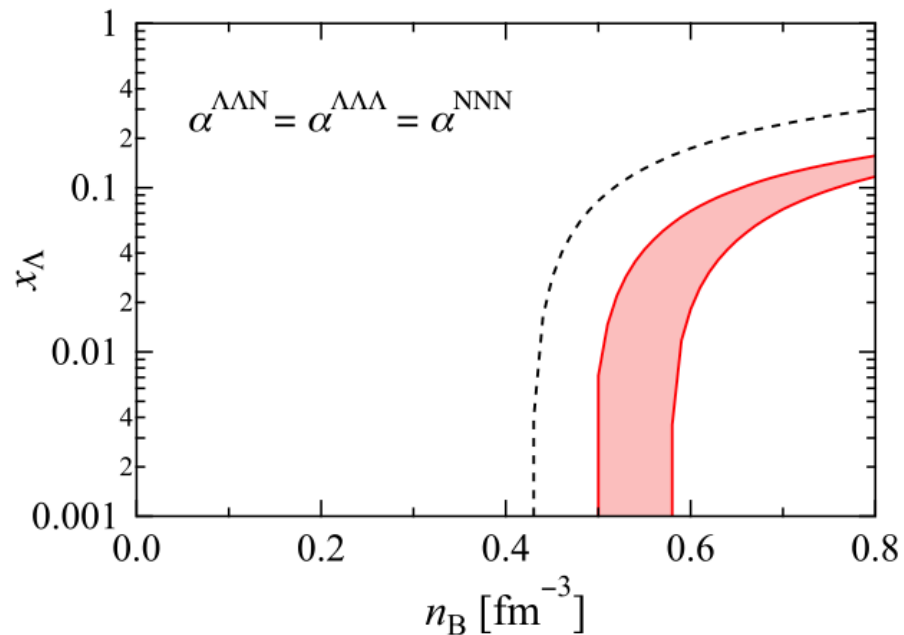
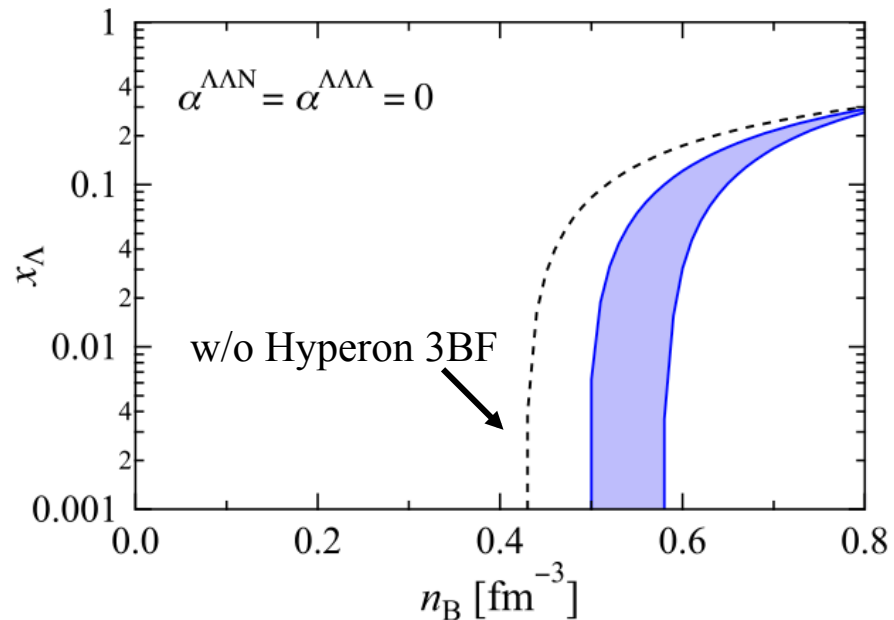
Particle fraction: $x_i = n_i/n_B$ ($i = p, \Lambda$)



Energy levels of the Λ single-particle major shells
(Rev. Mod. Phys. 88 (2016) 035004)

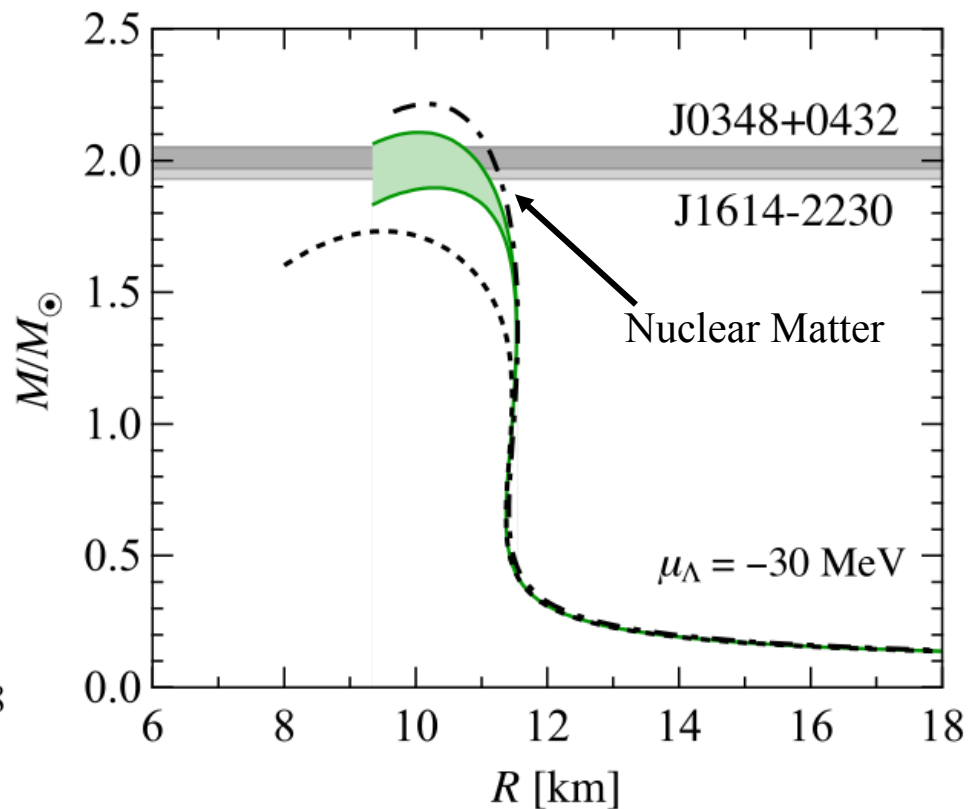
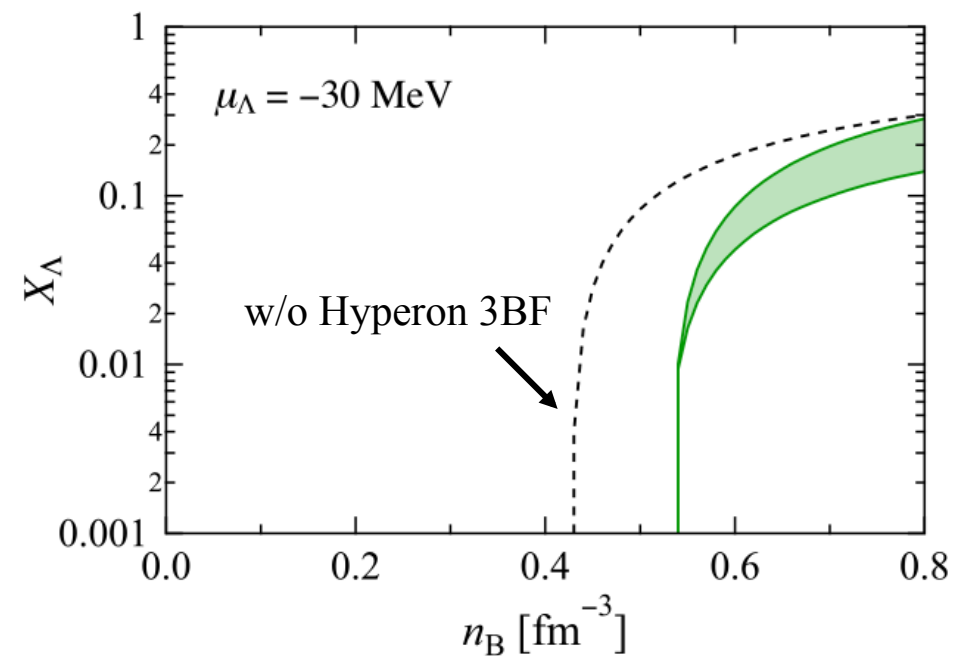
$$-32 \text{ MeV} \leq \mu_\Lambda \leq -28 \text{ MeV} \rightarrow 0.266 \leq \alpha^{\Lambda NN}/\alpha^{NNN} \leq 0.475$$

4. Application to Compact Stars



$\Lambda\Lambda N$ and $\Lambda\Lambda\Lambda$ three-body forces

- $\mu_\Lambda = -30 \text{ MeV}$ ($\alpha^{\Lambda NN} = 0.370 \alpha^{\text{NNN}}$)
- $0 \leq \alpha^{\Lambda\Lambda N} = \alpha^{\Lambda\Lambda\Lambda} \leq \alpha^{\text{NNN}}$



$\Lambda\Lambda N$ and $\Lambda\Lambda\Lambda$ three-body force:

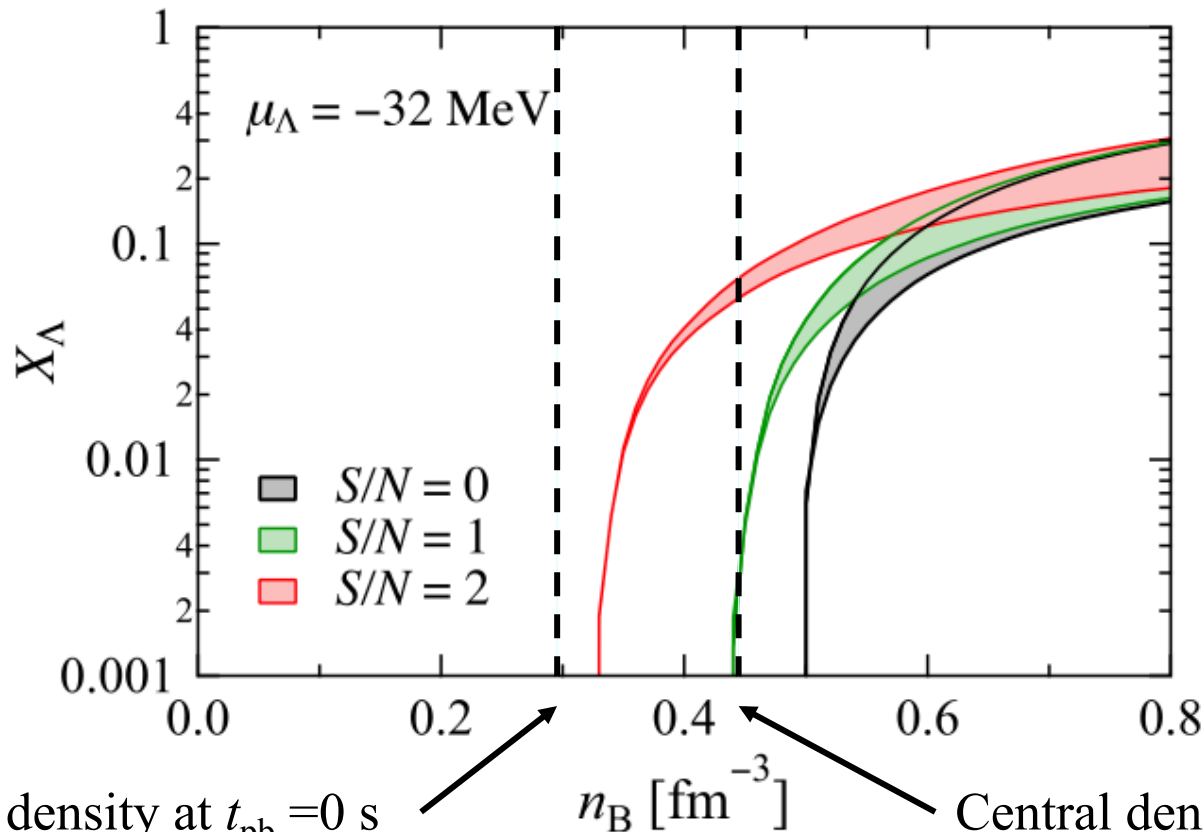
affect on the maximum mass of neutron stars (Important for HYPERON PUZZLE!?)

Hyperon mixing in supernova matter

Supernova matter

- Charge neutral and Isentropic matter (The entropy per baryon $S \sim 1-2$)
- Neutrino-free β -stable matter

$$(0 \leq \alpha^{\Lambda\Lambda N} = \alpha^{\Lambda\Lambda\Lambda} \leq \alpha^{\text{NNN}})$$



Central density at $t_{\text{pb}} = 0 \text{ s}$
(Core-collapse supernova)

Central density at $t_{\text{pb}} = 50 \text{ s}$
(PNS cooling)

Summary

We construct the EOS for nuclear matter including Λ hyperons at zero and finite temperatures by the variational method.

Application of the EOS to compact stars

- Λ NN three-body force:
affects on the single-particle potential and the onset density of Λ hyperon mixing
- $\Lambda\Lambda$ N and $\Lambda\Lambda\Lambda$ three-body forces:
affect on the maximum mass of neutron stars (Important for HYPERON PUZZLE!?)
- Λ hyperon fraction in supernova matter becomes larger at higher entropies.

Future Plans

- Construction of the EOS table for core-collapse simulations
- Taking into account mixing of other hyperons (Σ^- , Σ^0 , Σ^+ , Ξ^0 , Ξ^-)
- Employing more sophisticated baryon interactions (e.g. Nijmegen)

*→ we extend the cluster variational method to take into account coupled channels.
(Poster presentation in QNP 2018)*